

Cite this: *Chem. Sci.*, 2022, 13, 13442

All publication charges for this article have been paid for by the Royal Society of Chemistry

Ga and Zn increase the oxygen affinity of Cu-based catalysts for the CO_x hydrogenation according to *ab initio* atomistic thermodynamics†

Andreas Müller, ^a Aleix Comas-Vives ^{*bc} and Christophe Copéret ^{*a}

The direct hydrogenation of CO or CO₂ to methanol, a highly vivid research area in the context of sustainable development, is typically carried out with Cu-based catalysts. Specific elements (so-called promoters) improve the catalytic performance of these systems under a broad range of reaction conditions (from pure CO to pure CO₂). Some of these promoters, such as Ga and Zn, can alloy with Cu and their role remains a matter of debate. In that context, we used periodic DFT calculations on slab models and *ab initio* thermodynamics to evaluate both metal alloying and surface formation by considering multiple surface facets, different promoter concentrations and spatial distributions as well as adsorption of several species (O*, H*, CO* and CO₂*) for different gas phase compositions. Both Ga and Zn form an fcc-alloy with Cu due to the stronger interaction of the promoters with Cu than with themselves. While the Cu–Ga-alloy is more stable than the Cu–Zn-alloy at low promoter concentrations (<25%), further increasing the promoter concentration reverses this trend, due to the unfavoured Ga–Ga-interactions. Under CO₂ hydrogenation conditions, a substantial amount of O* can adsorb onto the alloy surfaces, resulting in partial dealloying and oxidation of the promoters. Therefore, the CO₂ hydrogenation conditions are actually rather oxidising for both Ga and Zn despite the large amount of H₂ present in the feedstock. Thus, the growth of a GaO_x/ZnO_x overlayer is thermodynamically preferred under reaction conditions, enhancing CO₂ adsorption, and this effect is more pronounced for the Cu–Ga-system than for the Cu–Zn-system. In contrast, under CO hydrogenation conditions, fully reduced and alloyed surfaces partially covered with H* and CO* are expected, with mixed CO/CO₂ hydrogenation conditions resulting in a mixture of reduced and oxidised states. This shows that the active atmosphere tunes the preferred state of the catalyst, influencing the catalytic activity and stability, indicating that the still widespread image of a static catalyst under reaction conditions is insufficient to understand the complex interplay of processes taking place on a catalyst surface under reaction conditions, and that dynamic effects must be considered.

Received 3rd June 2022
Accepted 18th September 2022

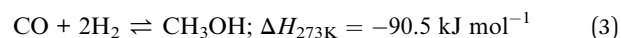
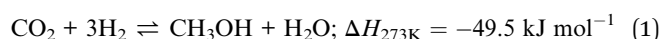
DOI: 10.1039/d2sc03107h

rsc.li/chemical-science

1 Introduction

Today, a large effort is directed at finding solutions to decrease the emission of carbon dioxide (CO₂) and its atmospheric concentration by conversion or sequestration processes since its atmospheric concentration directly relates to the observed climate change.¹ One of the possible routes towards a carbon-neutral society is based on the so-called methanol economy,^{2–4} where CO₂ is converted into methanol, a key chemical

intermediate and a fuel additive, *via* hydrogenation using “green” H₂, coming from renewable energy sources. This would lead to an overall decarbonisation of our economy (eqn (1)).^{5,6}



This reaction typically competes with the (reverse) water gas shift ((R)WGS) reaction (eqn (2)) as well as other reactions such as the hydrogenation of CO (eqn (3)) or the Sabatier (CO₂ methanation) reaction (eqn (4)). The choice of catalysts is crucial to drive the efficient use of hydrogen and the reactions towards the desired product, *e.g.* methanol.^{7,8}

^aDepartment of Chemistry and Applied Biosciences, ETH Zürich, 8093 Zurich, Switzerland. E-mail: coperet@ethz.ch; Tel: +41 44 633 93 94

^bInstitute of Materials Chemistry, TU Wien, 1060 Vienna, Austria. E-mail: aleix.comas@tuwien.ac.at

^cDepartament de Química, Universitat Autònoma de Barcelona, 08193 Cerdanyola del Vallès, Catalonia, Spain. E-mail: Aleix.Comas@uab.cat

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d2sc03107h>



In this context, some of the most promising catalysts are based on copper usually dispersed on or stabilized by oxide supports, e.g. Cu/ZnO/Al₂O₃. These catalysts were originally developed and are still used for the hydrogenation of CO to methanol. This process based on the conversion of syngas (mixture of H₂/CO at ratios between 2 : 1 and 3 : 1) is carried out in the presence of small quantities of CO₂ (a few %) at high temperatures (200 to 300 °C) and pressures (50 to 100 bar).^{9,10} The main advantage of such Cu-based catalysts is the absence of over-reduction of CO or CO₂ to hydrocarbons *via* the Sabatier reaction (4), which would be thermodynamically favoured and would waste large quantities of hydrogen in the feedstock to produce methane, a greenhouse gas and a molecule of lower value than methanol.¹¹

In recent years, large attention has been refocused on investigating these Cu-based catalysts for the hydrogenation of pure CO₂ or CO₂-rich gas feeds containing CO.^{12–54} This has led to the development of several processes towards large-scale application.^{55–59} Nevertheless, hydrogenation with CO₂-rich feeds usually suffers from product inhibition, reducing the overall activity of the catalyst.⁶⁰ Furthermore, low methanol selectivity is also observed in the low pressure range due to the competing RWGS-reaction.⁶¹ Finally, the production of large quantities of water causes catalyst deactivation, thus reducing the long-term catalyst stability.⁶²

This has led to refocused research efforts to increase the activity, selectivity and stability of Cu-based catalysts.^{63,64} Detailed mechanistic studies using Surface Organometallic Chemistry (SOMC) and thermolytic molecular precursors (TMP),^{65,66} among others, have highlighted the role of the interface and specific promoters between Cu nanoparticles and the oxide supports.^{67–70}

It was also shown that the presence of promoters like Ga or Zn can generate Cu–Ga- or Cu–Zn-alloy NPs upon treatment under H₂. The addition of these promoters also greatly increases methanol selectivity, in particular at high conversion and low pressure. Detailed *operando* X-ray Absorption Spectroscopy was able to show that under reaction conditions ($T = 230$ °C, 15 bar H₂, 5 bar CO₂), full and partial dealloying takes place for Ga and Zn in a partially reversible process upon switching from a CO₂/H₂ to a pure H₂ atmosphere, clearly showing that dynamic phenomena readily take place in these systems.^{71–74} Such process is reminiscent of what has been discussed for the Cu/ZnO/Al₂O₃ catalysts, further fuelling the debate regarding the nature of the catalytic active sites.^{75–81} Furthermore, it is consistent with the rich chemistry of Cu–Zn- and Cu–Ga-alloys, that display similar yet distinct alloying behaviour and phase diagrams (*vide infra*).^{82,83}

In parallel, computational chemistry has emerged as a powerful technique for understanding heterogeneous catalysts^{84–91} and also helped unravel the role of the support in methanol synthesis.^{68,92} We aim to understand from first principles calculations the state of Cu–Ga- and Cu–Zn-based alloys under a broad range of reaction conditions, ranging from a pure hydrogen atmosphere under which catalysts are typically prepared, to typical CO/CO₂ hydrogenation conditions. We will interrogate the stability of Cu–Ga- and Cu–Zn-alloys and their

propensity to interact with multiple relevant adsorbates (O*, H*, CO*, CO₂*) under these conditions, extending our previous work for pure Cu surfaces.⁹³ We will focus our study on the most stable facets of Cu-based alloy NPs promoted with Ga and Zn by modelling the gas phase composition and temperature *via ab initio* atomistic thermodynamics. Since promoters are usually the minor species, we will consider a fcc-solid solution of both alloys under reaction conditions.

2 Theoretical methods

2.1 Computational details

All Calculations were performed using the Vienna *Ab initio* Simulation Package (VASP)^{94–96} with a cut-off energy of 400 eV. The structures were optimized using the projector augmented wave (PAW) method (plane-wave basis set with pseudopotentials)^{97,98} using the official VASP pseudopotentials⁹⁹ and the Perdew–Burke–Ernzerhof (PBE) exchange–correlation functional.^{100,101} Default settings were used for VASP if not stated otherwise. Electronic occupancies were determined according to a Methfessel–Paxton scheme with an energy smearing of 0.2 eV. The convergence criteria for the electronic SCF-loop (EDIFF) was set to 10^{–5} eV. The ionic relaxation was considered converged when the norms of all the forces were smaller than 10 meV Å^{–1} (EDIFFG = –0.01). The ions were relaxed *via* the conjugate gradient algorithm (IBRION = 2). The Monkhorst–Pack *k*-point mesh was used to sample the Brillouin zone; 11 × 11 × 11 for bulk, 3 × 3 × 1 for the slabs, and only at the Γ -point for molecules. Typical input files used for the calculations are attached to the ESI.† The slab models were illustrated using Visualization for Electronic and Structural Analysis (VESTA), a 3D visualisation program for structural models, volumetric data and crystal morphologies.¹⁰²

2.2 *Ab initio* thermodynamics

The stability of all surfaces considering the adsorption at given conditions was calculated using *ab initio* thermodynamics.¹⁰³ We used the same approach as described in our last work (Müller *et al.*⁹³). The configurational entropy can safely be neglected since its contribution to the Gibbs free energy is less than 3 meV Å^{–1} for transition metal surfaces and a surface area per surface site of around 10 Å² as long as the temperature remains significantly lower than 1000 K.¹⁰⁴ For a more detailed discussion than in the main text, we refer the reader to the ESI.†

2.2.1 Chemical potential. The chemical potential of a gas atom at a temperature of 0 K and standard pressure, μ_i^0 at 0 K, is the electronic energy of the substance as calculated by DFT-based calculations. The temperature- and pressure-dependency of the chemical potential of an atom in the gas phase can be calculated using the thermochemical tables¹⁰⁵ and the ideal gas law:

$$\mu_{Z,\text{gas}} = \mu_{Z,\text{gas}}^0 + \Delta G_Z(T) + k_B T \ln \left\{ \frac{p_Z}{p^0} \right\} \quad (5)$$



For gases consisting of multiple different atoms, such as CO₂, the chemical potential of a single atom is defined as the difference of the chemical potential of the whole molecule and the chemical potential of the remaining molecule without that specific atom. For example, the chemical potential of O* in CO₂ is the difference between the chemical potential of CO₂ and the chemical potential of CO.

$$\begin{aligned}\mu_{\text{O,CO}_2} &= \mu_{\text{CO}_2} - \mu_{\text{CO}} \\ &= \left[\mu_{\text{CO}_2}^0 - \mu_{\text{CO}}^0 \right] + [\Delta G_{\text{CO}_2}(T) - \Delta G_{\text{CO}}(T)] + k_{\text{B}} T \ln \left\{ \frac{p_{\text{CO}_2}}{p_{\text{CO}}} \right\}\end{aligned}\quad (6)$$

Thus, the chemical potential for any atom/molecule Z directly in the gas phase or a in gas containing Z can be represented by an equation of the following form:

$$\mu_{\text{Z,gas}} = a_{\text{Z}} + b_{\text{Z}}(T) + c_{\text{Z}}(T,p) \quad (7)$$

where a_{Z} contains all the constant terms, $b_{\text{Z}}(T)$ contains all temperature-dependent but pressure-independent terms and $c_{\text{Z}}(T,p)$ contains all both temperature- and pressure-depending terms. For a fixed temperature T_{fix} , the expression in eqn (7) can be further simplified into the following equation depending linearly on the logarithm of the individual partial pressures of each gas:

$$\mu_{\text{Z,gas}} = d_{\text{Z}}(T_{\text{fix}}) + e_{\text{Z}}(T_{\text{fix}}) \cdot \left[\ln \left\{ \frac{p_{\text{Z}}}{p^0} \right\} \text{ or } \ln \left\{ \frac{p_{\text{M-Z}}}{p_{\text{M}}} \right\} \right] \quad (8)$$

where the chemical potential of an atom/molecule Z in the gas phase for a fixed temperature T_{fix} depends only on the partial pressure of the atoms/molecules in the gas phase.

2.2.2 Chemical potentials under reaction conditions. The conditions to consider for gas-phase mixtures are those corresponding to the standard experimental particle preparation (pure H₂ or similar reductive conditions), both CO₂ and CO hydrogenation testing conditions used in our group¹⁰⁶ and industrially useful CO₂ hydrogenation conditions⁹ as described in our last article and others.^{93,107} Assuming a small conversion to methanol (<10%), we showed that by using the simplified equation for the WGS equilibrium constant as a function of the temperature provided by Moe¹⁰⁸ ($K_{\text{WGS}}(T)$) and the RWGS conversion (ξ_{RWGS}) as single parameter for any fixed temperature, the partial pressures of H₂, CO, CO₂ and H₂O under reaction conditions can be calculated based on the initial partial pressures of the different gases ($p_{\text{CO}_2}^0$, p_{CO}^0 , $p_{\text{H}_2}^0$, $p_{\text{H}_2\text{O}}^0$). Assuming that only CO₂ and H₂O act as oxidising agents, eqn (9) can be used to calculate the total chemical potential of oxygen under reaction conditions ($\mu_{\text{O,tot}}$):

$$\mu_{\text{O,tot}} = \max(\mu_{\text{O,CO}_2,\text{CO}}, \mu_{\text{O,H}_2\text{O,H}_2}) - \mu_{\text{O,O}_2} \quad (9)$$

where the chemical potentials can be calculated using eqn (8) and the individual partial pressures based on the RWGS conversion (ξ_{RWGS}) can be calculated using the following equation:

$$p_{\text{Y}}(\xi_{\text{RWGS}}) = p_{\text{Y,init}} - (p_{\text{Y,init}} - p_{\text{Y}}^{\text{RWGS}}(T)) \cdot \xi_{\text{RWGS}} \quad (10)$$

Since $\mu_{\text{O,tot}}$ only depends on the ratio of the partial pressures of CO₂ to CO and H₂O to H₂ and μ_{CO_2} , μ_{CO} , $\mu_{\text{H}_2\text{O}}$ and μ_{H_2} depend linearly on the logarithm of the individual partial pressures (p_{CO_2} , p_{CO} , p_{H_2} , $p_{\text{H}_2\text{O}}$; see eqn (8)), rather small changes are expected by changing the absolute pressures without changing the relative pressures. By doubling the pressures from the initial conditions ($T = 230$ °C, 15 bar H₂, 5 bar CO₂ to $T = 230$ °C, 30 bar H₂, 10 bar CO₂) we slightly increase the initial μ_{H_2} and μ_{CO_2} and thus the initial μ_{O} . Considering at least minimal RWGS conversion, μ_{H_2} and μ_{CO_2} decrease while $\mu_{\text{H}_2\text{O}}$ and μ_{CO} increase resulting in practically the same observed chemical potentials as observed for the initial conditions ($T = 230$ °C, 15 bar H₂, 5 bar CO₂). The same trend is observed if the partial pressure is doubled again resulting in the industrially relevant conditions ($T = 230$ °C, 60 bar H₂, 20 bar CO₂, $p^{\text{tot}} = 80$ bar). Since μ_{O} depends only on the ratio of CO₂ to CO and H₂O to H₂, considering minimal RWGS conversion results in equal values for μ_{O} for the three different partial pressures since the ratios of p_{CO_2} to p_{CO} and $p_{\text{H}_2\text{O}}$ to p_{H_2} are equal in all three cases (compare eqn (8)). A selection of conditions with explicit values for the different chemical potentials are shown in Table S1 in the ESI,[†] while the changes for the oxygen chemical potentials for selected starting conditions are shown in the ESI (Fig. S1–S4[†]).

More significant changes are observed if the composition of the feedstock changes. If we replace some of the CO₂ in the feedstock by CO, μ_{O} decreases. Fig. 1 shows μ_{O} under reaction conditions depending on the molar ratio of CO₂ in the feedstock (X_{CO_2}) with the colour gradient indicating the progressing RWGS conversion: from low (yellow) to high (red). Unless a pure CO-feed is present, some CO₂ and H₂ are always converted to CO and H₂O by the RWGS-reaction mainly due to the absence of H₂O in the feedstock. Since H₂O is a stronger oxidant than CO₂ unless no CO is present in the feedstock, μ_{O} increases with increasing RWGS conversion.

If some of the CO₂ in the feed is replaced by CO (50%), the initial μ_{CO} increases quite significantly. Since some CO₂ is always converted into CO by the RWGS-reaction, μ_{CO} increases further due to the progressing RWGS, but this effect is smaller than the effect of changing the initial gas phase composition. For the mixed feed, the μ_{H_2} under reaction conditions increases slightly compared to the pure CO₂ feed due to the lower consumption of H₂ by the RWGS-reaction since large quantities of CO are already present in the feed gas. Moreover, less H₂O is produced by the RWGS-reaction. Together with the lower CO₂ partial pressure, this decreases μ_{O} from the range of -3.11 eV to -2.92 eV (corresponding to an equivalent O₂ partial pressure below 10^{-48} bar at 500 K) to the range of -3.34 eV to -3.01 eV (corresponding to an equivalent O₂ partial pressure below 10^{-50} bar). This results are again practically independent of the different initial pressures if minimal RWGS conversion is assumed ($\xi_{\text{RWGS}} > 0.01$). Thus, the reaction mixture becomes slightly less oxidising and as a result, less oxygen is expected to adsorb on the surface.

If the CO₂ concentration further decreases to around 5%, the same trend continues for the active atmosphere under reaction conditions: μ_{H_2} increases very slightly, taking practically the same value as for the case of 50% CO, μ_{CO} increases quite



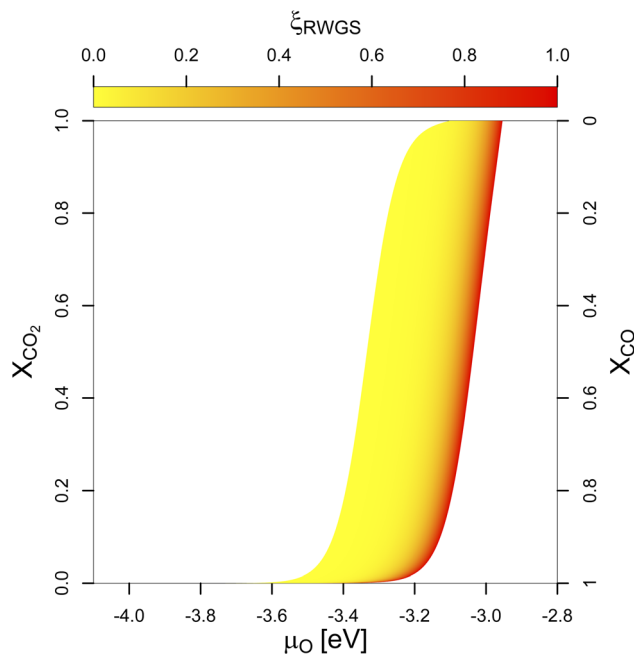


Fig. 1 Chemical potential of oxygen (μ_{O}) under CO_x hydrogenation conditions depending on the molar ratio of CO_2 in the feedstock (X_{CO_2}). The colour gradient from yellow to red indicates the progressing RWGS conversion (yellow indicates low conversion, red indicates higher RWGS conversion, respectively). μ_{O} does not depend linearly on the (R)WGS-conversion (compare Fig. S5 in the ESI†). For intermediate conditions ($0.1 > X_{\text{CO}_2} > 0.9$), μ_{O} under reaction conditions depends linearly on X_{CO_2} . For a pure CO_2 -feed ($X_{\text{CO}_2} = 1$), μ_{O} is much narrower and lies between -3.11 eV and -2.92 eV (corresponding to an equivalent O_2 partial pressure below 10^{-48} bar at 500 K). For a pure CO -feed ($X_{\text{CO}_2} = 0$), μ_{O} is much lower and lies between -4.07 to -3.93 eV (corresponding to an equivalent O_2 partial pressure below 10^{-69} bar). Removing the last quantities of CO_2 from the feedstock has the most drastic effect.

strongly, favouring the CO adsorption on the surface and μ_{O} decreases from the range of -3.34 eV to -3.01 eV to the range of -3.47 eV to -3.13 eV (corresponding to an equivalent O_2 partial pressure below 10^{-53} bar) due to a decrease of the CO_2 partial pressure and reduced water formation due to the RWGS-reaction. Overall, this further decreases oxygen adsorption.

The biggest change in terms of oxygen adsorption is observed if the last 5% of CO_2 is removed from the reaction conditions. In this case, the effect on μ_{H} and μ_{CO} are minimal. Nevertheless, due to the absence of both CO_2 and H_2O in the active atmosphere, the gas mixture becomes significantly less oxidising since no oxidant is present, with μ_{O} decreasing to a range of around -4.07 eV to -3.93 eV (corresponding to an equivalent O_2 partial pressure below 10^{-69} bar). Moreover, under these conditions, the WGS-reaction occurs rather than RWGS one.

The effect of the composition of the feedstock on μ_{O} shows that small quantities of CO_2 significantly increase μ_{O} under CO hydrogenation conditions. These results could explain why small quantities of CO_2 are needed in the syngas-feedstock for efficient CO hydrogenation: CO_2 can oxidise the promoter (Ga, Zn) under reaction conditions (*vide infra*), either directly as CO_2 or, more likely, indirectly in the form of water produced by the RWGS-

reaction, forming partial/full promoter surface oxides, with potential active sites for the CO_2 hydrogenation reaction at the interface of the metallic phase with the promoter (sub)oxides.

2.3 Mixing energy

The mixing energy of a bulk structure is the energy change upon mixing two elements relative to the individual bulk chemical potentials:

$$\Delta E_{n_X, n_Y}^{\text{mixing}} = E_{X_n Y_k, \text{bulk}} - n_X \cdot \mu_{X, \text{bulk}}^0 - n_Y \cdot \mu_{Y, \text{bulk}}^0 \quad (11)$$

2.4 Surface energy

We used the same approach described in our last work (Müller *et al.*⁹³) to calculate the surface energies of the evaluated systems (note that we calculate the Gibbs free surface energies). For a more detailed discussion, we refer the reader to the ESI.† The surface energy γ is defined as the energy needed to generate a surface from an equivalent bulk structure, normalized by the surface area of the resulting slab. It can be calculated for mono-metallic and multi-metallic systems:

$$\gamma_{n_X, n_Y} = \frac{E_{\text{tot, surface}} - E_{X_n Y_k, \text{bulk}}}{A} \quad (12)$$

If an atom Z is adsorbed from the gas phase (the source of adsorbed Z), the chemical potential of the atom in the gas phase is subtracted to calculate the surface energy in the presence of the adsorbate Z .

$$\gamma_{n_X, n_Y, n_Z} = \frac{E_{\text{surface}, n_X, n_Y, n_Z} - E_{X_n Y_k, \text{bulk}} - n_Z \cdot \mu_{Z, \text{gas}}}{A} \quad (13)$$

Inserting the expression in eqn (8) into (13) results in an expression depending linearly on the logarithm of the different partial pressures, obtaining the following equation:

$$\gamma_{n_X, n_Y, n_Z}(T_{\text{fix}}, p) = \frac{1}{A} \cdot f_{X, Y, Z}(T_{\text{fix}}) - \frac{1}{A} \cdot g_{X, Y, Z}(T_{\text{fix}}) \cdot \ln \left[\left\{ \frac{p_Z}{p^0} \right\} \text{ or } \ln \left\{ \frac{p_{M-Z}}{p_M} \right\} \right] \quad (14)$$

where the surface energy of a model $\gamma_{n_X, n_Y, n_Z}(T, p)$ for a fixed temperature T_{fix} depends only on the partial pressure of the adsorbate atoms/molecules in the gas phase. Under the investigated CO_x hydrogenation conditions, this approach works well for H^* and CO^* , while for O^* , eqn (9) can be used in combination with eqn (14).

2.5 Cu-alloy phase diagrams

The phase diagrams for the Cu-Ga- and Cu-Zn-systems reveal some of their key properties. Both are fully miscible and form alloys (intermetallics and solid solutions) with different structures depending on the composition. The phase diagrams of the two bimetallic mixtures are shown in Fig. 2 with the different crystal structures explained in Table 1.

For low promoter (Ga/Zn) concentrations, both structures crystallise in a solid solution in which the fcc-structure of pure



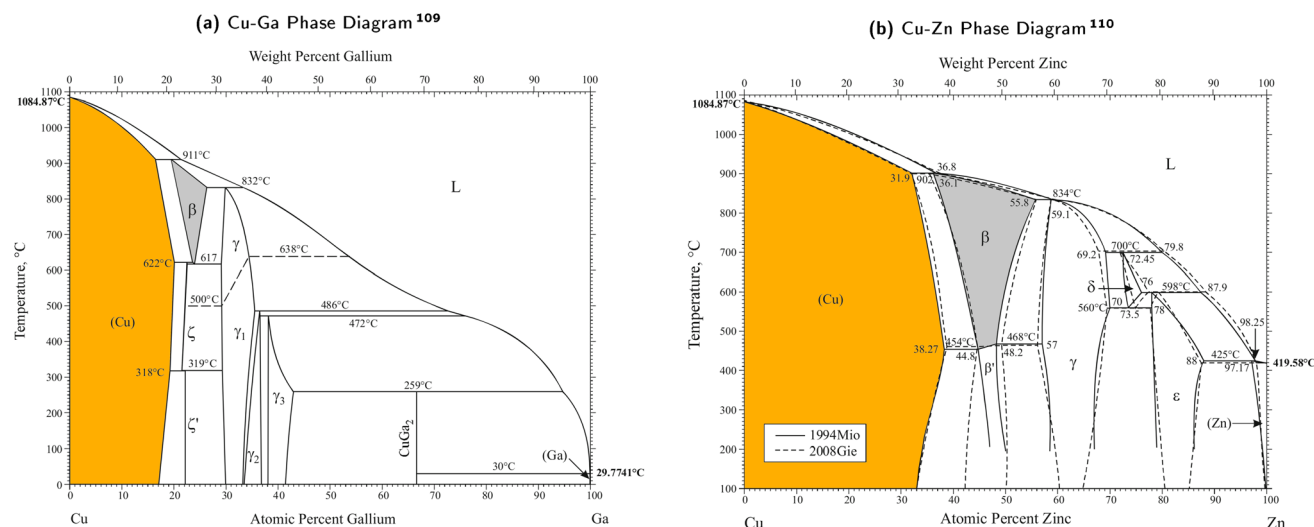


Fig. 2 Phase diagrams of the two bimetallic mixtures for the Cu–Ga- and Cu–Zn-system. Figures taken and adapted from Okamoto¹⁰⁹ (Cu–Ga phase diagram, (a)) and Okamoto¹¹⁰ (Cu–Zn phase diagram, (b)). For low concentrations of Ga/Zn, both structures crystallise in a solid solution with the fcc-structure of pure Cu modified by exchanging Cu with Ga/Zn (α -CuGa/-CuZn, coloured in orange). At higher concentrations, different crystal structures are obtained for the two systems (see Table 1).

Table 1 Crystal structures of the two bimetallic systems: Cu–Ga (Fig. 2a) and Cu–Zn (Fig. 2b)

Cu–Ga		Cu–Zn	
Phase	Crystal structure	Phase	Crystal structure
(Cu)	fcc-Cu (with Ga)	(Cu)	fcc-Cu (with Zn)
β	bcc-Cu (with Ga)	β	bcc-Cu (with Zn, disordered)
ζ	hcp-Cu (with Ga)	β'	bcc-Cu (with Zn, ordered)
ζ'	$\text{Cu}_{0.778}\text{Ga}_{0.222}$	γ	Cu_5Zn_8
$\gamma, \gamma_1, \gamma_2, \gamma_3$	Cu_9Ga_4 (multiple variations)	δ	bcc-Zn (with Cu)
CuGa_2	CuGa_2	ε	hcp-Zn (with Cu)
(Ga)	Ga-structures	(Zn)	hcp-Zn

Cu is modified by exchanging Cu with Ga/Zn. The Cu–Ga-system forms this stable fcc-solid solution up to a molar fraction of around 0.15 to 0.2 atomic% of Ga (α -CuGa, orange area in Fig. 2a), while the fcc-solid solution is much more stable for the Cu–Zn-systems up to a molar fraction of around 0.35 to 0.4 Zn (α -CuZn, orange area in Fig. 2b). For higher promoter concentrations, the effect of Ga/Zn becomes increasingly dominant, resulting in different crystal structures such as the bcc-solid solution (grey areas in Fig. 2a and b), the hcp-solid solutions, the stoichiometric ζ' -phase ($\text{Cu}_{0.778}\text{Ga}_{0.222}$) or even more complex systems at higher promoter concentrations (such as the γ -phase or its variations). In this article, we focus on evaluating the composition ranges within the α -phases of both systems (α -CuGa, orange area in Fig. 2a and α -CuZn, orange area in Fig. 2b).

3 Results and discussion

3.1 Bulk structures and alloy formation

We started by optimising the bulk structure of reference compounds, namely fcc-Cu, Cu_2O , CuO, hcp-Zn, w-ZnO

(Wurtzite ZnO), fcc- Cu_3Zn , bcc-CuZn, β -Ga, β - Ga_2O_3 , fcc- Cu_3Ga , bcc-CuGa, hcp-CuGa and Cu_9Ga_4 as well as relevant gas phase molecules (CO_2 , CO, H_2O , H_2 , O_2). Then, the oxide formation energies of Cu_2O , CuO, w-ZnO and β - Ga_2O_3 and the alloy formation energies for the fcc- Cu_3Zn , bcc-CuZn, fcc- Cu_3Ga , bcc-CuGa, hcp-CuGa and Cu_9Ga_4 were calculated with respect to their constituents and compared with published computational and experimental results.

The calculated formation enthalpies for copper(i) oxide (Cu_2O) and copper(ii) oxide (CuO) are $\Delta_f H(\text{Cu}_2\text{O})^{\text{calc}} = -1.25$ eV (-121 kJ mol⁻¹) and $\Delta_f H(\text{CuO})^{\text{calc}} = -1.12$ eV (-108 kJ mol⁻¹), which are in acceptable agreement with the experimental data of $\Delta_f H(\text{Cu}_2\text{O})^{\text{exp}} = -1.76$ eV (-170 kJ mol⁻¹) and $\Delta_f H(\text{CuO})^{\text{exp}} = -1.62$ eV (-156 kJ mol⁻¹).¹¹¹ The calculated formation enthalpies for wurtzite-like zinc oxide (w-ZnO) and β - Ga_2O_3 are $\Delta_f H(\text{w-ZnO})^{\text{calc}} = -2.91$ eV (-280 kJ mol⁻¹) and $\Delta_f H(\beta\text{-Ga}_2\text{O}_3)^{\text{calc}} = -9.41$ eV (-908 kJ mol⁻¹), which are again in acceptable agreement with experimental data of $\Delta_f H(\text{w-ZnO})^{\text{exp}} = -3.63$ eV (-351 kJ mol⁻¹) and $\Delta_f H(\beta\text{-Ga}_2\text{O}_3)^{\text{exp}} = -11.29$ eV (-1089 kJ mol⁻¹).¹¹⁴



The differences between calculated and experimental formation energies (0.51 eV for Cu₂O, 0.50 eV for CuO, 0.72 eV for w-ZnO and 1.88 eV for β-Ga₂O₃) are most likely due to the wrong description of the energy of the O₂ molecule and metal oxides when using the GGA-based PBE exchange-correlation functional. As a consequence, the formation energies of transition metal oxides are underestimated by a mean absolute error of *ca.* 0.83 eV per oxygen atom, while the error in the formation energies of CuO, Cu₂O and w-ZnO are 0.49, 0.53 and 0.77 eV per oxygen atom respectively in comparison to the reported values.¹¹⁵ The reported difference for the bulk formation energies per oxygen atom are almost identical to our results. Thus, since O₂ is poorly described using the PBE exchange-correlation functional, adding 1 eV to the calculated energy of the O₂ molecule decreases the difference between calculated and experimental formation energies to 0.01 eV for Cu₂O ($\Delta_f H(\text{Cu}_2\text{O})^{\text{calc,cor}} = -1.75$ eV), 0.00 eV for CuO ($\Delta_f H(\text{CuO})^{\text{calc,cor}} = -1.62$ eV), 0.22 eV for w-ZnO ($\Delta_f H(\text{w-ZnO})^{\text{calc,cor}} = -3.41$ eV) and 0.38 eV for β-Ga₂O₃ ($\Delta_f H(\beta\text{-Ga}_2\text{O}_3)^{\text{calc,cor}} = -10.91$ eV); the error in the formation energies is reduced without over-stabilising the bulk oxides. For further analysis we will thus use the corrected energy of O₂.

The calculated alloy formation enthalpies for fcc-Cu₃Ga and fcc-Cu₃Zn are $\Delta_f H(\text{fcc-Cu}_3\text{Ga})^{\text{calc}} = -0.39$ eV (-37 kJ mol⁻¹) and $\Delta_f H(\text{fcc-Cu}_3\text{Zn})^{\text{calc}} = -0.28$ eV (-27 kJ mol⁻¹). The formation of the fcc-Cu₃Ga alloy is thus more favoured in comparison to the fcc-Cu₃Zn one. The calculated alloy formation enthalpies for bcc-CuGa and bcc-CuZn are $\Delta_f H(\text{bcc-CuGa})^{\text{calc}} = 0.014$ eV (1.4 kJ mol⁻¹) and $\Delta_f H(\text{bcc-CuZn})^{\text{calc}} = -0.19$ eV (-18 kJ mol⁻¹). Only the formation energy of bcc-CuZn is negative, as expected from the alloy phase diagrams of the Cu-Ga- and the Cu-Zn-system (Fig. 2a and b).^{82,83} This indicates that bcc-CuZn is stable at low temperatures, while bcc-CuGa is only stable at high temperatures. The calculated alloy formation enthalpy for hcp-CuGa is $\Delta_f H(\text{hcp-CuGa})^{\text{calc}} = -0.066$ eV (-6.4 kJ mol⁻¹), thus being more stable than the bcc-CuGa structure (as expected from the phase diagram in Fig. 2a). The calculated alloy formation enthalpy for Cu₉Ga₄ is $\Delta_f H(\text{Cu}_9\text{Ga}_4)^{\text{calc}} = -1.58$ eV (-152 kJ mol⁻¹), which makes Cu₉Ga₄ a remarkably stable alloy. The good agreement of the calculated alloy formation energies with the experimental phase diagrams gives confidence in the chosen computational methodology. The calculated structural parameters are also in good agreement with the experimental ones (see Table 2).

3.1.1 Bulk mixing energy. The mixing energies for the fcc-Cu-Ga- and fcc-Cu-Zn-systems are calculated according to eqn (11) and are shown in Fig. 3. The mixing energy is strictly monotonic decreasing when adding additional Ga/Zn to the fcc-lattice for molar fractions up to 0.25 (in a somewhat linear fashion). The relative stabilisation is stronger for the Cu-Ga-system than for the Cu-Zn-system (Fig. 3a *vs.* Fig. 3b)). The Cu-Ga-system shows a total mixing energy lower (more negative) than -3.0 eV for 25% Ga, while the Cu-Zn-system shows a total mixing energy slightly lower than -2.0 eV for 25% Zn. Nevertheless, when the molar fraction increases to values above 0.25, different energy trends are observed for the Cu-Zn and Cu-Ga-system; this change is associated with the direct interaction between two promoter atoms of the same kind (Ga-Ga-*vs.* Zn-Zn-contacts). Hence, the Cu-Ga-system becomes less stable (less negative) with the mixing energy increasing linearly with the molar fraction. In contrast, the mixing energy for the Cu-Zn-system further decreases (becomes more negative) to a minima around a molar fraction of 0.5, but the decrease is less pronounced for Zn molar fractions in between 0.25 and 0.5 than for Zn molar fractions lower than 0.25. Surprisingly, the alloying energy at the minima (Zn molar ratio of 0.5) is equal to $\Delta_f H(\text{fcc-Cu}_{16}\text{Zn}_{16})^{\text{calc}} = -2.70$ eV (-260 kJ mol⁻¹), which is still smaller than the alloying energy for a Ga molar ratio of 0.25 ($\Delta_f H(\text{fcc-Cu}_{24}\text{Ga}_8)^{\text{calc}} = -3.13$ eV (-302 kJ mol⁻¹)).

In short, we show that mixing fcc-Cu with either Zn or Ga by exchanging small amounts of Cu with Zn or Ga stabilises the system by decreasing its total energy when compared to the individual bulk structures. The interaction of Cu with Zn or Ga is thus significantly stronger than the Zn-Zn- and Ga-Ga-interactions. Moreover, the differences in the mixing energy for molar fractions above 0.25 clearly show that Ga-Ga-interactions are less favourable than the Zn-Zn-interactions. These trends and the associated thermodynamic implications can help rationalising the phase diagrams in Fig. 2: Ga-Ga-bonds are rather weak and destabilise the fcc-Cu₃Ga structure. Thus, other crystal structures with less bonds in general, but also less Ga-Ga-bonds such as the ζ or ζ' structure (see Fig. 2a) are observed rather than structures with a higher mean coordination for the metals. The discrepancy between the observed change of the crystal lattice and the expected change based on the mixing energy depending on the Ga molar fraction (Fig. 3a) is partially due to the Cu-Ga system being a solid solution instead of an intermetallic system. Accordingly, Ga is

Table 2 Table with all relevant structural parameters of the optimized bulk structures in this article. The good agreement with experimental data for the structural parameters gives further confidence about the chosen computational methodology

Structure	Calculated structural parameters	Reference structural parameters
fcc-Cu	$a = 3.630$ Å	$a = 3.597 \pm 0.004$ Å (ref. 111)
Cu ₂ O	$a = 4.272$ Å	$a = 4.268$ Å (ref. 111)
CuO	$a = 4.620$ Å, $b = 3.515$ Å, $c = 5.141$ Å, $\beta = 98.28^\circ$	$a = 4.684$ Å, $b = 3.423$ Å, $c = 5.129$ Å, $\beta = 99.54^\circ$ (ref. 111)
hcp-Zn	$a = b = 2.683$ Å, $c = 4.812$ Å	$a = b = 2.665$ Å, $c = 4.947$ Å (ref. 112)
α-Ga	$a = 4.597$ Å, $b = 7.761$ Å, $c = 4.593$ Å	$a = 4.520$ Å, $b = 7.663$ Å, $c = 4.526$ Å (ref. 113)
ZnO	$a = b = 3.256$ Å, $c = 5.237$ Å	$a = b = 3.250$ Å, $c = 5.207$ Å (ref. 114)
β-Ga ₂ O ₃	$a = 12.280$ Å, $b = 3.070$ Å, $c = 5.867$ Å, $\beta = 103.78^\circ$	$a = 12.225$ Å, $b = 3.040$ Å, $c = 5.809$ Å, $\beta = 103.8^\circ$ (ref. 114)



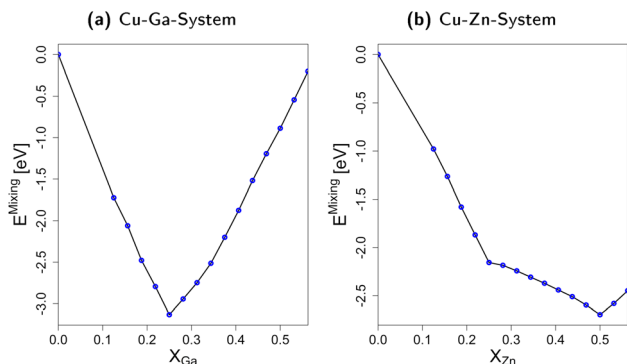


Fig. 3 Mixing energies for both the fcc Cu–Ga- (a) and fcc Cu–Zn-system (b) based on the molar fraction of Ga/Zn. The blue points indicated data points from calculations of explicit bulk structures. The difference in behaviour for a molar fraction above 0.25 Ga/Zn indicates the difference between Ga–Ga and Zn–Zn interactions.

statistically distributed in the fcc-Cu-lattice meaning direct Ga–Ga-bonds are already expected for lower Ga molar fractions than in the ideal, fully ordered crystal. In contrast, the Zn–Zn-bonds are significantly more stable, resulting in the α -brass structure being observed for Zn molar fractions of up to around 0.35–0.4 (see Fig. 2b), when a transition to the bcc-structure (β -Brass) is observed.

3.2 Slab models

We recently generated Cu slab models representing the different surfaces of Cu nanoparticles (NPs).⁹³ Here, we used an analogous approach to investigate Cu–Ga and Cu–Zn surface structures. The bulk structure of fcc-Cu₃Ga (ideal structure of α -Cu₃Ga) was cut at different crystal planes to generate specific surface facets, namely the low index (100)-, (110)- and (111)-facets. We ensured for each facet that the surface layer consists of four atoms. The slab model of all facets consists of 8 layers (4 atoms per layer, 32 atoms in total). Top-view images of the selected slabs and the specific adsorption sites evaluated are depicted in Fig. 4. A single colour is used for all atoms, since the atomic distribution changes for every facet.

Each of the facets has different adsorption sites. The (100)-facet has three distinct adsorption sites (Fig. 4a): the first one is a hollow four-fold site on top of the subsurface atomic site (red square). The second site is on the top of an atom of the top-most surface (red circle) and the third site is a bridge site between two atoms of the top-most layer (red rectangle).

The (110)-facet has four distinct adsorption sites (Fig. 4b): the first one is a four-fold site on the hollow site of the top layer of the surface (blue square), the second site is on top of an atom of the top-most layer (blue circle), the third one is on the short bridging site between two atoms (blue rectangle), while the fourth site is on the long bridge site between two surface atoms of the top-most layer (blue rectangle with black point).

The (111)-facet has four distinct adsorption sites (Fig. 4c): the first one is a three-fold hollow fcc-site (green triangle), the second site is a three-fold hollow hcp-site (green triangle with black point), the third site is on top of an atom of the top-most

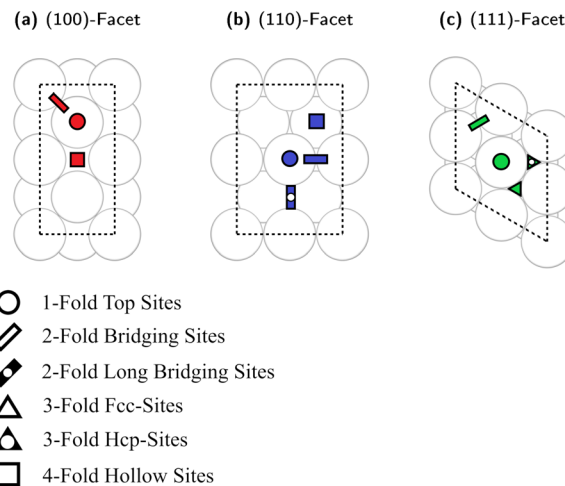


Fig. 4 Top view onto the slab models used for surface energy calculations with one colour for all atoms of the models. The adsorption sites are indicated by specific colours (one colour for each facet) and symbols, which are explained in the legend on the bottom. (a) (100)-facet. (b) (110)-facet. (c) (111)-facet.

layer (green circle), and the fourth site is a bridge site between two atoms of the top-most layer (green rectangle).

For all investigated facets, there are multiple variations of each distinct adsorption sites for different promoter concentrations and spatial distributions. For instance, if one atom in the surface layer of the (111)-facet is exchanged for Ga, there are now two different variations for each initial adsorption site (either adjacent to Ga or not adjacent to Ga). Similar variations exist for all investigated facets with multiple different variations for each adsorption site for all promoter concentrations, which must also be considered when investigating the different slabs.

3.2.1 Promoter atom distribution. Different promoter atom concentrations were investigated for every evaluated facet for each alloy (12.5–43.75% Zn for the Cu–Zn-system, 12.5–31.25% for the Cu–Ga-system) and compared to the pure Cu model. For low promoter concentrations (up to 25%, (sub)stoichiometric fcc-alloys), the promoter atoms are evenly distributed within the Cu-fcc structure to have the best possible representation of the fcc-Cu₃Zn/Cu₃Ga structure, *i.e.* ensuring the best dilution of the promoters in the fcc-Cu-lattice. For higher promoter concentrations (more than 25%, superstoichiometric fcc-alloys), additional Cu was replaced by the promoters, again ensuring the best possible dilution of the promoters in the slab models by minimising promoters as nearest neighbours of each other.

Moreover, we also investigated slab models with different spatial distribution of the two metals inside the slab models for different promoter concentrations (*i.e.*, partial or total separation of the two metals in the slab). This includes partial enrichment of the promoter on the relaxed surface (*e.g.*, surface layers with 50% promoter concentrations instead of the only 25%), the diminishment of the promoters on the optimised surface (*e.g.*, surface layers without promoters instead of the 25% in the ideal structure), or the formation of one or multiple surface layers of promoters on top of the pure copper structure



or (sub)stoichiometric fcc-alloys, among others. This enabled us to better represent the possible slab models, including fully alloyed and dealloyed structures and intermediate situations.

We fixed the four bottom layers as either stoichiometric fcc-alloys (Cu_3Ga , Cu_3Zn , Cu- or Ga-/Zn-terminated) or pure fcc-Cu layers to correctly subtract the surface energy of the bottom layers. A selection of some of the used slab models for the Cu–Zn system is shown in Fig. 5. The Cu–Ga-system slab models are analogous to the Cu–Zn- ones with Zn exchanged for Ga and the

lattice parameters modified according to the values obtained from bulk fcc- Cu_3Ga , ensuring that the promoter concentration is not too high (lower or equal to 31.25% Ga).

3.2.2 Surface-adsorption of gas-phase adsorbates. The adsorption of atomic oxygen (O^*), atomic hydrogen (H^*), carbon monoxide (CO^*), and carbon dioxide (CO_2^*) was investigated for all Cu–Ga and Cu–Zn surfaces by placing the adsorbate on a selection of the previously described surface sites while ensuring that the most distinct variations due to

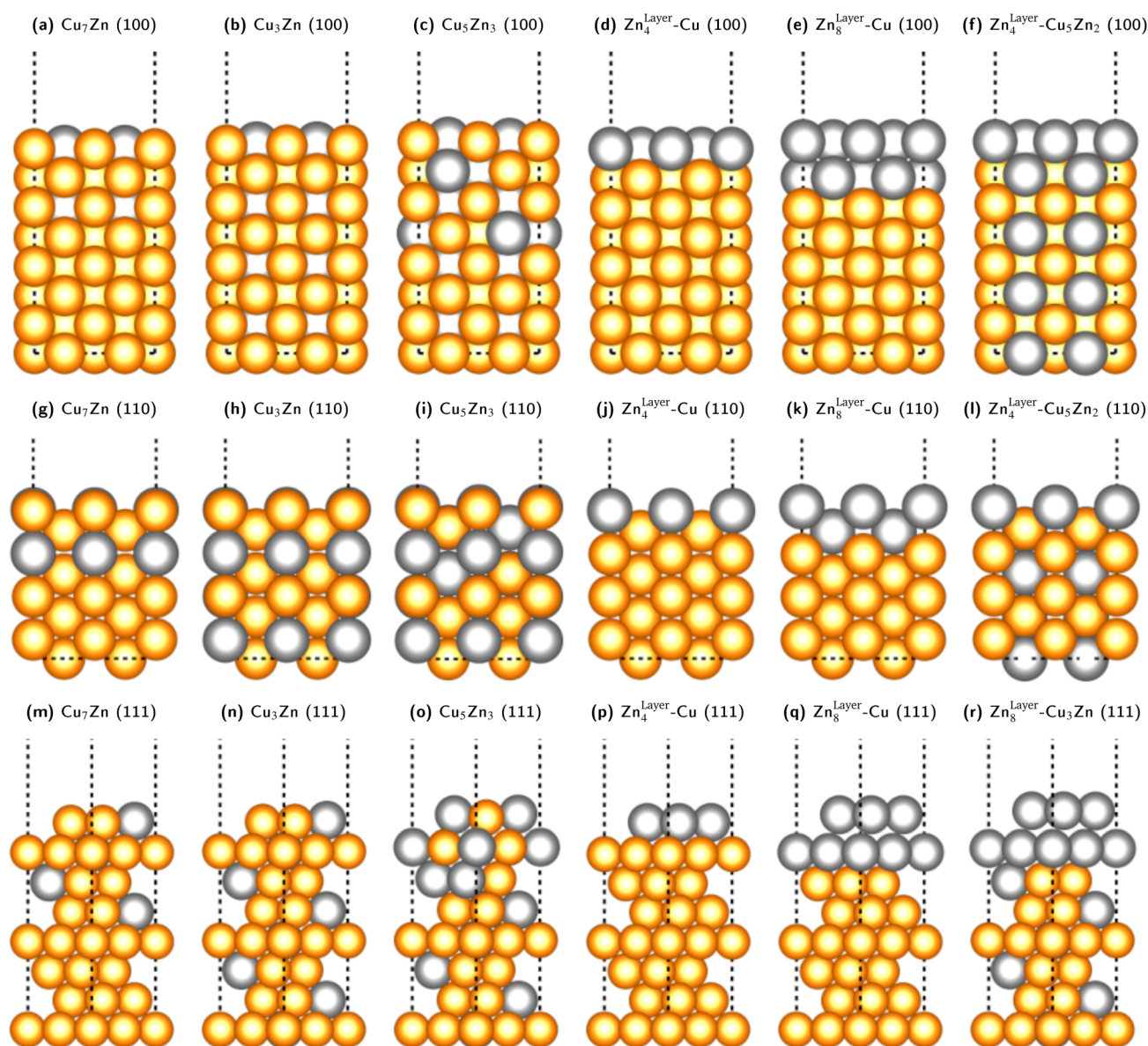


Fig. 5 Selection of some of the used slab models for the Cu–Zn system with different promoter concentrations and distributions in the slab model for the (100)-, (110)- and (111)-facet. The same models with Zn exchanged for Ga were used for the Cu–Ga-system, but structures with Ga-content higher than 31% were not investigated since the bulk structure would not be fcc based on the mixing energy shown in Fig. 3. The depicted models are: (100)-facet: (a) substoichiometric fcc- Cu_3Zn (12.5% Zn). (b) Stoichiometric fcc- Cu_3Zn (25% Zn). (c) Superstoichiometric fcc- Cu_3Zn (37.5% Zn). (d) One Zn-layer on pure fcc-Cu. (e) Two Zn-layers on pure fcc-Cu. (f) One Zn-layer on stoichiometric fcc- Cu_3Zn . (110)-facet: (g) substoichiometric fcc- Cu_3Zn (12.5% Zn). (h) Stoichiometric fcc- Cu_3Zn (25% Zn). (i) Superstoichiometric fcc- Cu_3Zn (37.5% Zn). (j) One Zn-layer on pure fcc-Cu. (k) Two Zn-layers on pure fcc-Cu. (l) One Zn-layer on stoichiometric fcc- Cu_3Zn . (111)-facet: (m) substoichiometric fcc- Cu_3Zn (12.5% Zn). (n) Stoichiometric fcc- Cu_3Zn (25% Zn). (o) Superstoichiometric fcc- Cu_3Zn (37.5% Zn). (p) One Zn-layer on pure fcc-Cu. (q) Two Zn-layers on pure fcc-Cu. (r) Two Zn-layers on stoichiometric fcc- Cu_3Zn .



different promoter concentrations and spatial distributions were studied, followed by an optimisation. While studying all variations of all specific surface sites is not possible due to the high dimensionality of the accessible parameter space, we have ensured a good representation by selecting many structurally different surface sites, including different configurations of the same adsorption site. The adsorption of multiples of the same adsorbate to investigate different coverages was addressed analogously by testing multiple distinct structures. The total number of adsorbates for each structure was lower or equal to four H*/CO* adsorbed, which corresponds to a H*/CO* surface coverage between 0.25 and 1 monolayer (ML), lower or equal to one adsorbed CO₂* (corresponding to a CO₂* surface coverage of 0.25 ML) and lower or equal to six O* adsorbed, corresponding to an O* surface coverage between 0.25 and 1.5 ML. Since we assumed that oxygen could reconstruct the surface quite extensively, especially for high coverages, placing O* just onto the surfaces might not ensure the formation of these oxide-like overlayers. Thus, we also looked for subsurface variations of the previously explained O* adsorption sites to ensure extensive surface reconstruction: we placed oxygen atoms into different tetrahedral and octahedral sites of the subsurface layers of the alloy slab. We selected a large number of structurally distinct surface and subsurface sites for this purpose. The adsorption of

the different adsorbates onto the same facets with different promoters M (M = Ga, Zn) will be discussed first.

3.3 Surface stability diagrams – how chemical potential (gas phase composition) affects surface states

One-dimensional surface stability diagrams (SSDs) at 500 K (CO_x hydrogenation temperature) were constructed for each facet and promoter (Ga/Zn) atom distribution in the slab models considering the adsorption of oxygen (O*), hydrogen (H*), carbon monoxide (CO*) and carbon dioxide (CO₂*) to evaluate the surface state of the different model systems in different atmospheres. The presence of O* is particularly relevant for the CO/CO₂ hydrogenations^{71,72,116} in particular for a CO₂ rich atmosphere as it is rather oxidizing for Ga and Zn (*vide infra*); we will thus start with discussing O* adsorption on the surfaces.

3.3.1 Adsorption of atomic oxygen – O*. Fig. 6 shows the most stable surfaces of the stoichiometric alloys depending on the oxygen chemical potential (μ_{O}) and equivalent oxygen partial pressure with corrected surface energies for both the Cu–Zn-systems (Fig. 6 left) and Cu–Ga-systems (Fig. 6 right). The yellow-red area indicates the expected μ_{O} under CO₂ hydrogenation conditions ($T = 230$ °C, $p_{\text{H}_2} = 15$ bar, $p_{\text{CO}_2} = 5$ bar, $p_{\text{Ar}} = 5$ bar, -3.11 eV $< \mu_{\text{O}} < -2.92$ eV, corresponding to an equivalent O₂ partial pressure below 10⁻⁴⁸ bar) and the green-yellow area

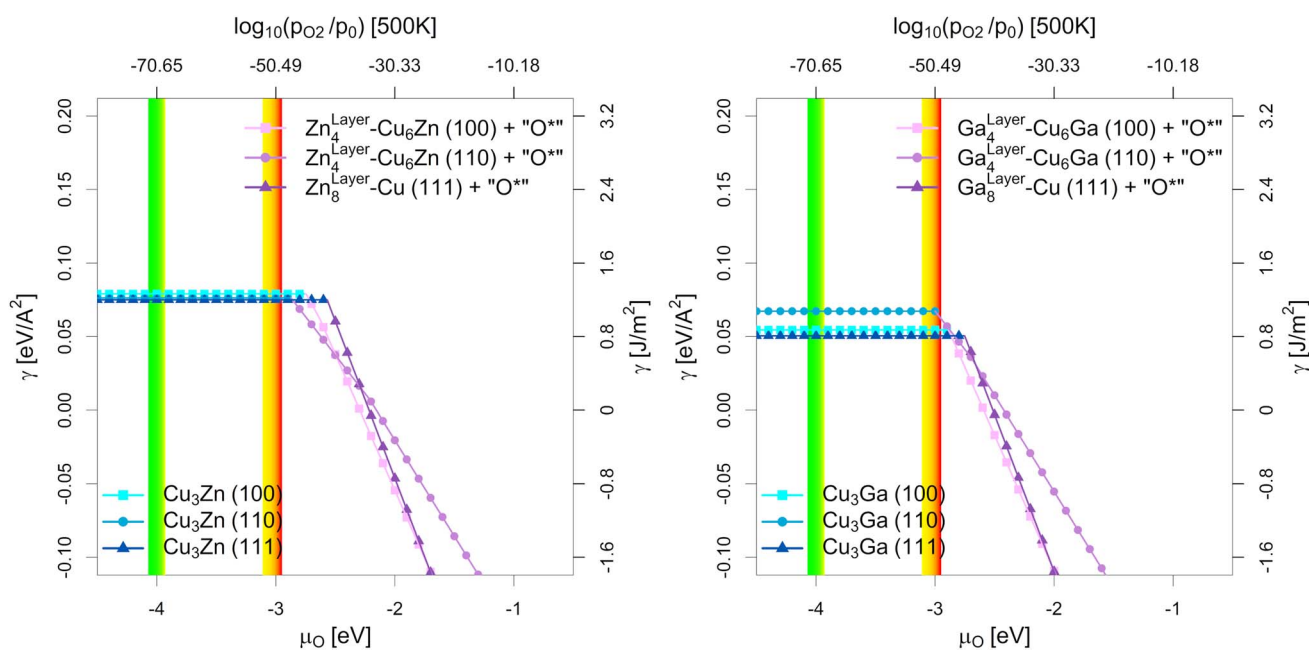


Fig. 6 Most stable surfaces of the stoichiometric fcc-Cu₃Ga/fcc-Cu₃Zn alloys (25% Ga/Zn) depending on the chemical potential of oxygen (μ_{O} in eV) and the equivalent oxygen partial pressure at 500 K (Cu–Zn-system on the left, Cu–Ga-system on the right). The investigated systems are the unlayered structure with stoichiometric fcc-Cu₃Ga/fcc-Cu₃Zn alloy (see Fig. 5b, h and n), the partially layered structure with one Ga/Zn surface layer on substoichiometric fcc-Cu₃Ga/fcc-Cu₃Zn alloy and the fully layered structure with two Ga/Zn surface layers on pure fcc-Cu (see Fig. 5e, k and q). The yellow-red area indicates the expected oxygen chemical potential (μ_{O}) under CO₂ hydrogenation conditions (-3.11 eV $< \mu_{\text{O}} < -2.92$ eV, corresponding to an equivalent O₂ partial pressure below 10⁻⁴⁸ bar) while the green-yellow area indicates the expected μ_{O} under CO hydrogenation conditions (-4.07 eV $< \mu_{\text{O}} < -3.93$ eV, equivalent to an O₂ partial pressure below 10⁻⁶⁹ bar) with the colour gradient indicating the progressing (R)WGS reaction (yellow for low conversion, red/green for high conversion; see Fig. 1 as well as Fig. S1, S4 and S5 in the ESI†). At low μ_{O} , the alloyed structures are more stable (blue part without incline), while the partially/fully layered structures are more stable at high μ_{O} due to surface oxidation (violet part with incline).



indicates the expected μ_{O} under CO hydrogenation conditions ($T = 230\text{ }^{\circ}\text{C}$, $p_{\text{H}_2} = 15\text{ bar}$, $p_{\text{CO}} = 5\text{ bar}$, $p_{\text{Ar}} = 5\text{ bar}$, $-4.07\text{ eV} < \mu_{\text{O}} < -3.93\text{ eV}$, corresponding to an equivalent O_2 partial pressure below 10^{-69} bar). The colour gradient indicates the progressing (R)WGS-reaction (yellow for low conversion, red/green for high conversion, compare Fig. 1). Note that the oxygen chemical potential does not depend linearly on the RWGS-conversion (compare Fig. S5 in the ESI†). The correction of the surface energy by correcting the energy of the O_2 molecule results in a shift of the μ_{O} where the oxidised surface is more stable than the reduced surface (called point of oxygen adsorption in the following text) by around 0.5 eV. Only the corrected surface energies are shown in Fig. 6; we refer the reader to the ESI† for the individual surface stability diagrams in oxygen atmosphere including the uncorrected surface energies.

As shown in Fig. 5, there are multiple configurations with different spatial distribution for the stoichiometric alloys (25% Ga/Zn): the fully alloyed structure with stoichiometric fcc- $\text{Cu}_3\text{Ga}/\text{fcc-Cu}_3\text{Zn}$ alloy (compare Fig. 5b, h and n), the partially layered structure with one Ga/Zn surface layer on substoichiometric fcc- $\text{Cu}_3\text{Ga}/\text{fcc-Cu}_3\text{Zn}$ alloy and the fully layered structure with two Ga/Zn surface layer on pure fcc-Cu (compare Fig. 5e, k and q). In vacuum (and thus under very low μ_{O} , such as CO hydrogenation conditions without CO_2), the alloyed, unlayered structures are the most stable configurations independently of the investigated facet and used promoter, since the overall alloying energy is negative. If the gas phase becomes more oxidising (μ_{O} increases), the partially and fully alloyed

structures show a significantly stronger stabilisation due to preferred interaction of the promoters with oxygen when compared to Cu. Thus, if μ_{O} increases enough (e.g. CO_2 hydrogenation conditions), dealloying and surface (sub)oxide formation is observed for both promoters for all investigated surface facets.

Comparing the points of oxygen adsorption, and thus dealloying and surface (sub)oxide formation (formation of $\text{GaO}_x/\text{ZnO}_x$) reveals some key features of the two promoters. The surface energies of the alloyed structures are rather similar in vacuum conditions, with the surface energy of the (111)-facet being the lowest while the one of the (110)-facet being the highest, mirroring the results of pure Cu.⁹³ Moreover, while both promoters stabilise oxygen adsorption, a stronger stabilisation of the dealloyed structure is observed for the Cu-Ga-system than for the Cu-Zn-system. The oxygen adsorption for the Cu-Ga-system is so strongly stabilised that full dealloying and surface oxidation becomes thermodynamically favoured under CO_2 hydrogenation conditions, especially for the (110)-facet, while the Cu-Zn-system should remain fully alloyed. Nevertheless, the (110)-facet is also the easiest facet to oxidise for the Cu-Zn-system, and the point of oxygen adsorption is rather close to the μ_{O} expected under CO_2 hydrogenation conditions.

To better understand the dependency of the oxygen adsorption and (sub)oxide formation on the promoter concentration in the fcc-Cu crystal lattice, we plotted the points of oxygen adsorption (lowest μ_{O} where the slab model with one or

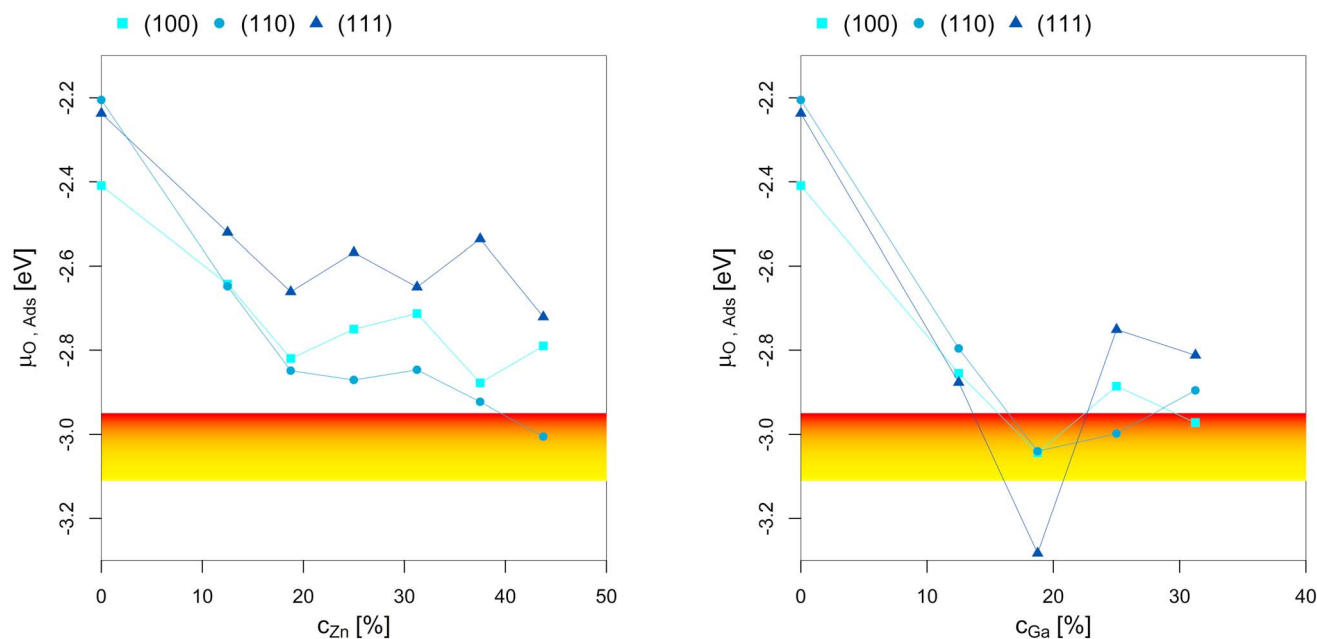


Fig. 7 Points of oxygen adsorption (lowest μ_{O} where the slab model with one or multiple adsorbed O^* is more stable than the clean surface), dealloying and surface oxidation depending on the Zn-/Ga-concentration in the fcc-Cu crystal lattice. The yellow-red area indicates the expected oxygen chemical potential (μ_{O}) under CO_2 hydrogenation conditions ($-3.11\text{ eV} < \mu_{\text{O}} < -2.92\text{ eV}$, corresponding to an equivalent O_2 partial pressure below 10^{-48} bar). The expected μ_{O} under CO hydrogenation conditions ($-4.07\text{ eV} < \mu_{\text{O}} < -3.93\text{ eV}$, equivalent to an O_2 partial pressure below 10^{-69} bar) is so low that it is not shown in the plots. The point of oxygen adsorption generally decreases towards lower μ_{O} for increasing promoter concentrations. For the Cu-Zn-system, a plateau is reached for intermediate promoter concentrations (around 25% Zn), with and additional decrease observed for the highly superstoichiometric alloys (around 40% Zn). For the Cu-Ga-system, a minimum is reached for around 20% Ga, with the point of oxygen adsorption increasing for higher promoter concentrations. Data on pure Cu taken from Müller *et al.*⁹³



multiple adsorbed O* is more stable than the clean surface) depending on the Zn-/Ga-concentration (Fig. 7). The yellow-red area indicates the expected μ_{O} under CO₂ hydrogenation conditions ($-3.11 \text{ eV} < \mu_{\text{O}} < -2.92 \text{ eV}$, corresponding to an equivalent O₂ partial pressure below 10^{-48} bar). The expected μ_{O} under CO hydrogenation conditions ($-4.07 \text{ eV} < \mu_{\text{O}} < -3.93 \text{ eV}$, equivalent to an O₂ partial pressure below 10^{-69} bar) is not shown in the plots due to its low value. Fig. 7 shows that the point of oxygen adsorption generally decreases towards lower μ_{O} for increasing promoter concentrations, especially when compared to pure Cu as reported in our last article.⁹³ The most stable surfaces of the alloys in an oxygen atmosphere with different stoichiometries and the individual surface stability diagrams in oxygen atmosphere are reported in the ESI.†

For the Cu-Zn-system at low promoter concentration ($0\% < c_{\text{Zn}} < 20\%$), the point of oxygen adsorption decreases linearly as a function of the promoter concentration. The surface oxidation is especially stabilised for the (110)-facet, while the two other facets show rather similar behaviour for the point of oxygen adsorption. For intermediate promoter concentrations ($20\% < c_{\text{Zn}} < 40\%$), a plateau is reached, with the point of oxygen adsorption being mostly independent of the promoter concentration. When the promoter concentration is further increased ($c_{\text{Zn}} > 40\%$), the point of oxygen adsorption decreases again, but only for the (110)-facet. The surface oxidation is so stabilized for this termination that partial dealloying is predicted under CO₂ hydrogenation reaction conditions, while fully reduced/alloyed surfaces are expected under CO hydrogenation reaction conditions.

For the Cu-Ga-system at low promoter concentrations ($0\% < c_{\text{Ga}} < 20\%$), the point of oxygen adsorption again decreases linearly as a function of the promoter concentration. However, the Cu-Ga-system shows several key differences compared to the Cu-Zn-system. First, the linear decrease has a much steeper slope compared to the Cu-Zn-system, indicating that small quantities of Ga have a much stronger effect on stabilising the surface oxidation when compared to Zn. Thus, dealloying and surface oxidation under CO₂ hydrogenation reaction conditions is expected for the Cu-Ga-system with a Ga-concentration of around 20% irrespective of the investigated facets, while the equivalent structures for the Cu-Zn-system remain alloyed and reduced under the same conditions. Second, while the trends for the (100)- and (110)-facet are equivalent for both the Cu-Zn- and Cu-Ga-systems, the (111)-facet shows such a strong stabilisation that fully oxidised surfaces are expected even at a μ_{O} lower than the one expected under CO₂ hydrogenation reaction conditions.

If the promoter concentration increases further, the stabilisation is reversed: oxygen adsorption and surface oxidation becomes less favoured when increasing the promoter concentration, showing that the Cu-Ga-system has an optimal promoter concentration at around 18.75%; a similar effect has not been observed for the Cu-Zn-system. This shows that the Cu-Ga-system is more easily dealloyed and oxidised under CO₂ hydrogenation conditions than the Cu-Zn-system; but while the oxygen adsorption on the Cu-Ga-system has a stronger dependency on the promoter (Ga) concentration, the Cu-Zn-system

seems less affected by higher promoter (Zn) concentrations, which is similar to what has been observed for the alloy formation energies (Fig. 3) and mirrors again the results of the experimental phase diagrams (Fig. 2).

3.3.2 Adsorption of hydrogen. Concerning the hydrogen adsorption, we plotted the points of hydrogen adsorption (lowest μ_{H} where the slab model with one or multiple adsorbed H* is more stable than the clean surface) depending on the Zn-/Ga-concentration in Fig. 8. We refer the reader to the ESI† for the individual plots of the most stable surfaces of the alloys with different stoichiometries in a hydrogen atmosphere and the individual surface stability diagrams. The red/green line indicates μ_{H} expected under CO₂ or CO hydrogenation conditions ($-0.235 \text{ eV} < \mu_{\text{H}} < -0.234 \text{ eV}$).

Hydrogen adsorption under CO₂ or CO hydrogenation conditions is exclusively located on the three-fold adsorption site on three Cu atoms on the (111)-facet, as shown for pure Cu.⁹³ As a result, adsorption of H* on the (111)-facet is observed for both the Cu-Ga- and Cu-Zn-systems as long as the three-fold adsorption site on three Cu atoms is present. This adsorption is only observed for promoter concentrations up to 25% and if the promoters are not enriched on the top-most layer. Less hydrogen adsorption is expected for higher promoter concentrations since the promoters disfavour the formation of surface hydrides. But this effect is rather weak when compared to the effect of the promoters on oxygen adsorption. As a result, hydrogen adsorption is still expected under CO_x hydrogenation reaction conditions, but the amount of H* decreases as a function of the promoter concentration. Thus, increasing the hydrogen partial pressure results in realloying, favouring the formation of additional three-fold adsorption sites composed by three Cu atoms. Nevertheless, this only affects the (111)-facet since the other facets are not modified in the presence of H₂.

In summary, while the promoters influence hydrogen adsorption, this effect is much weaker than the one observed for O*, since only one specific adsorption site of the Cu-based alloys adsorb H* under CO₂ or CO hydrogenation conditions. Thus, the promoters reduce the total amount of adsorbed H*, but a partial H*-coverage is still observed under CO_x hydrogenation conditions.

The feedstock has a much stronger effect on the surface state: under only CO and H₂, no surface oxidation and surface layer formation is observed, and only hydrogen is adsorbed. In contrast, if the feedstock contains large quantities of CO₂ (and H₂), surface oxidation and dealloying (by surface oxide layer formation) becomes possible for both promoters, and this effect is stronger for the Cu-Ga-alloys than for the Cu-Zn- ones; however, hydrogen adsorption and realloying is still possible under these conditions. Thus, both oxygen adsorption, surface oxidation and dealloying (by surface oxide layer formation) as well as hydrogen adsorption and realloying are thermodynamically feasible under CO₂ hydrogenation conditions, which might result in a highly dynamic system under reaction conditions due to the competing oxidation and reduction processes. Such highly dynamic processes have already been observed experimentally using *operando* TEM while studying the oxidation of hydrogen on a Cu catalyst.¹¹⁷



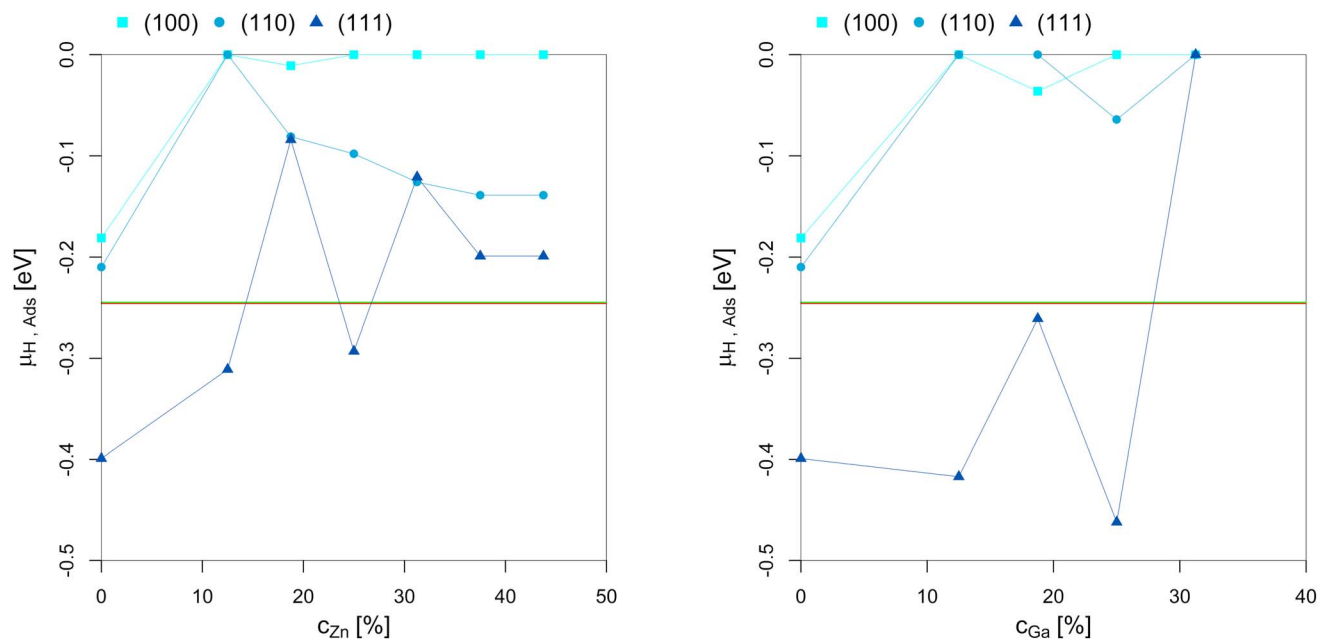


Fig. 8 Points of hydrogen adsorption (lowest μ_{H} where the slab model with one or multiple adsorbed H^* is more stable than the clean surface) depending on the Zn-/Ga-concentration in the fcc-Cu crystal lattice. The red/green line indicates μ_{H} expected under CO_2 or CO hydrogenation conditions ($-0.235 \text{ eV} < \mu_{\text{H}} < -0.234 \text{ eV}$). The overall dependency of the hydrogen adsorption on the promoter concentrations is rather low since H_2 adsorption in the range of the μ_{H} expected under CO_2 or CO hydrogenation conditions is exclusively located on the three-fold adsorption site on three Cu atoms on the (111)-facet.⁹³ As long as this three-fold adsorption site is present on the (111)-facet (as long as the promoter concentration is 25% or lower), H^* adsorption under reaction conditions is observed. Nevertheless, promoter atoms in close proximity still decrease hydrogen adsorption.

3.3.3 Adsorption of CO. The points of CO adsorption (lowest μ_{CO} where the slab model with one or multiple adsorbed CO^* is more stable than the clean surface) depending on the Zn-/Ga-concentration are shown in Fig. 9. We refer the reader to the ESI† for the individual plots. The yellow-red area indicates μ_{CO} expected under CO_2 hydrogenation conditions ($X_{\text{CO}_2} = 1$, $-2.00 \text{ eV} < \mu_{\text{CO}} < -0.96 \text{ eV}$) with the colour gradient indicating the progressing RWGS-reaction (yellow for low conversion, red for high conversion; μ_{CO} increases with increasing the RWGS conversion). The green line indicates the expected μ_{CO} under CO hydrogenation conditions ($X_{\text{CO}_2} = 0$, $\mu_{\text{CO}} = -0.882$) with no changes for μ_{CO} with progressing WGS-conversion.

CO adsorption shows a similar trend to hydrogen adsorption: CO^* adsorbs almost exclusively on the adsorption sites composed by only Cu, however, the facet-dependency is much weaker for CO adsorption than for hydrogen adsorption, *i.e.* CO adsorption is also observed for facets other than the (111)-facet. Similarly to hydrogen adsorption, the CO^* -coverage on the surface decreases when the promoter contraction increases. Nevertheless, in contrast to the formation of surface H^* , the promoters also influence the point of CO adsorption, predicting CO^* -free surfaces for both the Cu-Ga- and Cu-Zn-systems under CO_2 hydrogenation conditions for small quantities of either promoter. Higher promoter concentrations (>25%) further decrease CO adsorption, meaning CO^* -free surfaces are predicted even under CO hydrogenation conditions.

CO adsorption on the other facets follows the same trend, but the initial decrease in CO adsorption upon adding promoters is

stronger than for the (111)-facet. Again, this can be explained by the spatial distribution of the promoter in the slabs: Zn or Ga added to the (100)- and (110)-facet of Cu are preferentially located on the surface, preventing CO adsorption on the preferred sites composed by only Cu atoms. Thus, both the (100)- and (110)-facet of the CuGa- and CuZn-alloys are expected to remain free of adsorbed CO^* even under CO hydrogenation reaction conditions, while the equivalent structure of pure Cu is expected to be covered with a substantial amount of CO^* .

In summary, we showed that both Ga and Zn significantly decrease CO adsorption under CO_x hydrogenation reaction conditions, especially if the promoters are located on the outmost layer of the slab. The (111)-facet of the alloys is the only facet that can still adsorb CO^* under CO_x hydrogenation reaction conditions for low promoter concentrations (<25%), and this adsorption is located exclusively on the three-fold adsorption site composed by three Cu atoms. Since hydrogen adsorption is also located on the same adsorption site, CO and hydrogen adsorption compete under reaction conditions, with the latter being thermodynamically slightly more favoured over CO adsorption. As a result, CO in the active atmosphere helps promoting realloying, and might increase the dynamics of the system under reaction conditions (both oxygen adsorption, surface oxidation and dealloying and hydrogen/CO adsorption and realloying being simultaneously possible). Nevertheless, if the atmospheric concentration of CO becomes to high, CO^* can replace large quantities of adsorbed H^* on the particle surface, thus presumably reducing the overall activity of the catalysts.



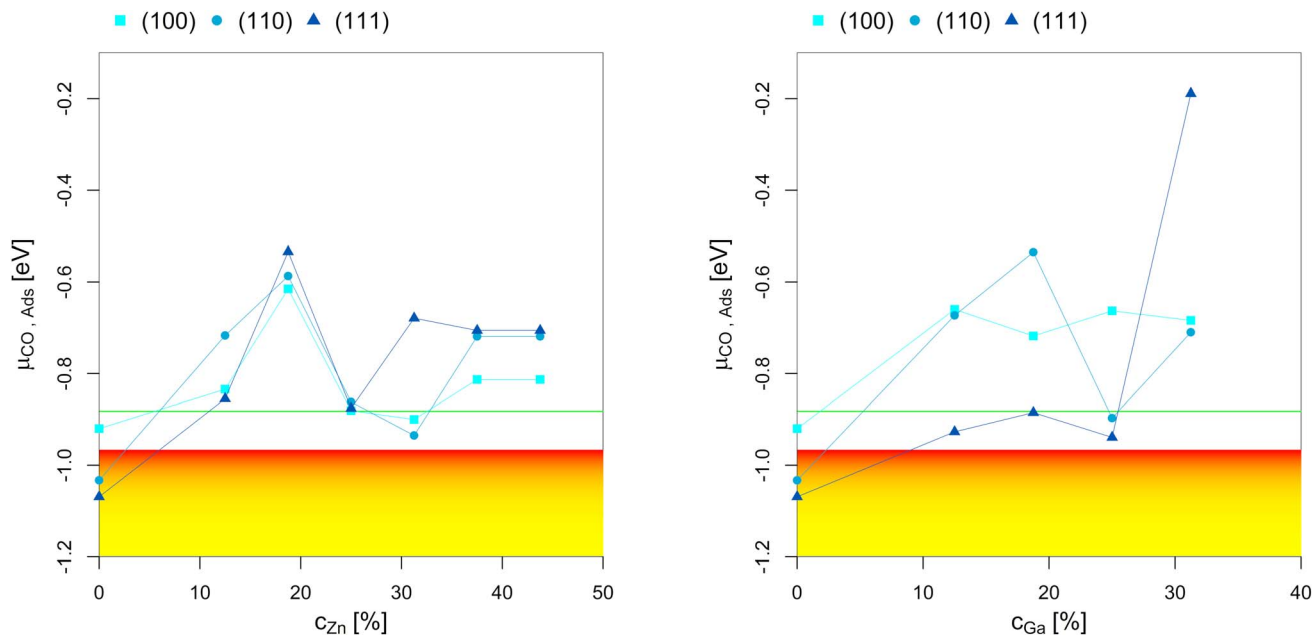


Fig. 9 Points of CO adsorption (lowest μ_{CO} where the slab model with one or multiple adsorbed CO^* is more stable than the clean surface) depending on the Zn-/Ga-concentration in the fcc-Cu crystal lattice. The yellow-red area indicates μ_{CO} expected under CO_2 hydrogenation conditions ($X_{\text{CO}_2} = 1$, $-2.00 \text{ eV} < \mu_{\text{CO}} < -0.96 \text{ eV}$) with the colour gradient indicating the progressing RWGS-reaction (yellow for low conversion, red for high conversion; μ_{CO} increases with increasing RWGS-conversion). The green line indicates the expected μ_{CO} under CO hydrogenation conditions ($X_{\text{CO}_2} = 0$, $\mu_{\text{CO}} = -0.882$) with no changes for μ_{CO} with progressing WGS-conversion.

3.3.4 Adsorption of CO_2 . Similarly to what was found for pure Cu,⁹³ none of the investigated stoichiometries and facets adsorbs CO_2 in their metallic state. But if the promoters are oxidised to their respective (sub)oxides, CO_2 adsorption at the metal-oxide interface should be facilitated, similarly to what was shown for Cu NPs supported on non-reducible oxide supports.^{68,92} For a more in-depth discussion of CO_2 adsorption under CO_2 and CO hydrogenation conditions, we refer the reader to the ESI.†

3.4 Effect of gas phase composition – comparing CO_2 vs. CO hydrogenation conditions

The oxygen chemical potential (μ_{O}) at which oxygen starts adsorbing depending on the promoter concentration (Fig. 7) shows that for μ_{O} corresponding to CO_2 hydrogenation conditions ($X_{\text{CO}_2} = 1$, $X_{\text{CO}} = 0$, $-3.11 \text{ eV} < \mu_{\text{O}} < -2.92 \text{ eV}$ corresponding to $p_{\text{O}_2} < 10^{-48}$ bar at 500 K, yellow-red area), the surface facets of the Cu-Ga-alloy slabs are at least partially covered with O^* under reaction conditions. Thus, while the pure Cu surfaces are expected to be fully reduced under reaction conditions,⁹³ adding Ga to the slab significantly increases the oxygen affinity of the system. As a result, a partial or full Ga oxide overlayer is predicted under CO_2 hydrogenation conditions for the Cu-Ga alloys. This can significantly influence the catalyst reactivity (from increased selectivity to reduced activity⁷²). The formation of GaO_x -island on metallic Cu under CO_2 hydrogenation conditions has indeed been observed experimentally for selected model systems using STEM and XPS.¹¹⁸ In contrast, if μ_{O} decreases significantly, as predicted for pure CO hydrogenation conditions ($X_{\text{CO}_2} = 0$, $X_{\text{CO}} = 1$, -4.07 eV

$< \mu_{\text{O}} < -3.93 \text{ eV}$, equivalent to $p_{\text{O}_2} < 10^{-69}$ bar at 500 K), oxygen adsorption is no longer favourable. For a mixture of CO_2 and CO ($0.1 < X_{\text{CO}_2} < 0.9$), μ_{O} increases linearly as a function of X_{CO_2} (as described in Fig. 1). As a result, if small quantities of CO_2 are present in the gas feed, oxygen adsorption is possible under reaction conditions, but depends quite strongly on X_{CO_2} .

Zn shows a similar effect as Ga as a promoter in the slab; it significantly increases the oxygen affinity of the alloy when compared to pure Cu. But the Cu-Zn-system is more reluctant towards oxygen adsorption than the Cu-Ga-system. Thus, while the Cu-Zn-system is expected to form Zn (sub)-oxide layers on top of the slab models under CO_2 hydrogenation reaction conditions, less surface oxidation is predicted based on our calculations. As a result, partially reduced Zn layers and alloyed surfaces are expected under CO_2 hydrogenation conditions, which is consistent with experimental results.⁷¹ For the effect of the feedstock composition, the same trend as for the Cu-Ga-system is observed: Reducing the amount of CO_2 in the feedstock decreases μ_{O} , which results in less adsorbed O^* and more reduced surfaces.

The effect of the other adsorbates (H^* and CO^*) is opposite to O^* : the $\mu_{\text{H}}/\mu_{\text{CO}}$ at which hydrogen/CO starts adsorbing depends on the promoter concentration. Fig. 8 and 9 show that adding Ga or Zn to the slab surface decreases the adsorption of H^* and CO^* on the slab. In addition, since the adsorption is preferred on sites composed only by Cu, adding H_2/CO to the feedstock decreases the dealloying and surface oxidation (as observed for oxygen adsorption) while stabilising the reduced and alloyed slabs. As a result, both dealloying and realloying processes are thermodynamically possible under CO_x



hydrogenation conditions with the preferred state depending on the gas phase compositions (associated with μ_{O} , μ_{H} and μ_{CO}): more CO_2 results in higher μ_{O} and more oxidised surfaces, while less CO_2 leads to a lower μ_{O} and more reduced surfaces. Nevertheless, as shown in Fig. 1, the initial decrease in μ_{O} due to replacing CO_2 with CO is rather small (the same applies to μ_{H} and μ_{CO}), resulting in small changes of the surface state of the catalyst if the gas phase still contains small quantities of CO_2 (compare Table S1 in the ESI†). Thus, unless all CO_2 is removed from the feedstock, surface oxidation remains possible under CO_x hydrogenation conditions.

3.5 Comparison of Ga and Zn as promoters

Comparing Ga and Zn as promoters for Cu reveals key similarities and differences. Solid solutions for both Ga and Zn in Cu are known, both having an fcc-crystal structure for high enough Cu concentrations (Fig. 2) with the promoters (Ga/Zn) being statistically distributed by randomly replacing Cu in the fcc-Cu crystal lattice with the promoters. The top-most layer of the surface is rich in Zn or Ga (in the respective alloys) and both promoters facilitate oxygen adsorption when compared to pure copper, favouring partial or full dealloying and forming partial or fully oxidized $\text{GaO}_x/\text{ZnO}_x$ surface layers. Also, both promoters decrease hydrogen and CO adsorption, especially in the absence of surface sites composed only of Cu atoms.

However, our study also highlights the differences between Zn and Ga as promoters. The stabilisation due to alloying is much larger for the fcc-CuGa solid solution in comparison to the equivalent fcc-CuZn solid solution for low promoter (Ga/Zn) concentration, as seen in the alloy energies in Fig. 3. But stabilisation due to the formation of a fcc-solid solution is only observed for Ga-contents below 25%, while there is a net stabilisation for the Cu-Zn-system up to a Zn-content of 50%. Partial or full surface oxidation under CO_2 hydrogenation reaction conditions is expected for the Cu-Ga-system (depending on the feedstock compositions), while the Cu-Zn-system is expected to be partially oxidised at most.

4 Conclusions

In this article, we used DFT calculations in combination with *ab initio* atomistic thermodynamics to investigate the surface coverage of multiple key adsorbates on several surfaces for both the fcc-CuGa and fcc-CuZn solid solution with different stoichiometric compositions evaluating different spatial distributions of the two metals in the slab. Our calculations predict that both systems form fcc-solid solutions when the promoter concentration is low enough, with fcc-CuZn being stable for much higher Zn-concentrations than its corresponding Ga-equivalent, in agreement with the experimental phase diagrams. The outmost surface layers of the built surface models are the most impactful on the surface energy and surface coverage, while the subsurface layers have a rather low overall influence on them.

We also show that when Ga and Zn are on the surface, oxygen adsorption is more favoured compared to pure copper

due to formation of surface (sub)oxides. This is especially pronounced for the Cu-Ga-system, in which the stabilisation is so strong that under pure CO_2 and mixed CO_2/CO hydrogenation reaction conditions (small but significant quantities of CO_2 present in the gas phase), partial to full dealloying and oxidation of all surface Ga to GaO_x is possible. This effect is much weaker for the Cu-Zn-system, so that only partial dealloying and oxidation to surface ZnO_x is possible under pure CO_2 and mixed CO_2/CO hydrogenation reaction conditions with some Zn retained in its metallic form, while all Ga is at least partially oxidised and no metallic Ga is expected. In contrast, the pure Cu-system was shown to be fully reduced under reaction conditions.⁹³ This trend agrees well with experimental data, where Ga is fully oxidised and Zn is partially oxidised under pure CO_2 and mixed CO_2/CO hydrogenation reaction conditions.^{71,72}

Regarding hydrogen/CO adsorption, we showed that the stability of surface H^*/CO^* is lower on both the fcc-CuGa and fcc-CuZn solid solutions in comparison to pure Cu, being especially pronounced in the absence of adsorption sites composed by only Cu atoms (observed for high promoter concentrations close to the surface), since both adsorbates exclusively adsorb on surface sites adjacent to multiple Cu atoms under CO_2 and CO hydrogenation reaction conditions. As a result, both hydrogen and CO adsorption result in surface reduction and thus realloying of the surface slabs, but the stability of the resulting surface H^*/CO^* depends strongly on the active atmosphere. Finally, the adsorption of molecular CO_2 is not observed on any of the investigated facets/stoichiometries.

Concerning the influence of the active atmosphere, we showed that the CO/CO_2 hydrogenation conditions, usually perceived as rather reducing due to the presence of large quantities of hydrogen gas, are actually quite oxidising for both Zn and especially for Ga. This oxidising potential can be tuned by modifying the ratio of CO_2 to CO in the feed gas. Only pure CO/H_2 -feeds are reducing enough so that fully reduced alloys are thermodynamically preferred under reaction conditions, while the addition of small quantities of CO_2 increases the oxygen chemical potential (μ_{O}) so much that the formation of ZnO_x and GaO_x surface species become thermodynamically possible, likely forming a dynamic system enabling oxidation and reduction processes to occur at the same time.

The use of small quantities of CO_2 to accelerate the hydrogenation of CO in an industrial setting is well-known and is mainly attributed to the influence of CO_2 onto the reaction kinetics. We propose an alternative and complementary explanation based on thermodynamic properties: having small amounts of CO_2 in the feed increases the oxygen chemical potential, resulting in a more oxidising atmosphere, which favours the formation of Zn or Ga (sub)oxides, and thereby generates interfacial sites between Cu and MO_x directly on the metal surface, which could promote methanol formation similarly to the mechanism observed for the interface of Cu with non-reducible supports.^{68,92} This site formation depends on the gas phase composition, since the partial pressures of the individual adsorbates determine their specific chemical



potentials under reaction conditions and thus their interplay with the catalytic surface. Without any CO₂, no active sites are formed, while a pure CO₂ feed is so oxidising that a closed oxide overlayer can form without any accessible Cu sites remaining, an aspect especially pronounced for Ga. Thus, carefully tuning the reaction conditions is of high importance for both the catalytic activity and stability under reaction conditions, at par with catalyst composition and particle preparation.

In conclusion, we found that Ga and Zn as promoters in Cu-based CO/CO₂ hydrogenation significantly influence the interaction of the catalyst with the gas phase. Both fcc-CuGa and fcc-CuZn solid solutions show increased oxygen adsorption but reduced hydrogen and CO adsorption with respect to pure Cu, especially for high promoter concentrations. The gas phase composition can in return strongly interplay with the particle surface and form a dynamic equilibrium relevant for the reaction mechanism. Nevertheless, such effects are impossible to capture using static DFT, highlighting the need for dynamic calculations such as *ab initio* molecular dynamics of discrete Cu alloy particles to investigate the dynamics of the catalyst structure under reaction conditions. We are currently pursuing this effort and will disclose our work in due course.

Data availability

The data used in this manuscript (including structures as .cif-files, the corresponding energies in a .txt-file as well as typical INCAR and KPOINTS files used in the calculations) are available on Zenodo (a general-purpose open repository developed under the European OpenAIRE program and operated by CERN) at <https://doi.org/10.5281/zenodo.7267564>.

Author contributions

CRedit (Contributor Roles Taxonomy)¹¹⁹ was used for standardised contribution descriptions: Andreas Müller: data curation, formal analysis, investigation, methodology, software, validation, visualization, writing (original draft), writing (review & editing). Aleix Comas-Vives: data curation, conceptualization, funding acquisition, methodology, supervision, validation, writing (review & editing). Christophe Copéret: conceptualization, funding acquisition, project administration, resources, supervision, writing (review & editing).

Conflicts of interest

There are no conflicts of interest to declare.

Acknowledgements

This research is supported by the Sinergia Project of the Swiss National Science Fonds (SNSF fond number: CRSII5_183495). All DFT calculations were performed using the EULER cluster at ETH Zurich managed by the HPC team. Andreas Müller thanks Jan Alfke, Dr María Tejada Serrano, Dr Gina Noh and Mélissa Zacharias for the discussion on the catalytic system and suggestions of further calculations and on the visual

representation of the data. Aleix Comas-Vives thanks the Spanish MEC and the European Social Fund (RyC-2016-19930) and the Spanish “Ministerio de Ciencia, Innovación y Universidades” (PGC2018-100818-A-I00).

Notes and references

- S. Solomon, G.-K. Plattner, R. Knutti and P. Friedlingstein, *Proc. Natl. Acad. Sci.*, 2009, **106**, 1704–1709.
- G. A. Olah, *Angew. Chem., Int. Ed.*, 2005, **44**, 2636–2639.
- A. Goeppert, M. Czaun, J.-P. Jones, G. K. Surya Prakash and G. A. Olah, *Chem. Soc. Rev.*, 2014, **43**, 7995–8048.
- S. Simon Araya, V. Liso, X. Cui, N. Li, J. Zhu, S. L. Sahlin, S. H. Jensen, M. P. Nielsen and S. K. Kær, *Energies*, 2020, **13**, 596.
- A. Álvarez, A. Bansode, A. Urakawa, A. V. Bavykina, T. A. Wezendonk, M. Makkee, J. Gascon and F. Kapteijn, *Chem. Rev.*, 2017, **117**, 9804–9838.
- A. Álvarez, M. Borges, J. J. Corral-Pérez, J. G. Olcina, L. Hu, D. Cornu, R. Huang, D. Stoian and A. Urakawa, *ChemPhysChem*, 2017, **18**, 3135–3141.
- K. Stangeland, H. Li and Z. Yu, *Ind. Eng. Chem. Res.*, 2018, **57**, 4081–4094.
- C. V. Miguel, M. A. Soria, A. Mendes and L. M. Madeira, *J. Nat. Gas Sci. Eng.*, 2015, **22**, 1–8.
- M. Behrens, F. Studt, I. Kasatkin, S. Köhl, M. Hävecker, F. Abild-Pedersen, S. Zander, F. Girgsdies, P. Kurr, B.-L. Kniep, M. Tovar, R. W. Fischer, J. K. Nørskov and R. Schlögl, *Science*, 2012, **336**, 893–897.
- S. Polierer, D. Guse, S. Wild, K. Herrera Delgado, T. N. Otto, T. A. Zevaco, M. Kind, J. Sauer, F. Studt and S. Pitter, *Catalysts*, 2020, **10**, 816.
- R. Horn and R. Schlögl, *Catal. Lett.*, 2015, **145**, 23–39.
- S. Dang, H. Yang, P. Gao, H. Wang, X. Li, W. Wei and Y. Sun, *Catal. Today*, 2019, **330**, 61–75.
- K. Jalama, *Catal. Rev.*, 2017, **59**, 95–164.
- H. Yang, C. Zhang, P. Gao, H. Wang, X. Li, L. Zhong, W. Wei and Y. Sun, *Catal. Sci. Technol.*, 2017, **7**, 4580–4598.
- O. A. Ojelade and S. F. Zaman, *Catal. Surv. Asia*, 2020, **24**, 11–37.
- O. Tursunov, L. Kustov and A. Kustov, *Oil Gas Sci. Technol.*, 2017, **72**, 30.
- W. Li, H. Wang, X. Jiang, J. Zhu, Z. Liu, X. Guo and C. Song, *RSC Adv.*, 2018, **8**, 7651–7669.
- A. Pustovarenko, A. Dikhtiarenko, A. Bavykina, L. Gevers, A. Ramírez, A. Russkikh, S. Telalovic, A. Aguilar, J.-L. Hazemann, S. Ould-Chikh and J. Gascon, *ACS Catal.*, 2020, **10**, 5064–5076.
- A. Ramirez, L. Gevers, A. Bavykina, S. Ould-Chikh and J. Gascon, *ACS Catal.*, 2018, **8**, 9174–9182.
- A. Dokania, A. Ramirez, A. Bavykina and J. Gascon, *ACS Energy Lett.*, 2019, **4**, 167–176.
- D. Stoian, A. Bansode, F. Medina and A. Urakawa, *Catal. Today*, 2017, **283**, 2–10.
- A. Bansode and A. Urakawa, *ACS Catal.*, 2014, **4**, 3877–3880.
- B. Tidona, C. Koppold, A. Bansode, A. Urakawa and P. Rudolf von Rohr, *J. Supercrit. Fluids*, 2013, **78**, 70–77.



- 24 A. Álvarez, M. Borges, J. J. Corral-Pérez, J. G. Olcina, L. Hu, D. Cornu, R. Huang, D. Stoian and A. Urakawa, *ChemPhysChem*, 2017, **18**, 3087.
- 25 G. Liu, P. Poths, X. Zhang, Z. Zhu, M. Marshall, M. Blankenhorn, A. N. Alexandrova and K. H. Bowen, *J. Am. Chem. Soc.*, 2020, **142**, 7930–7936.
- 26 J.-Y. Liu, X.-Q. Gong, R. Li, H. Shi, S. B. Cronin and A. N. Alexandrova, *ACS Catal.*, 2020, **10**, 4048–4058.
- 27 J.-Y. Liu, X.-Q. Gong and A. N. Alexandrova, *J. Phys. Chem. C*, 2019, **123**, 3505–3511.
- 28 W. J. Durand, A. A. Peterson, F. Studt, F. Abild-Pedersen and J. K. Nørskov, *Surf. Sci.*, 2011, **605**, 1354–1359.
- 29 E. L. Kunkes, F. Studt, F. Abild-Pedersen, R. Schlögl and M. Behrens, *J. Catal.*, 2015, **328**, 43–48.
- 30 M. S. Duyar, C. Tsai, J. L. Snider, J. A. Singh, A. Gallo, J. S. Yoo, A. J. Medford, F. Abild-Pedersen, F. Studt, J. Kibsgaard, S. F. Bent, J. K. Nørskov and T. F. Jaramillo, *Angew. Chem., Int. Ed.*, 2018, **57**, 15045–15050.
- 31 S. Polierer, J. Jelic, S. Pitter and F. Studt, *J. Phys. Chem. C*, 2019, **123**, 26904–26911.
- 32 J. Tai, Q. Ge, R. J. Davis and M. Neurock, *J. Phys. Chem. B*, 2004, **108**, 16798–16805.
- 33 J. Medina-Ramos, W. Zhang, K. Yoon, P. Bai, A. Chemburkar, W. Tang, A. Atifi, S. S. Lee, T. T. Fister, B. J. Ingram, J. Rosenthal, M. Neurock, A. C. T. van Duin and P. Fenter, *Chem. Mater.*, 2018, **30**, 2362–2373.
- 34 S. N. Steinmann, C. Michel, R. Schwiedernoch, J.-S. Filhol and P. Sautet, *ChemPhysChem*, 2015, **16**, 2307–2311.
- 35 W. Guo, C. Michel, R. Schwiedernoch, R. Wischert, X. Xu and P. Sautet, *Organometallics*, 2014, **33**, 6369–6380.
- 36 M. J. Kolb, D. Loffreda, P. Sautet and F. Calle-Vallejo, *J. Catal.*, 2021, **395**, 136–142.
- 37 V. D. Dasireddy and B. Likozar, *Renewable Energy*, 2019, **140**, 452–460.
- 38 Y. Ma, J. Wang, K. R. Goodman, A. R. Head, X. Tong, D. J. Stacchiola and M. G. White, *J. Phys. Chem. C*, 2020, **124**, 22158–22172.
- 39 J. Kim, B. B. Sarma, E. Andrés, N. Pfänder, P. Concepción and G. Prieto, *ACS Catal.*, 2019, **9**, 10409–10417.
- 40 W. Wang, Z. Qu, L. Song and Q. Fu, *J. Catal.*, 2020, **382**, 129–140.
- 41 M. S. Tameh, A. K. Dearden and C. Huang, *J. Phys. Chem. C*, 2018, **122**, 17942–17953.
- 42 R. M. Palomino, P. J. Ramírez, Z. Liu, R. Hamlyn, I. Waluyo, M. Mahapatra, I. Orozco, A. Hunt, J. P. Simonovis, S. D. Senanayake and J. A. Rodriguez, *J. Phys. Chem. B*, 2018, **122**, 794–800.
- 43 G. Noh, E. Lam, D. T. Bregante, J. Meyet, P. Šot, D. W. Flaherty and C. Copéret, *Angew. Chem., Int. Ed.*, 2021, **60**, 9650–9659.
- 44 T. P. Araújo, A. H. Hergesell, D. Faust-Akl, S. Büchele, J. A. Stewart, C. Mondelli and J. Pérez-Ramírez, *ChemSusChem*, 2021, **14**, 2914–2923.
- 45 M. Zabilskiy, K. Ma, A. Beck and J. A. van Bokhoven, *Catal. Sci. Technol.*, 2021, **11**, 349–358.
- 46 M. Zabilskiy, V. L. Sushkevich, D. Palagin, M. A. Newton, F. Krumeich and J. A. van Bokhoven, *Nat. Commun.*, 2020, **11**, 2409.
- 47 A. Ye, Z. Li, J. Ding, W. Xiong and W. Huang, *ACS Catal.*, 2021, **11**, 10014–10019.
- 48 C. Göbel, S. Schmidt, C. Froese, Q. Fu, Y.-T. Chen, Q. Pan and M. Muhler, *J. Catal.*, 2020, **383**, 33–41.
- 49 Y. Lu, R. Zhang, B. Cao, B. Ge, F. F. Tao, J. Shan, L. Nguyen, Z. Bao, T. Wu, J. W. Pote, B. Wang and F. Yu, *ACS Catal.*, 2017, **7**, 5500–5512.
- 50 N. D. Nielsen, J. Thrane, A. D. Jensen and J. M. Christensen, *Catal. Lett.*, 2020, **150**, 1427–1433.
- 51 Z. Zhang, X. Chen, J. Kang, Z. Yu, J. Tian, Z. Gong, A. Jia, R. You, K. Qian, S. He, B. Teng, Y. Cui, Y. Wang, W. Zhang and W. Huang, *Nat. Commun.*, 2021, **12**, 4331.
- 52 C. Du, P. Lu and N. Tsubaki, *ACS Omega*, 2020, **5**, 49–56.
- 53 S. Dey, G. C. Dhal, D. Mohan and R. Prasad, *Adv. Compos. Hybrid Mater.*, 2019, **2**, 626–656.
- 54 W. U. Khan, L. Baharudin, J. Choi and A. C. K. Yip, *ChemCatChem*, 2021, **13**, 111–120.
- 55 Clariant, *Clariant's MegaMax® 800 Catalyst Achieved an Outstanding Performance at CNOOC'S Davy Methanol Plant in Hainan, China*, 2021, <https://www.clariant.com/en/Corporate/News/2021/07/Clariant's-MegaMax-reg-800-catalyst-achieved-an-outstanding-performance-at-CNOOC's-Davy-metha>.
- 56 K. Girod, H. Lohmann and S. Kaluza, *Chem. Ing. Tech.*, 2021, **93**, 850–855.
- 57 J. Matthey, *Decarbonising Chemicals Production*, 2021, <https://matthey.com/en/about-us/transitions/decarbonising-chemicals-production>.
- 58 S. Energy, *A New Hydrogen Reality: Fuel from Wind and Water*, 2021, <https://www.siemens-energy.com/global/en/offerings/renewable-energy/hydrogen-solutions/haruni.html>.
- 59 C. R. International, *Recycling CO₂ to Produce Methanol - A Question of Balance*, 2021, <https://www.carbonrecycling.is/co2-methanol>.
- 60 F. Studt, M. Behrens, E. L. Kunkes, N. Thomas, S. Zander, A. Tarasov, J. Schumann, E. Frei, J. B. Varley, F. Abild-Pedersen, J. K. Nørskov and R. Schlögl, *ChemCatChem*, 2015, **7**, 1105–1111.
- 61 J. Nakamura, T. Fujitani and Y.-S. Choi, *Top. Catal.*, 2003, **22**, 277–285.
- 62 M. B. Fichtl, D. Schlereth, N. Jacobsen, I. Kasatkin, J. Schumann, M. Behrens, R. Schlögl and O. Hinrichsen, *Appl. Catal., A*, 2015, **502**, 262–270.
- 63 A. Urakawa, *Nat. Catal.*, 2021, **4**, 447–448.
- 64 A. Beck, M. Zabilskiy, M. A. Newton, O. Safonova, M. G. Willinger and J. A. van Bokhoven, *Nat. Catal.*, 2021, **4**, 488–497.
- 65 C. Copéret, A. Fedorov and P. A. Zhizhko, *Catal. Lett.*, 2017, **147**, 2247–2259.
- 66 C. Copéret, *Acc. Chem. Res.*, 2019, **52**, 1697–1708.
- 67 S. Tada, F. Watanabe, K. Kiyota, N. Shimoda, R. Hayashi, M. Takahashi, A. Nariyuki, A. Igarashi and S. Satokawa, *J. Catal.*, 2017, **351**, 107–118.



- 68 K. Larmier, W.-C. Liao, S. Tada, E. Lam, R. Verel, A. Bansode, A. Urakawa, A. Comas-Vives and C. Copéret, *Angew. Chem., Int. Ed.*, 2017, **56**, 2318–2323.
- 69 K. Fujiwara, S. Tada, T. Honma, H. Sasaki, M. Nishijima and R. Kikuchi, *AICHE J.*, 2019, **65**, e16717.
- 70 S. R. Docherty and C. Copéret, *J. Am. Chem. Soc.*, 2021, **143**, 6767–6780.
- 71 E. Lam, G. Noh, K. Larmier, O. V. Safonova and C. Copéret, *J. Catal.*, 2021, **394**, 266–272.
- 72 E. Lam, G. Noh, K. W. Chan, K. Larmier, D. Lebedev, K. Searles, P. Wolf, O. V. Safonova and C. Copéret, *Chem. Sci.*, 2020, **11**, 7593–7598.
- 73 A. M. Hengne, D. J. Yuan, N. S. Date, Y. Saih, S. P. Kamble, C. V. Rode and K.-W. Huang, *Ind. Eng. Chem. Res.*, 2019, **58**, 21331–21340.
- 74 A. M. Hengne, K. D. Bhatte, S. Ould-Chikh, Y. Saih, J. M. Basset and K.-W. Huang, *ChemCatChem*, 2018, **10**, 1360–1369.
- 75 V. D. Dasireddy and B. Likozar, *Renewable Energy*, 2019, **140**, 452–460.
- 76 E. L. Kunkes, F. Studt, F. Abild-Pedersen, R. Schlögl and M. Behrens, *J. Catal.*, 2015, **328**, 43–48.
- 77 J. He, D. Laudenschleger, J. Schittkowski, A. Machoke, H. Song, M. Muhler, R. Schlögl and H. Ruland, *Chem. Ing. Tech.*, 2020, **92**, 1525–1532.
- 78 Y. Slotboom, M. Bos, J. Pieper, V. Vrieswijk, B. Likozar, S. Kersten and D. Brillman, *Chem. Eng. J.*, 2020, **389**, 124181.
- 79 J. Thrane, S. Kuld, N. D. Nielsen, A. D. Jensen, J. Sehested and J. M. Christensen, *Angew. Chem., Int. Ed.*, 2020, **59**, 18189–18193.
- 80 A. Pavlišić, M. Huš and B. Likozar, *J. Cleaner Prod.*, 2020, **275**, 122958.
- 81 A. Pavlišić, M. Huš and B. Likozar, *J. Cleaner Prod.*, 2020, **275**, 122958.
- 82 J.-B. Li, L. Ji, J. Liang, Y. Zhang, J. Luo, C. Li and G. Rao, *Calphad*, 2008, **32**, 447–453.
- 83 M. Kowalski and P. J. Spencer, *J. Phase Equilib.*, 1993, **14**, 432–438.
- 84 B. W. J. Chen, L. Xu and M. Mavrikakis, *Chem. Rev.*, 2021, **121**, 1007–1048.
- 85 L. Grajciar, C. J. Heard, A. A. Bondarenko, M. V. Polynski, J. Meeprasert, E. A. Pidko and P. Nachtigall, *Chem. Soc. Rev.*, 2018, **47**, 8307–8348.
- 86 C.-R. Chang, Z.-Q. Huang and J. Li, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2016, **6**, 679–693.
- 87 Á. Morales-García, F. Viñes, J. R. B. Gomes and F. Illas, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2021, **11**, e1530.
- 88 X. Zhang, T. E. Sewell, B. Glatz, S. Sarupria and R. B. Getman, *Catal. Today*, 2017, **285**, 57–64.
- 89 N. Artrith, *Matter*, 2020, **3**, 985–986.
- 90 M. T. Darby, M. Stamatakis, A. Michaelides and E. C. H. Sykes, *J. Phys. Chem. Lett.*, 2018, **9**, 5636–5646.
- 91 X. Li, R. Chiong and A. J. Page, *J. Phys. Chem. Lett.*, 2021, **12**, 5156–5162.
- 92 E. Lam, J. J. Corral-Pérez, K. Larmier, G. Noh, P. Wolf, A. Comas-Vives, A. Urakawa and C. Copéret, *Angew. Chem., Int. Ed.*, 2019, **58**, 13989–13996.
- 93 A. Müller, A. Comas-Vives and C. Copéret, *J. Phys. Chem. C*, 2021, **125**, 396–409.
- 94 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1993, **47**, 558–561.
- 95 G. Kresse and J. Furthmüller, *Comput. Mater. Sci.*, 1996, **6**, 15–50.
- 96 G. Kresse and J. Furthmüller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169–11186.
- 97 P. E. Blöchl, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **50**, 17953–17979.
- 98 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1999, **59**, 1758–1775.
- 99 G. Kresse and J. Hafner, *J. Phys.: Condens. Matter*, 1994, **6**, 8245–8257.
- 100 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- 101 J. Paier, R. Hirschl, M. Marsman and G. Kresse, *J. Chem. Phys.*, 2005, **122**, 234102.
- 102 K. Momma and F. Izumi, *J. Appl. Crystallogr.*, 2011, **44**, 1272–1276.
- 103 K. Reuter and M. Scheffler, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2001, **65**, 035406.
- 104 K. Reuter and M. Scheffler, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2003, **68**, 045407.
- 105 M. W. Chase, *NIST-JANAF Thermochemical Tables*, American Institute of Physics, 4th edn, 1998.
- 106 A. Roussey, P. Gentile, D. Lafond, E. Martinez, V. Jousseume, C. Thieuleux and C. Copéret, *J. Mater. Chem. C*, 2013, **1**, 1583–1587.
- 107 T. Reichenbach, M. Walter, M. Moseler, B. Hammer and A. Bruix, *J. Phys. Chem. C*, 2019, **123**, 30903–30916.
- 108 J. M. Moe, *Chem. Eng. Prog.*, 1962, **58**, 33–36.
- 109 H. Okamoto, *J. Phase Equilib. Diffus.*, 2016, **37**, 350–362.
- 110 H. Okamoto, *J. Phase Equilib. Diffus.*, 2018, **39**, 87–100.
- 111 G. Brauer, *Handbuch der Präparativen Anorganischen Chemie*, Ferdinand Enke, Stuttgart, 1978, vol. 3.
- 112 A. F. Wells, *Structural Inorganic Chemistry*, Clarendon Press, Oxford University Press, 5th edn, 1984.
- 113 B. D. Sharma and J. Donohue, *Z. Kristallogr. Cryst. Mater.*, 1962, **117**, 293–300.
- 114 W. M. Haynes, *CRC Handbook of Chemistry and Physics*, CRC Press, 2nd edn, 2011.
- 115 J. Yan and J. K. Nørskov, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2013, **88**, 245204.
- 116 E. Frei, A. Gaur, H. Lichtenberg, L. Zwiener, M. Scherzer, F. Girgsdies, T. Lunkenbein and R. Schlögl, *ChemCatChem*, 2020, **12**, 4029–4033.
- 117 X. Huang, T. Jones, A. Fedorov, R. Farra, C. Copéret, R. Schlögl and M.-G. Willinger, *Adv. Mater.*, 2021, **33**, 2101772.
- 118 J.-Q. Zhong, S. Shaikhutdinov and B. Roldan Cuenya, *J. Phys. Chem. C*, 2021, **125**, 1361–1367.
- 119 CRediT – Contributor Roles Taxonomy, 2019, <https://credit.niso.org/>.

