Characterisation of multipartite entanglement beyond the single-copy paradigm

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* But not always possible, \exists undistillable entanglement, in particular, all PPT-entangled states are undistillable

Quantum information processing with multiple copies

But typical scenario in QIP involves multiple parties 1, 2, 3, ..., N

 \mathcal{B}_{Λ}

 \mathcal{B}_6

 $|\mathcal{A}_2|$

 $\mathcal{A}_{\it A}$

This is the situation we consider here, we ask: What is the structure of genuine multipartite Base of the structure of genuine multipartite base of the structure of genuine multipartite structure structure of genuine multipartite structure structure of genuine multipartite structure stru

In particular, if individual copies are biseparable, B is joint state biseparable with parition C_3 $A_1B_1C_1...A_2B_2C_2...A_NB_NC_N...?$

Hayata Ygmasaki, Simon Morelli, Markuş Miethlinger, Jessica Bavaresco, NF, BMarcus Huber, Quantum 6, 695 (2022)

C. Palazuelos and J. I. de Vicente, Quantum 6, 735 (2022)
 and multiple copies A, B, C, ...



Genuine Multipartite Entanglement

Pure state $|\Phi^{(k)}\rangle$ separable w.r.t. to k-partition $\mathcal{A}_1|\mathcal{A}_2|\dots|\mathcal{A}_k$ if $|\Phi^{(k)}\rangle = \bigotimes_{i=1}^k |\phi_{\mathcal{A}_i}\rangle$

Mixed state $\rho^{(k)}$ is called k-separable if it can be written as a convex combination of pure states that are separable w.r.t. to some k-partition $\rho^{(k)} = \sum_{i} p_i |\Phi_i^{(k)}\rangle \langle \Phi_i^{(k)}|$



k-separability does not imply separability w.r.t. any specific partition except when $ho^{(k)}$ is pure or when k=N (fully separable)

Here: States that are separable w.r.t. any partition: partition separable For k = 2: biseparable

For k=1: genuinely k-partite entangled (here, just "GME")





Activation of Genuine Multipartite Entanglement

Why would it NOT be? Consider tripartite biseparable state

 $\rho_{\text{bisep}} = p \,\rho_{\mathcal{A}_1} \otimes \rho_{\mathcal{A}_2 \mathcal{A}_3} + (1-p) \,\rho_{\mathcal{A}_1 \mathcal{A}_2} \otimes \rho_{\mathcal{A}_3}$

If we take two copies: \mathcal{A}, \mathcal{B}

 $\rho_{\text{bisep}}^{\otimes 2} = p^2 \,\rho_{\mathcal{A}_1} \otimes \rho_{\mathcal{A}_2 \mathcal{A}_3} \otimes \rho_{\mathcal{B}_1} \otimes \rho_{\mathcal{B}_2 \mathcal{B}_3} + (1-p)^2 \,\rho_{\mathcal{A}_1 \mathcal{A}_2} \otimes \rho_{\mathcal{A}_3} \otimes \rho_{\mathcal{B}_1 \mathcal{B}_2} \otimes \rho_{\mathcal{B}_3} \\ + p(1-p) \left[\,\rho_{\mathcal{A}_1} \otimes \rho_{\mathcal{A}_2 \mathcal{A}_3} \otimes \rho_{\mathcal{B}_1 \mathcal{B}_2} \otimes \rho_{\mathcal{B}_3} + \rho_{\mathcal{A}_1 \mathcal{A}_2} \otimes \rho_{\mathcal{A}_3} \otimes \rho_{\mathcal{B}_1} \otimes \rho_{\mathcal{B}_2 \mathcal{B}_3} \, \right]$

Some terms not necessarily separable w.r.t. ${\cal A}_1|{\cal A}_2{\cal A}_3$, ${\cal A}_2|{\cal A}_1{\cal A}_3$ or ${\cal A}_1{\cal A}_2|{\cal A}_3$

Two copies might be GME

Indeed, example for such 2-copy activatable states have been found before [1] But is this structure tensor-stable? [1] Marcus Huber and Martin Plesch, Phys. Rev. A 83, 062321 (2011). The Tools Here, we examine GME activation in more detail Consider one-parameter family of isotropic N-qubit GHZ states $\rho(p) = p | \text{GHZ}_N \rangle \langle \text{GHZ}_N | + (1-p) \frac{1}{2^N} \mathbb{1}_{2^N} \quad \text{with} \quad | \text{GHZ}_N \rangle = \frac{1}{\sqrt{2}} (|0\rangle^{\otimes N} + |1\rangle^{\otimes N})$

We can then leverage two results:

For any N-qubit state ρ_X in X-form, a nonzero value of the genuinely multipartite concurrence $C_{\text{GM}}(\rho_X)$ [2] provides a necessary and sufficient GME criterion [3].

For any two states ρ and σ in \mathcal{H} , the Hadamard map $\mathcal{E}_{\circ}[\rho \otimes \sigma] = \frac{\rho \circ \sigma}{\operatorname{Tr}(\rho \circ \sigma)} \in \mathcal{H}$, can be implemented via SLOCC [4].

[2] Rafsanjani, Huber, Broadbent, and Eberly, Phys. Rev. A 86, 062303 (2012), arXiv:1208.2706.
[3] Ma, Chen, Chen, Spengler, Gabriel, and Huber, Phys. Rev. A 83, 062325 (2011), arXiv:1101.2001.
[4] Lami and Huber, J. Math. Phys. 57, 092201 (2016), arXiv:1603.02158.

The Works

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We observe: (1) isotropic N-qubit GHZ states are in X-form (2) Hadamard (Schur) product preserves this X-form (3) Hadamard (Schur) product can be implemented via SLOCC (4) SLOCC cannot create GME from biseparable states (5) Nonzero GM concurrence detects GME for states in X-form For one copy: If $C_{\text{GM}}(\rho(p)) > 0 \longrightarrow \rho(p)^{\otimes 2}$ GME \longrightarrow Condition: $p > p_{\text{GME}}^{(1)}(N) := \frac{2^{N-1}-1}{2^N-1}$ For two copies: $\rho(p)^{\otimes 2} \longrightarrow \mathcal{E}_{\circ}[\rho(p) \otimes \rho(p)]$ If $C_{\mathrm{GM}}(\mathcal{E}_{\circ}[\rho(p)^{\otimes 2}]) > 0 \implies \rho(p)^{\otimes 2}$ GME For k copies: If $C_{\text{GM}}(\mathcal{E}_{\circ}^{\circ(k-1)}[\rho(p)^{\otimes k}]) > 0 \implies \rho(p)^{\otimes k}$ GME Family of k-copy GME thresholds: $p > p_{GME}^{(k)}(N) := \frac{\sqrt[k]{2^{N-1}-1}}{2^{N-1} + \sqrt[k]{2^{N-1}-1}}$





The Conjectures

Our results lead us to two main conjectures

There exists a hierarchy of states with k-copy activatable GME, i.e., for all $k \ge 2$ there exists a biseparable but not partition-separable state ρ such that $\rho^{\otimes k-1}$ is biseparable, but $\rho^{\otimes k}$ is GME.

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GME may be activated for any biseparable but not partition-separable state of any number of parties as $k \to \infty$.

Confirmation of Conjectures

While we have been battling reviewers and the pandemic...

... our conjectures have been proven by our talented colleagues in Madrid

C. Palazuelos and J. I. de Vicente, Quantum **6**, 735 (2022)

They show that:

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Carlos Palazuelos

Julio I. de Vicente

(1) An n-partite state is GME-activatable if and only if it is not partition separable.

(2) For any $k \in \mathbb{N}$ and n (with n > 2), there exists an n-partite GME-activatable state such that k copies of it are not GME.









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