

J_c degradation in CCs for fusion magnets by small defects and its mitigation

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Funded by the European Union

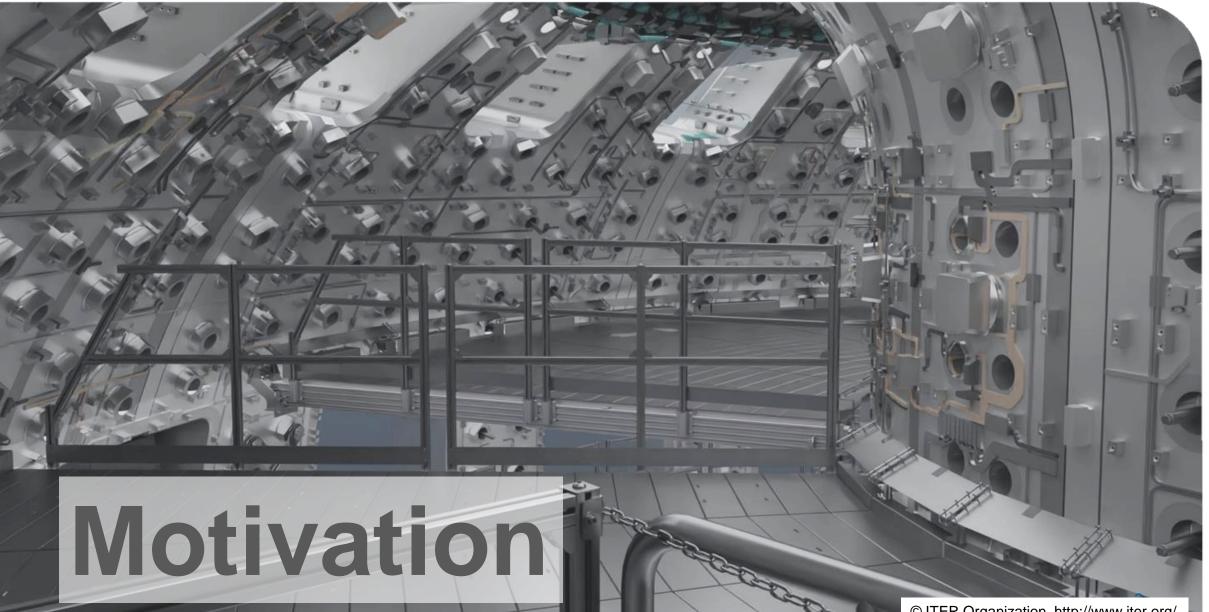






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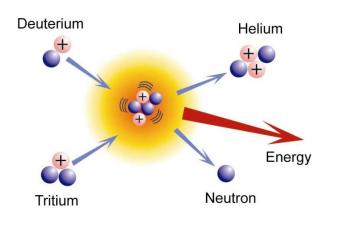




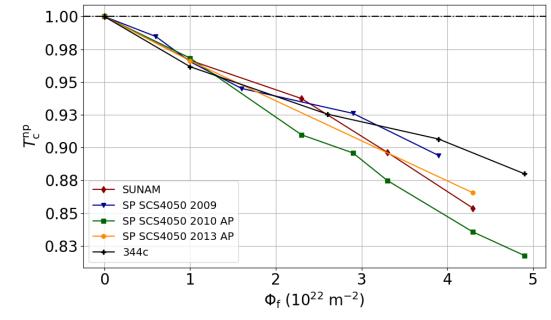
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Nuclear Fusion



Normalized transition temperature



Fast neutron fluence

scattering is pair breaking in d-wave superconductors

- decrease of transition temperature, T_c
- decrease of superfluid density, ρ_s

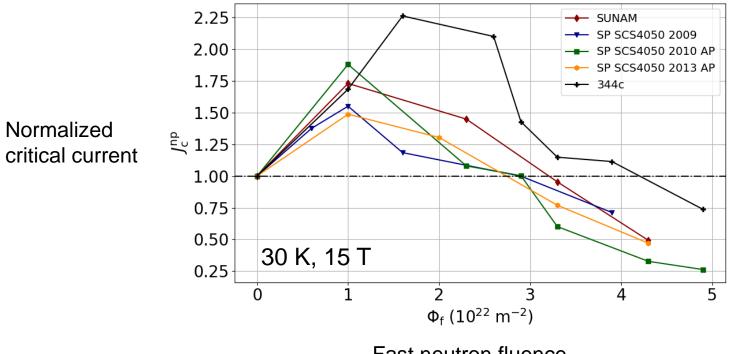
small fraction of the fusion neutrons reach the magnets





introduced defects

- enhance pinning
- increase scattering of charge carriers

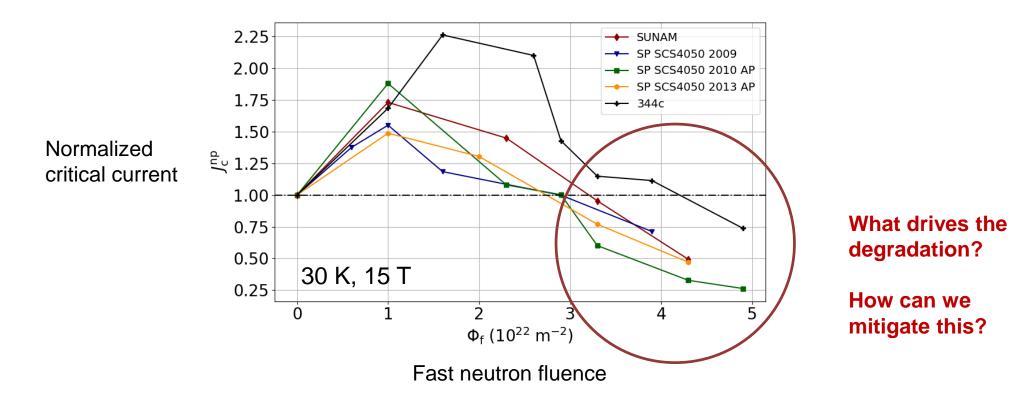


Fast neutron fluence



introduced defects

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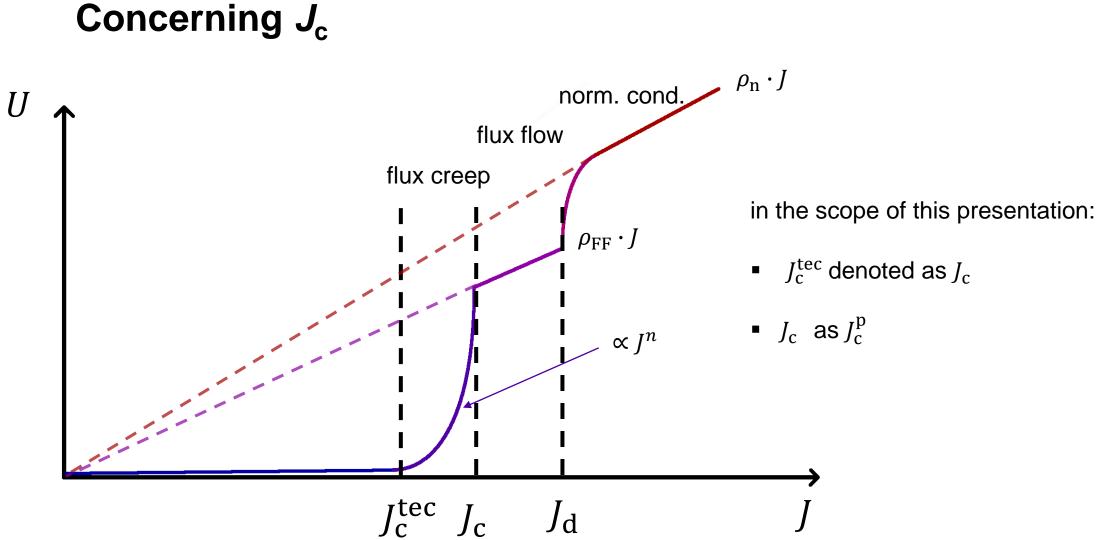




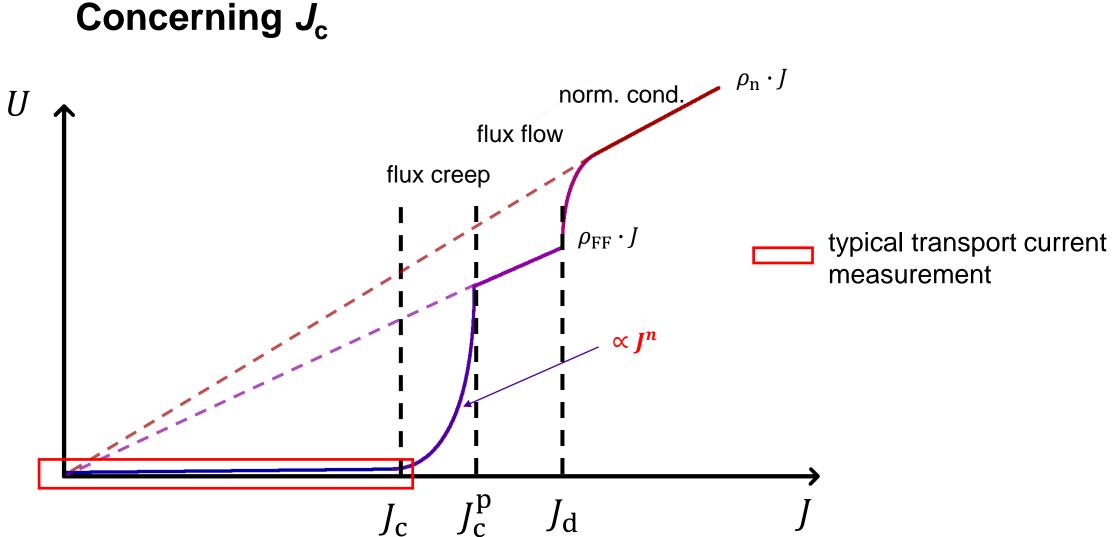
Experimental



Transport current measurements



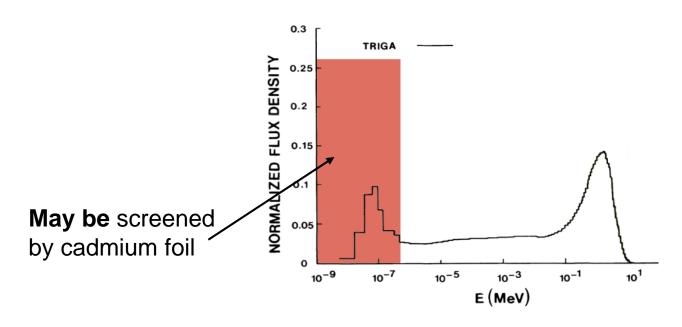
Transport current measurements



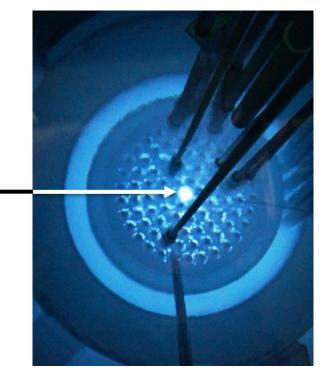


TRIGA MARK II at TU Wien

- irradiation in the central irradiation facility
- fast / thermal neutron flux $3.2 / 4 \times 10^{16} \text{ m}^{-2} \text{ s}^{-1}$
- irradiation with and **without** thermal (< 0.55 eV) neutrons

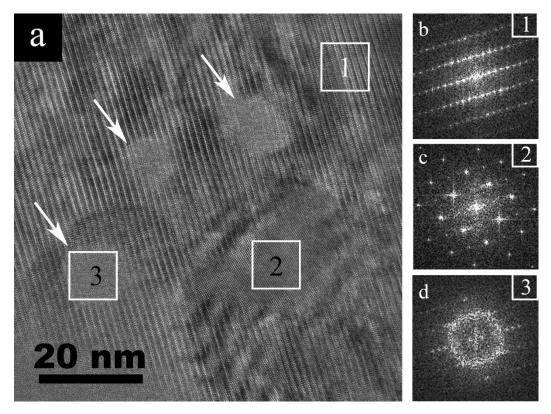


< 70 C at sample

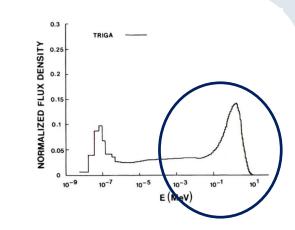


TRIGA MARK II – experimental fission reactor

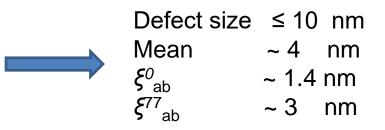
Volume defects



left – TEM picture of neutron induced defects **right** – FFT of selected regions ¹



- 1. Undisturbed GdBCO
- 2. Crystalline BZO rod
- 3. Amorphous cascade



Only large defects visible in TEM

[1] with friendly permission by Yatir Linden, See: Analysing neutron radiation damage in YBa2Cu3O7–x high-temperature superconductor tapes, Department of Materials, University of Oxford, Oxford, UK



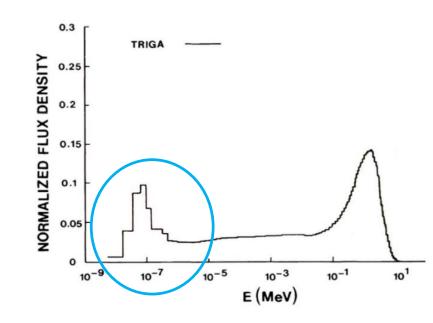


And the rest?

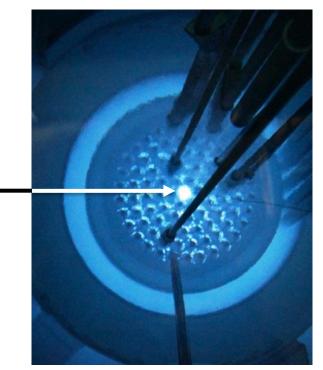


TRIGA MARK II at TU Wien

- irradiation in the central irradiation facility
- fast / thermal neutron flux 3.2 / 4 x 10¹⁶ m⁻² s⁻¹
- irradiation with and without thermal (< 0.55 eV) neutrons

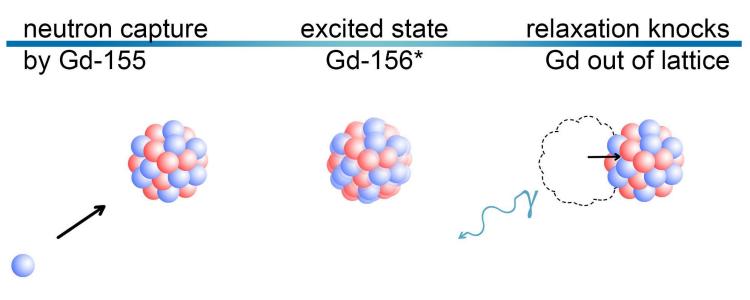


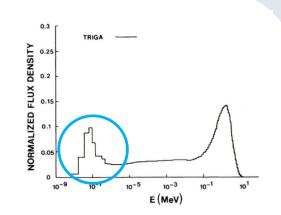
< 70 C at sample



TRIGA MARK II – experimental fission reactor

Point-like defects – Frenkel Pairs

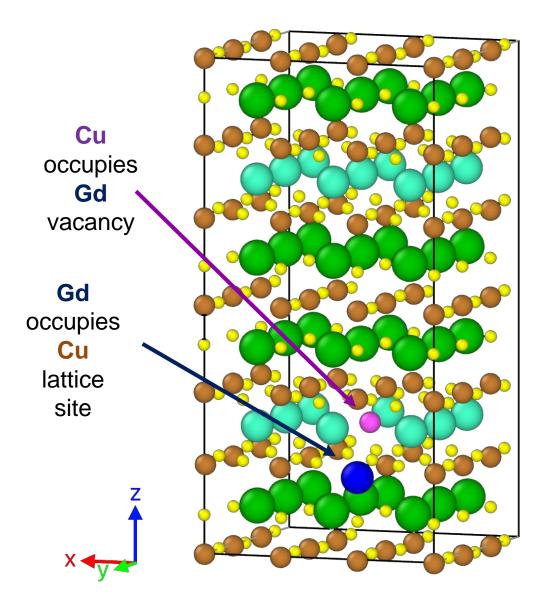


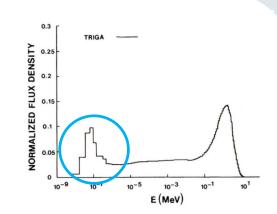


K.E. Sickafus et al., Phys. Rev. B 46 (1992) 11862

- very high defect densities achievable
- add to fast neutron induced defects

Point-like defects – Frenkel Pairs



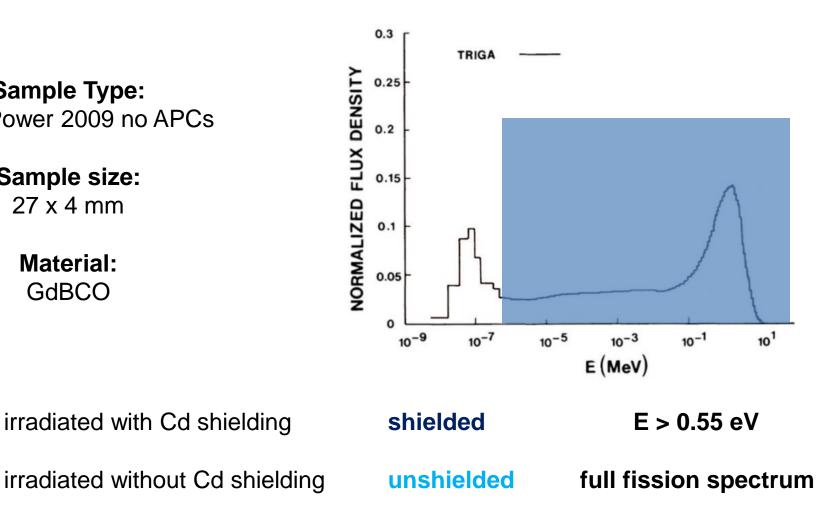


- picture shows one of the potential defects
- distortion in the Oxygen lattice very localized
- very stable distortions of the lattice

*data acquired by MDS using YBCO potential

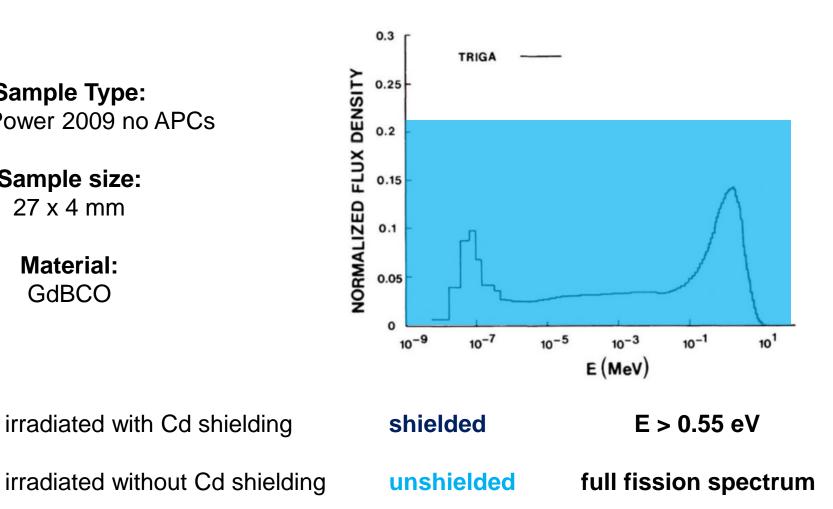
Two samples – two colours

Sample Type: SuperPower 2009 no APCs Sample size: 27 x 4 mm **Material:** GdBCO

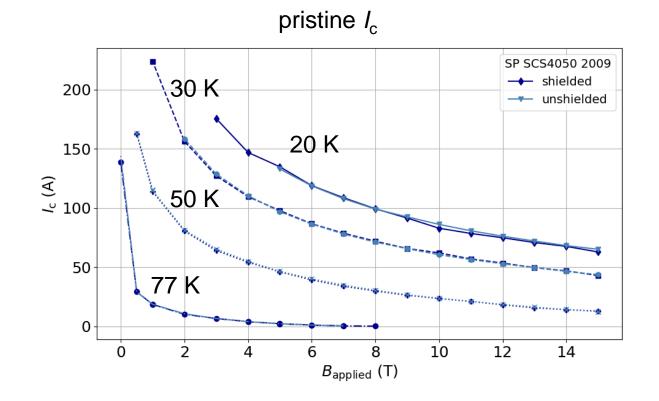


Two samples – two colours

Sample Type: SuperPower 2009 no APCs Sample size: 27 x 4 mm **Material:** GdBCO



Two nearly identical samples



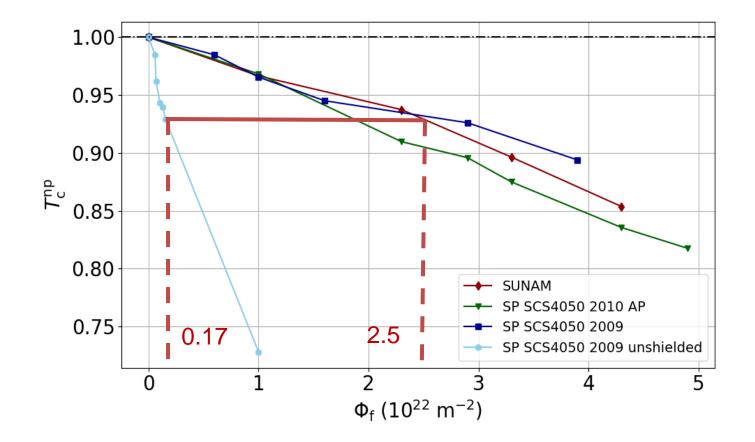
- practically identical $I_c(B,T)$ behavior
- same pristine $T_c = 93 \text{ K}$
- sample homogeneity checked by hall scans at self-field & 77 K







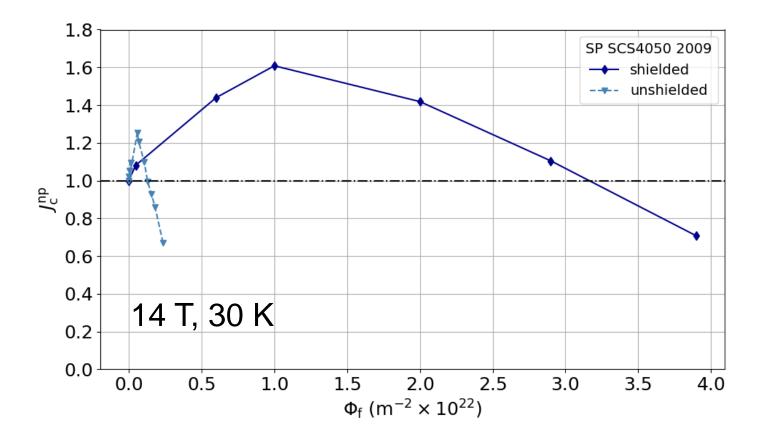
Influence of thermal neutrons: T_c



 $T_{\rm c}$ degrades ~13-15 x faster due to Gd-point defects



\square Influence of thermal neutrons: J_c

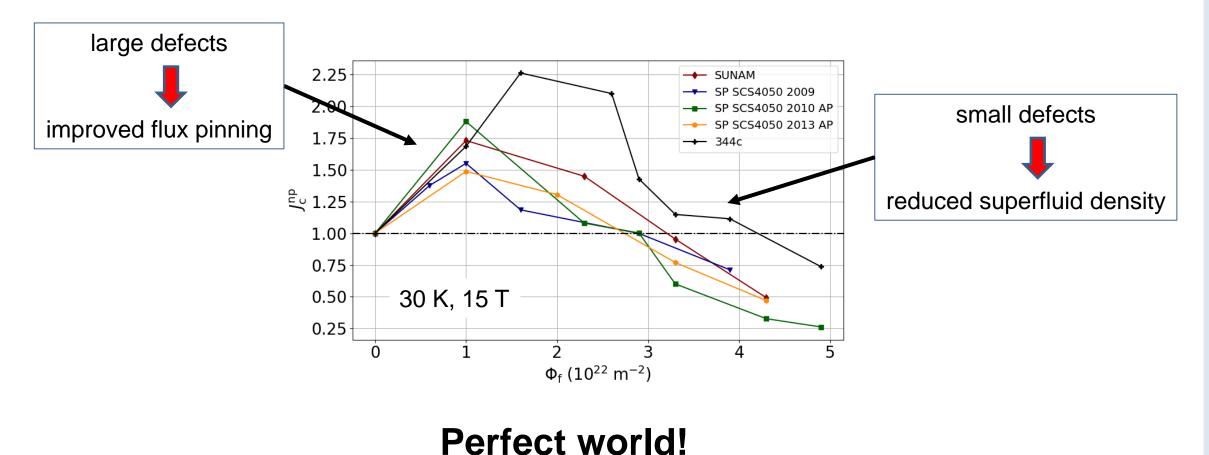


- maximum occurs at much lower neutron fluences
- $J_{\rm c}$ at maximum is smaller
- degradation much steeper

Influence of defect size

high scattering rate - high density of **small** defects

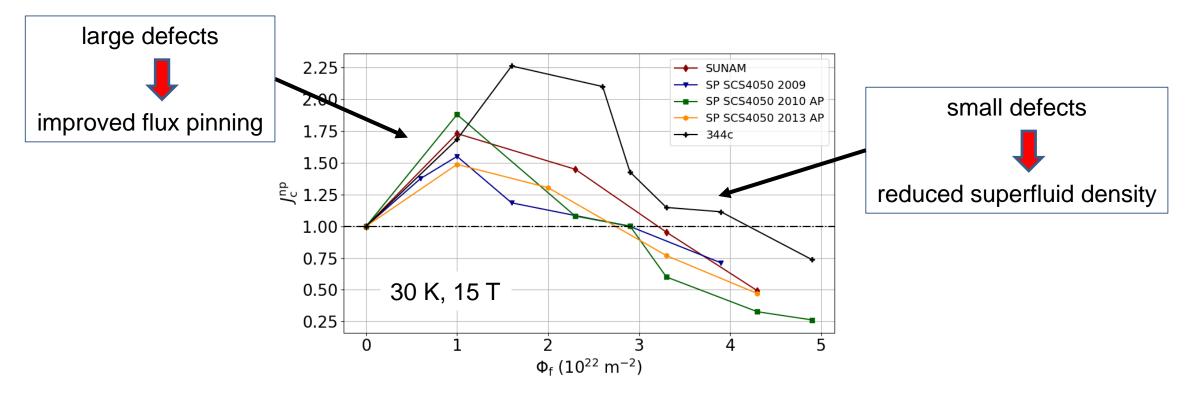
size of pinning centers match the superconducting coherence length: large defects



Influence of defect size

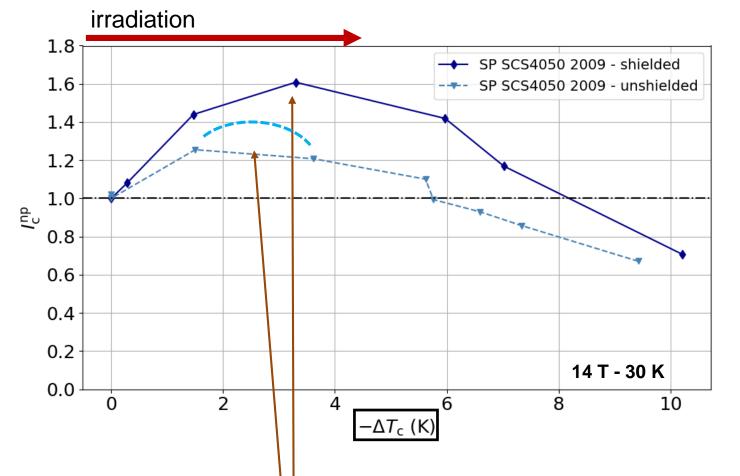
high scattering rate - high density of small defects

size of pinning centers match the superconducting coherence length: large defects



Oversimplified picture?

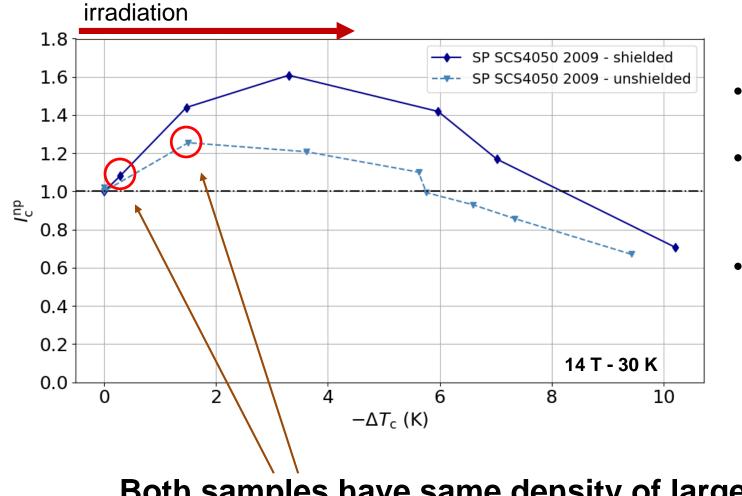
\square Influence of thermal neutrons: J_c



- maximum occurs at similar $T_{\rm c}$
- small defects are efficient pinning centers < 40 K!
- $T_{\rm c}$ is efficient disorder parameter (decrease of superfluid density)

Does J_c increase due to large fast neutron induced defects?

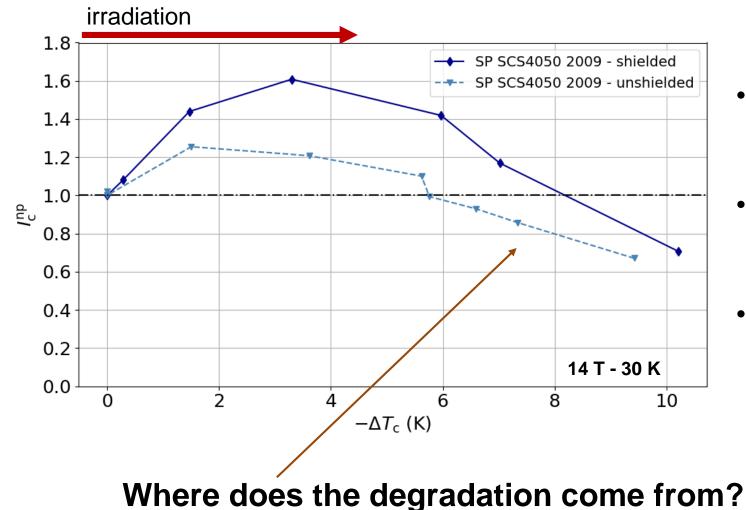
Influence of thermal neutrons: J_{c}



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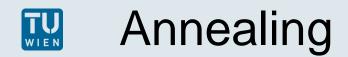
Both samples have same density of large cascades

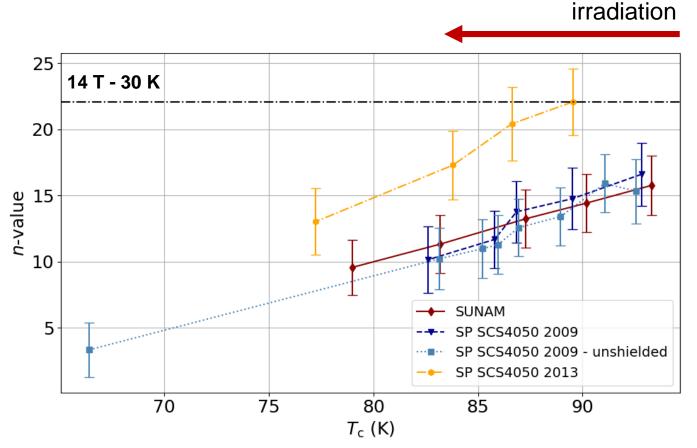
Influence of thermal neutrons: J_c



- extremely different defect size distribution
- almost equivalent slope in degrading branch
- same slope in proton irradiated samples!

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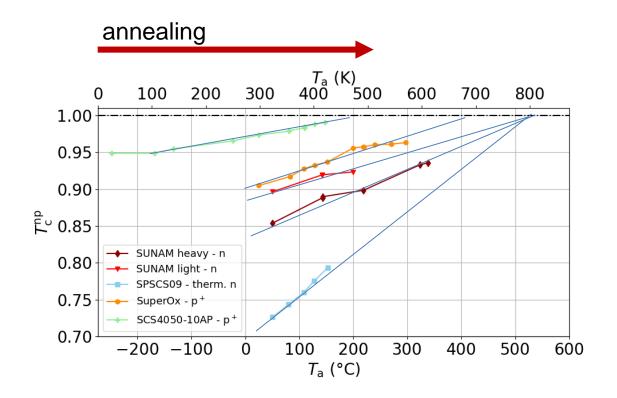
• n - value degrades linearly with T_c

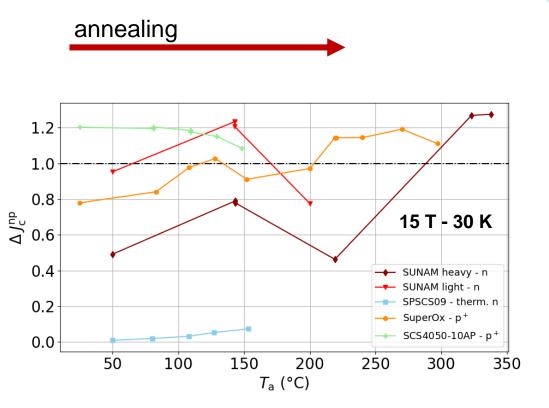
- no change in slope between
 shielded and unshielded sample
- sample with higher starting n
 exhibits steeper slope
- degradation of condensation energy reduces T_c , I_c and n

 $n \propto T_c \propto J_c^p \propto \sigma_s \dots$ superfluid density



Annealing

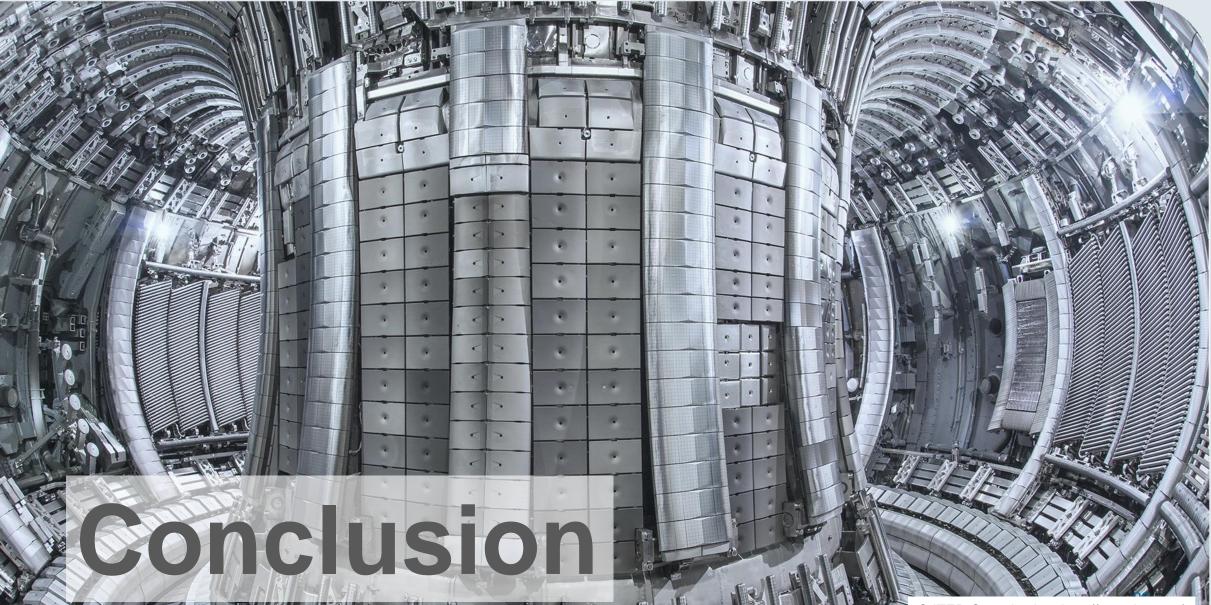




- $T_{\rm c}$ regenerates linearly with $T_{\rm a}$
- even at cryogenic temperatures after reaching onset (p⁺ irradiation)

- $J_{\rm c}$ regenerates non monotonic
- *J*_c only regenerates if maximum has been exceeded





Conclusions

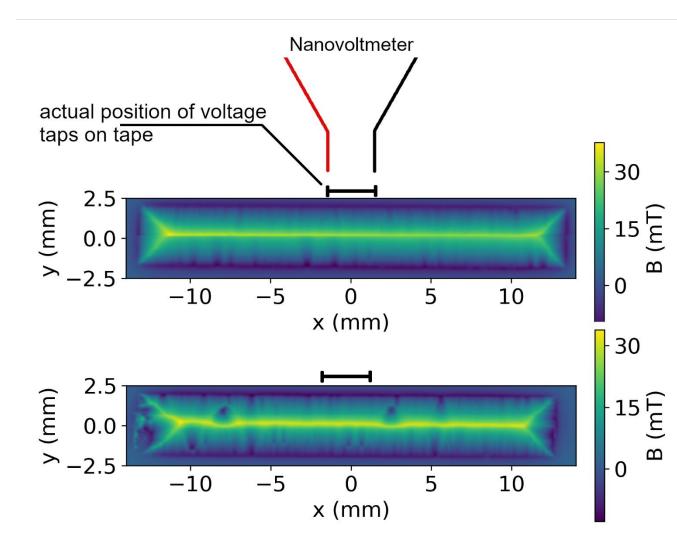
- Pair breaking by scattering decreases T_c linearly with neutron fluence (defect density).
- Point-like disorder / strain field enhances pinning at low temperatures (< 40 K) and high magnetic fields.
- Decrease of J_c at high defect density driven by the decrease of condensation energy.
- Defects, which are responsible for degradation of J_c are stable at elevated temperatures and don't anneal easier.
- Annealing can be an effective way to increase lifetime of magnets but is no simple "cure-all".





Appendix

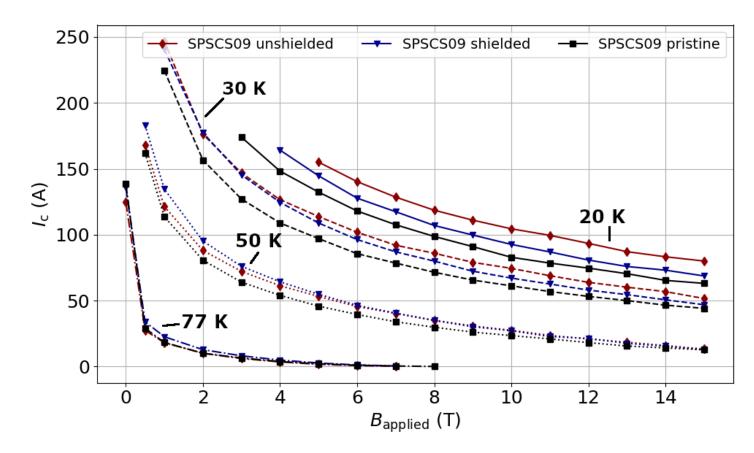
Two nearly identical samples



- sample consistency checked by hall scans
- profile at self-field & 77 K
- voltage taps in low defect areas
- slight differences in signal due to probe – sample distance



Influence of Gd point defects: J_c

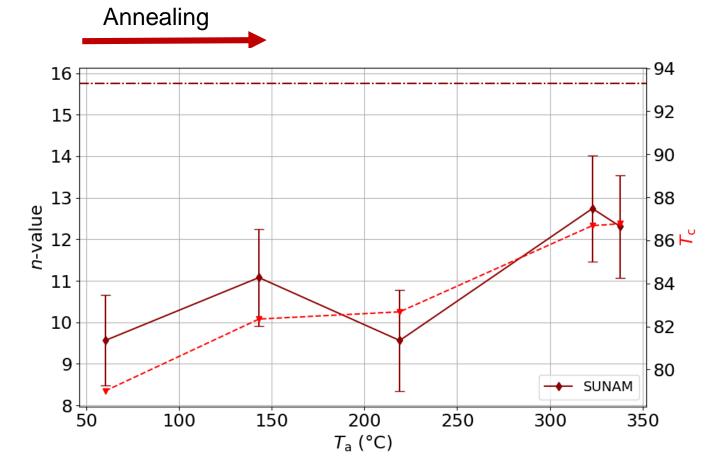


irradiated to similar fluence

unshielded sample:

- smaller T_c
- similar density of large defects
- larger J_c at low temperatures and high fields
- lower J_c at low fields (crossover)

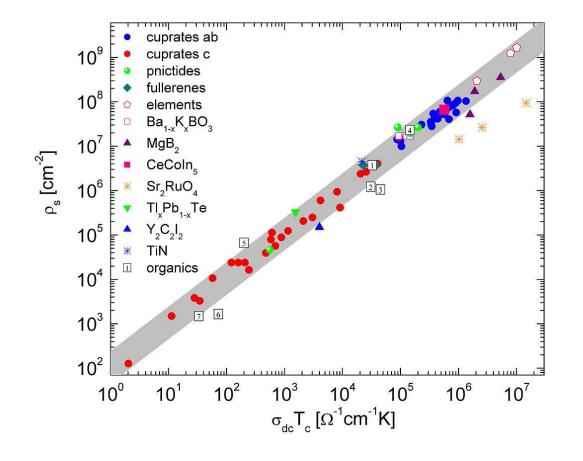
Displaced Gd atoms lead to efficient pinning below about 30 K!



- $E_{\rm c}$ increases ~linearly with $T_{\rm a}$
- J_c^{p} increases monotonically with E_c
- n-value however does not and leads to non monotonic J_c annealing

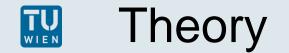


Homes' scaling law



- though logarithmic, the superfluid ٠ density scales with σ_{dc} and T_{c}
- many orders of magnitude
- many different materials





$$E_{\rm c} = \frac{1}{\lambda^2 \xi^2} = \frac{1}{\lambda^2} H_{\rm c2}$$

$$\frac{1}{\lambda^2} = \rho_s \qquad \xi = \sqrt{\xi_0 l}$$

$$\xi_0 \propto k_b T_{\rm c} \qquad \rho_{\rm DC} \propto \frac{1}{l}$$

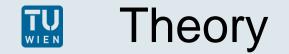
$$E_{\rm c} = \rho_s \frac{1}{\xi_0 l} \propto \rho_s \rho_{DC} T_c$$

 $\rho_{\rm s} \dots$ superfluid density

 $\sigma_{dc} = \rho_{dc}^{-1}$... normal state conductivity ξ_0 ... clean limit coherence length

l ... mean free path

 $H_{c_2} \propto \rho_{DC} T_c$



$$n \propto U_0$$
$$U_0 \propto E_c$$
$$E_c = \rho_s \frac{1}{\xi_0 l} \propto \rho_s \rho_{DC} T_c \propto n$$

 $\rho_{\rm s}$... superfluid density

 $\sigma_{dc} = \rho_{dc}^{-1}$... normal state conductivity ξ_0 ... clean limit coherence length l ... mean free path E_c ... condensation energy U_0 ... pinning energy