



CANNEX

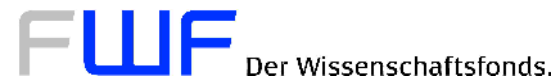


Casimir And Non-Newtonian force EXperiment

# Parallel plate force metrology: Status and Perspectives

PIERS 2023  
July 03, 2023, Prague

René Sedmik



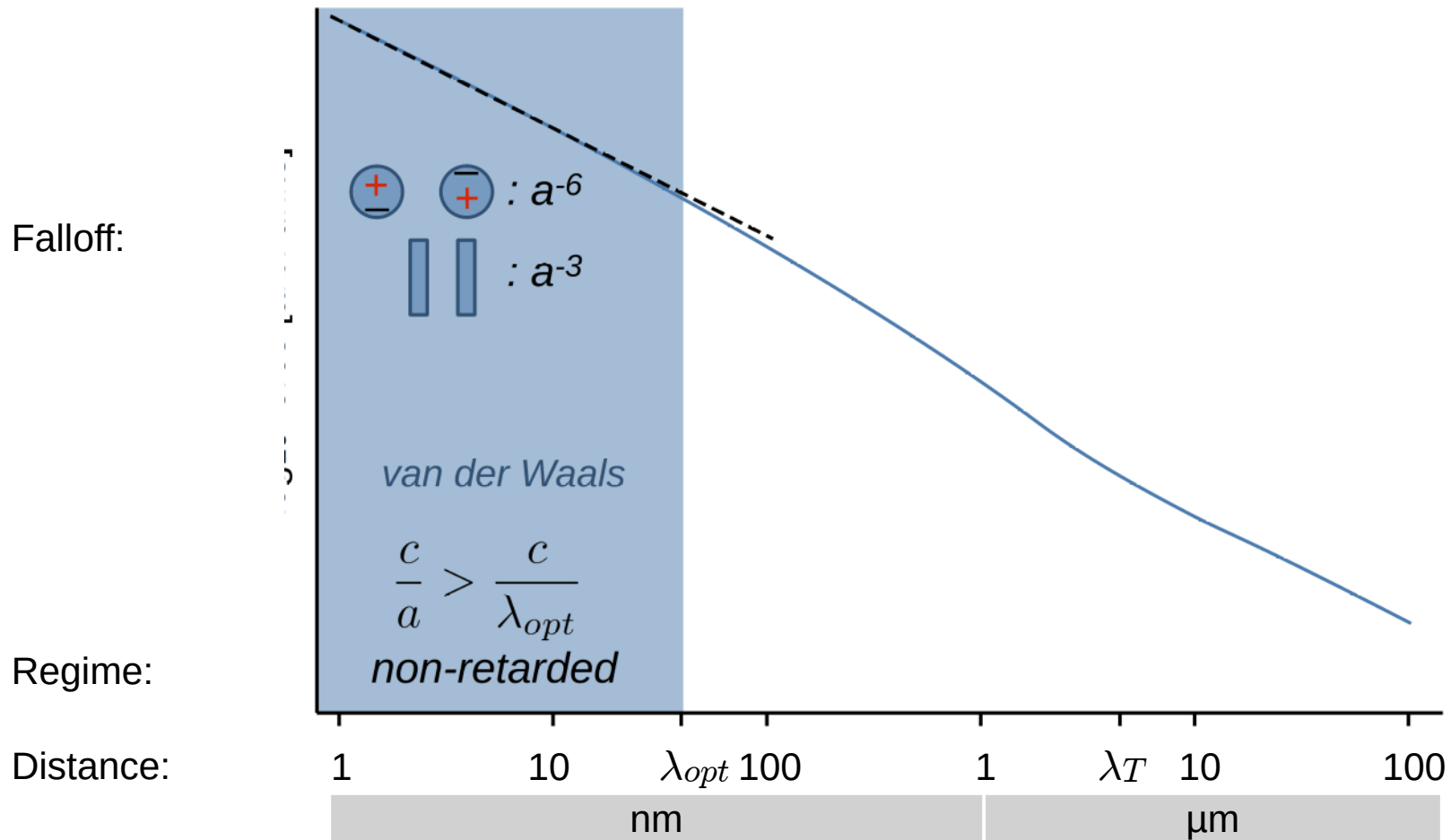
A little repetition

+

Thermal effects

# Morphology of dipole interactions:

Dominant source: vacuum fluctuations:



Distance  $a$

Thermal wavelength:  $\lambda_T = \frac{\hbar c}{k_B T}$

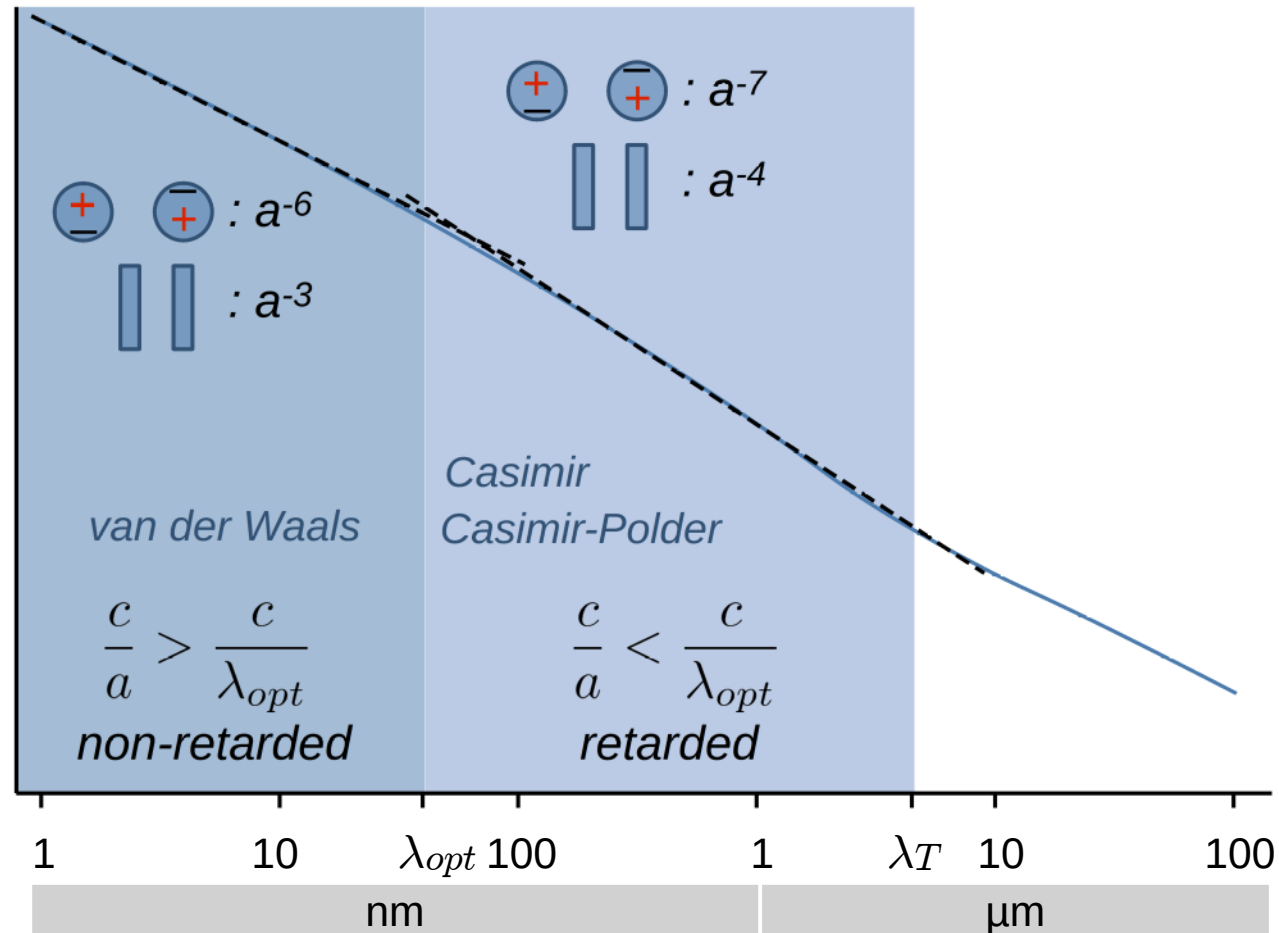
# Morphology of dipole interactions:

Dominant source:

vacuum  
fluctuations:

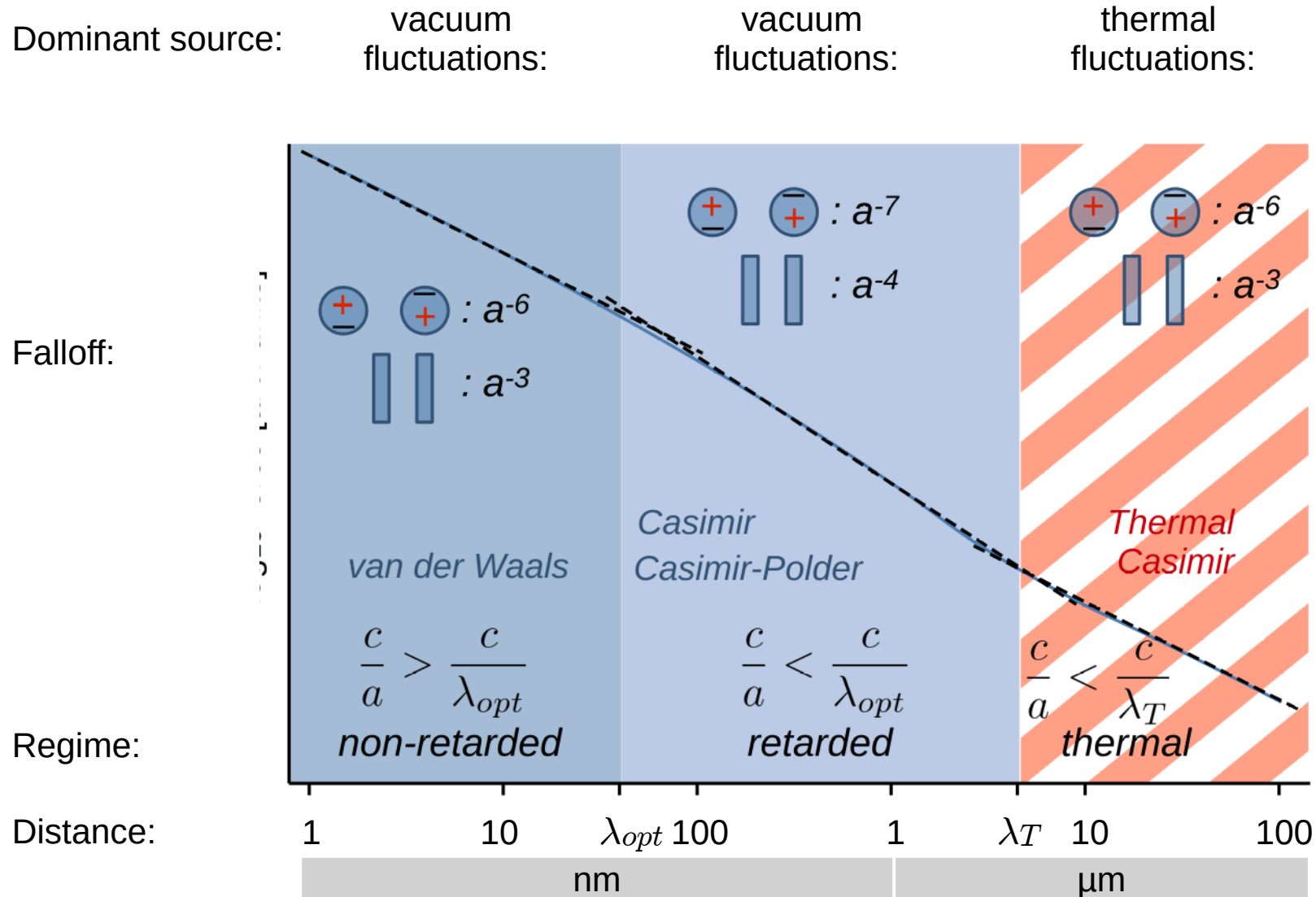
vacuum  
fluctuations:

Falloff:



Thermal wavelength:  $\lambda_T = \frac{\hbar c}{k_B T}$

# Morphology of dipole interactions:

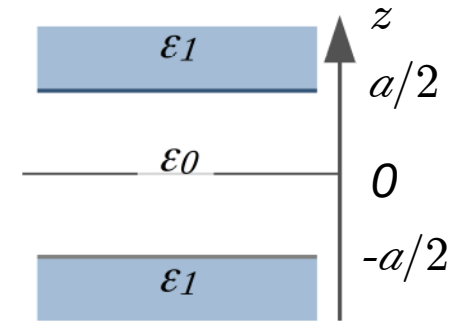


Thermal wavelength:  $\lambda_T = \frac{\hbar c}{k_B T}$

# Starting point: Lifshitz theory

Energy per area between plane parallel plates,  $T = 0$ :

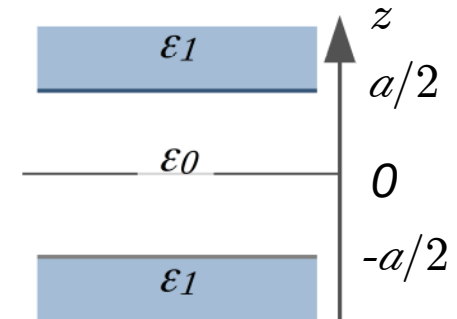
$$\frac{E_{\text{ren}}(a)}{A} = \frac{\hbar}{(2\pi)^2 c^2} \int_1^\infty dp \int_0^\infty d\xi p \xi^2 \left[ \ln \frac{\Delta_\perp(i\xi, a)}{\Delta_{\perp, \infty}(i\xi)} + \ln \frac{\Delta_\parallel(i\xi, a)}{\Delta_{\parallel, \infty}(i\xi)} \right],$$



$$\frac{\Delta_\perp(p, i\xi, a)}{\Delta_{\perp, \infty}(p, i\xi)} = 1 - \left( \frac{K_1 \epsilon_0(i\xi) - K_0 \epsilon_1(i\xi)}{K_1 \epsilon_0(i\xi) + K_0 \epsilon_1(i\xi)} \right)^2 e^{-2a \frac{\xi}{c} K_0}, \quad K_j(p, i\xi) = \sqrt{p^2 - 1 + \epsilon_j(i\xi)}$$

# Starting point: Lifshitz theory

Energy per area between plane parallel plates,  $T = 0$ :



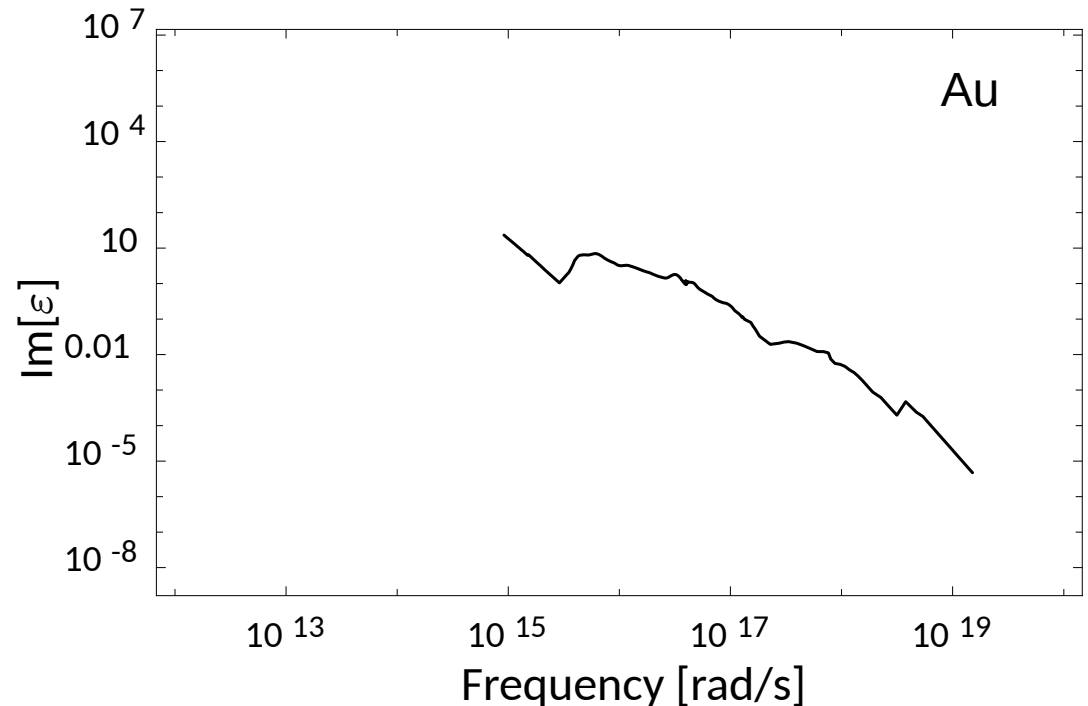
$$\frac{E_{\text{ren}}(a)}{A} = \frac{\hbar}{(2\pi)^2 c^2} \int_1^\infty dp \int_0^\infty d\xi p \xi^2 \left[ \ln \frac{\Delta_\perp(i\xi, a)}{\Delta_{\perp, \infty}(i\xi)} + \ln \frac{\Delta_\parallel(i\xi, a)}{\Delta_{\parallel, \infty}(i\xi)} \right],$$

$$\frac{\Delta_\perp(p, i\xi, a)}{\Delta_{\perp, \infty}(p, i\xi)} = 1 - \left( \frac{K_1 \epsilon_0(i\xi) - K_0 \epsilon_1(i\xi)}{K_1 \epsilon_0(i\xi) + K_0 \epsilon_1(i\xi)} \right)^2 e^{-2a \frac{\xi}{c} K_0}, \quad K_j(p, i\xi) = \sqrt{p^2 - 1 + \epsilon_j(i\xi)}$$

Kramers Kronig transformation:

$$\epsilon(i\xi) = 1 + \frac{2}{\pi} \int_0^\infty d\omega \frac{\omega \text{Im}\epsilon(\omega)}{\omega^2 + \xi^2}$$

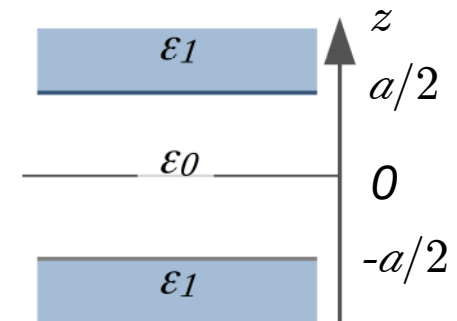
Ellipsometric data:  $\text{Im}(\epsilon) = 2nk$



# Starting point: Lifshitz theory

Energy per area between plane parallel plates,  $T = 0$ :

$$\frac{E_{\text{ren}}(a)}{A} = \frac{\hbar}{(2\pi)^2 c^2} \int_1^\infty dp \int_0^\infty d\xi p \xi^2 \left[ \ln \frac{\Delta_\perp(i\xi, a)}{\Delta_{\perp, \infty}(i\xi)} + \ln \frac{\Delta_\parallel(i\xi, a)}{\Delta_{\parallel, \infty}(i\xi)} \right],$$



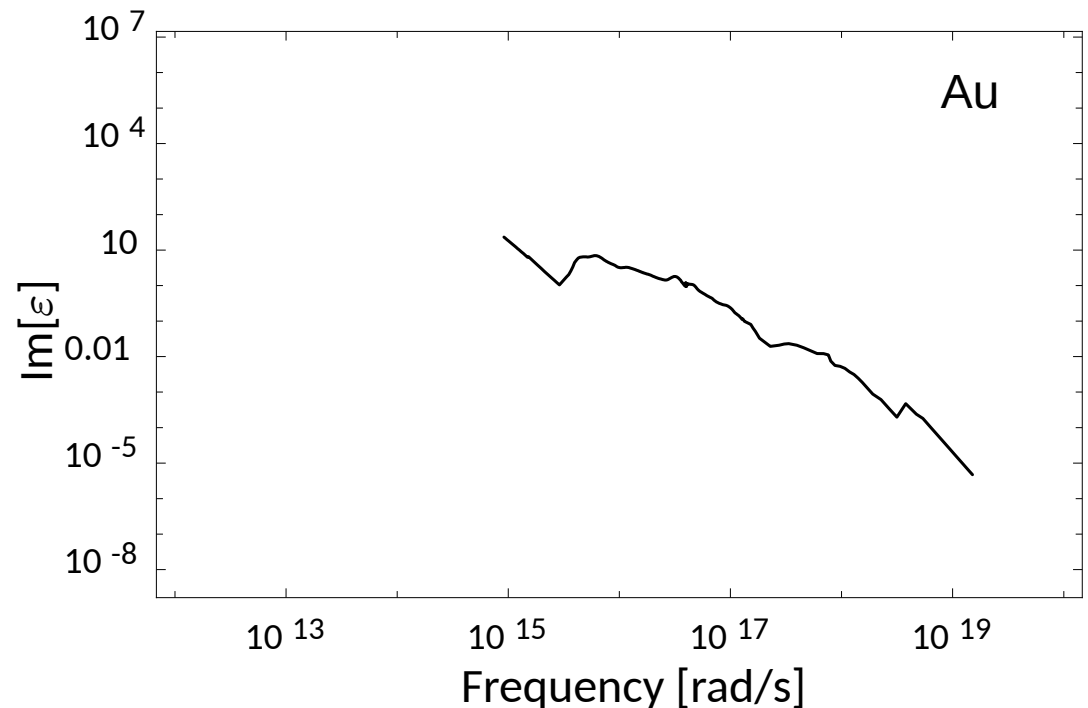
$$\frac{\Delta_\perp(p, i\xi, a)}{\Delta_{\perp, \infty}(p, i\xi)} = 1 - \left( \frac{K_1 \varepsilon_0(i\xi) - K_0 \varepsilon_1(i\xi)}{K_1 \varepsilon_0(i\xi) + K_0 \varepsilon_1(i\xi)} \right)^2 e^{-2a \frac{\xi}{c} K_0}, \quad K_j(p, i\xi) = \sqrt{p^2 - 1 + \varepsilon_j(i\xi)}$$

Kramers Kronig transformation:

$$\varepsilon(i\xi) = 1 + \frac{2}{\pi} \int_0^\infty d\omega \frac{\omega \text{Im}\varepsilon(\omega)}{\omega^2 + \xi^2}$$

Ellipsometric data:  $\text{Im}(\varepsilon) = 2nk$

**Need to extrapolate  
to zero frequency**





# Starting point: The Drude/Plasma debate (Casimir puzzle)

Dissipation at zero frequency or not?

Drude

$$\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}$$

plasma

$$\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2}$$

+ oscillator terms

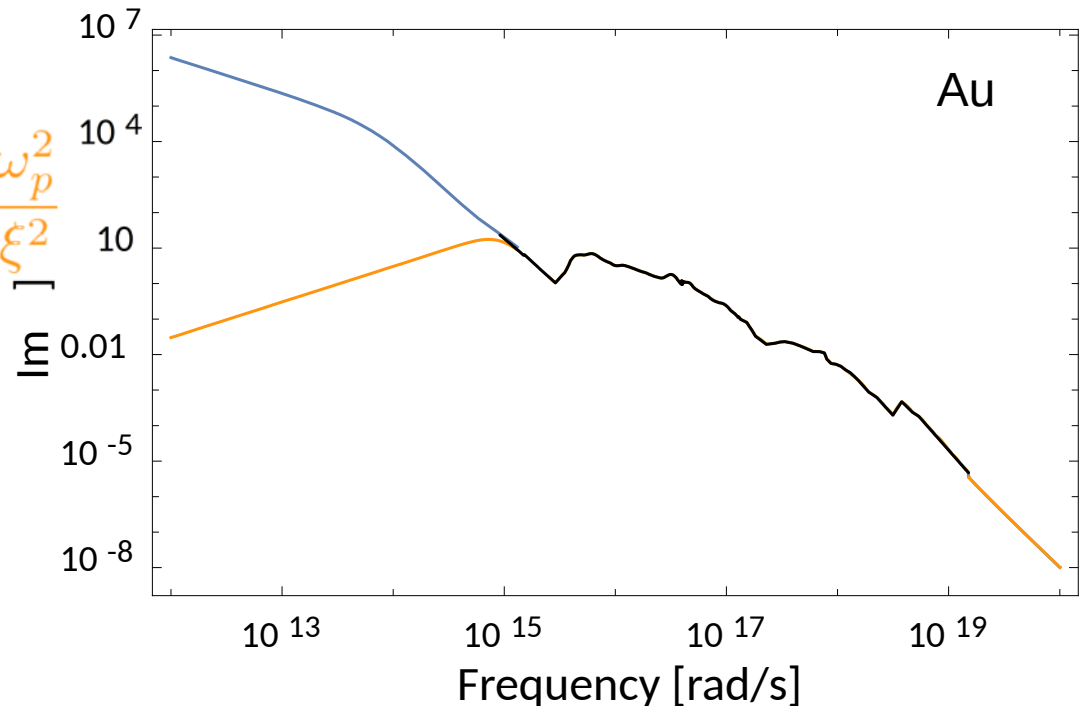
$\omega_p$  plasma frequency  
 $\gamma$  relaxation frequency (dissipation)

Kramers Kronig transformation:

$$\epsilon(i\xi) = 1 + \frac{2}{\pi} \int_0^{\infty} d\omega \frac{\omega \text{Im}\epsilon(\omega)}{\omega^2 + \xi^2} + \frac{\omega_p^2}{\xi^2}$$

Ellipsometric data:  $\text{Im}(\epsilon) = 2nk$

**Need to extrapolate to zero frequency**



# Starting point: The Drude/Plasma debate (Casimir puzzle)

Dissipation at zero frequency or not?

$$\text{Drude} \\ \epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}$$

$$\text{plasma} \\ \epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2}$$

+ oscillator terms

$\omega_p$  plasma frequency  
 $\gamma$  relaxation frequency (dissipation)

Physics?

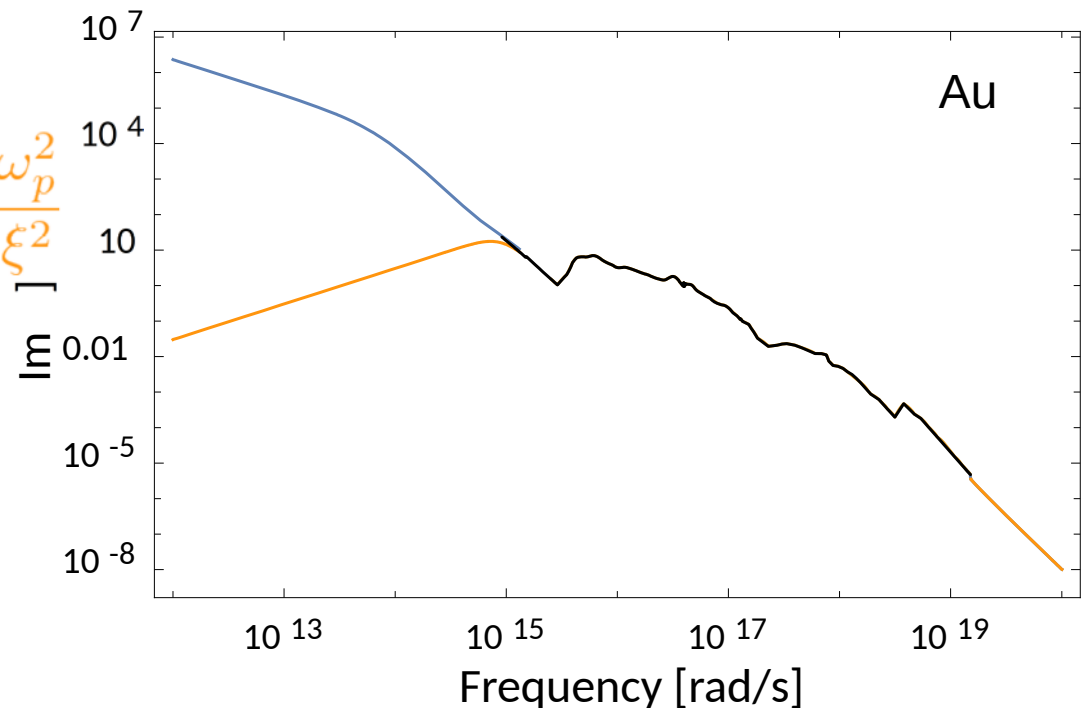
- Neither of the two models are derived from first principles.
- None of these models can be a description for a real material

Kramers Kronig transformation:

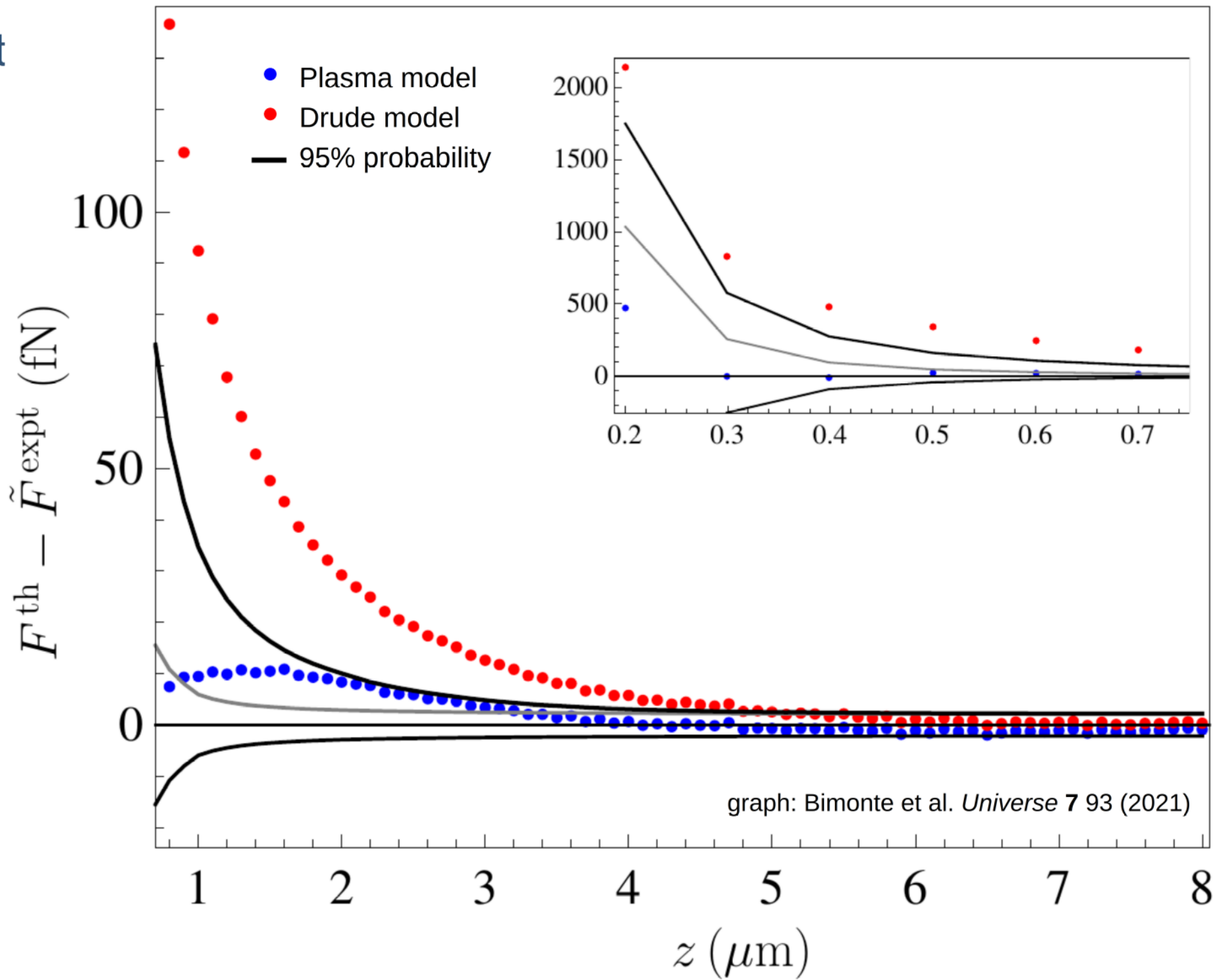
$$\epsilon(i\xi) = 1 + \frac{2}{\pi} \int_0^{\infty} d\omega \frac{\omega \text{Im}\epsilon(\omega)}{\omega^2 + \xi^2} + \frac{\omega_p^2}{\xi^2}$$

Ellipsometric data:  $\text{Im}(\epsilon) = 2nk$

Need to extrapolate to zero frequency



St



# Starting point: The Drude/Plasma debate (Casimir puzzle)

## Dissipation at zero frequency or not? Drude vs. plasma debate (still)

<p>Drude</p> $\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}$	<p>plasma</p> $\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2}$	+ oscillator terms	$\omega_p$	plasma frequency
			$\gamma$	relaxation frequency (dissipation)

- Recent data suggests:**
- No dissipation for vacuum fluctuations (short distance)!  
Bimonte *et al.*, Phys. Rev. B **93**, 184434 (2016)  
Banishev *et al.*, PRL **110**, 137401 (2013)  
Bimonte *et al.*, Universe **7** 93 (2021)  
Liu *et al.*, Phys. Rev. A **100**, 052511(2019)
  - Situation at large distance ( $\gtrsim 5 \mu\text{m}$ ) unclear

## Physics?

- Neither of the two models are derived from first principles.
- None of these models can be a description for a real material

# Starting point: The Drude/Plasma debate (Casimir puzzle)

## Dissipation at zero frequency or not? Drude vs. plasma debate (still)

<p style="text-align: center;">Drude</p> $\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}$	<p style="text-align: center;">plasma</p> $\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2}$	+ oscillator terms	$\omega_p$ plasma frequency $\gamma$ relaxation frequency (dissipation)
--	---	--------------------	--

- Recent data suggests:**
- No dissipation for vacuum fluctuations (short distance)!  
Bimonte *et al.*, Phys. Rev. B **93**, 184434 (2016)  
Banishev *et al.*, PRL **110**, 137401 (2013)  
Bimonte *et al.*, Universe **7** 93 (2021)  
Liu *et al.*, Phys. Rev. A **100**, 052511(2019)
  - Situation at large distance ( $\gtrsim 5 \mu\text{m}$ ) unclear

## Physics?

- Neither of the two models are derived from first principles.
- None of these models can be a description for a real material

## But ...

- (DC) Dissipation is a proven concept (e.g. Johnson noise)
- ... captured by the Drude model

**Q: Why is this still a topic worth investigating?**

# Starting point: The Drude/Plasma debate (Casimir puzzle)

## Dissipation at zero frequency or not? Drude vs. plasma debate (still)

<p style="text-align: center;">Drude</p> $\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega}$	<p style="text-align: center;">plasma</p> $\epsilon_r(\omega) = \epsilon_0 - \frac{\omega_p^2}{\omega^2}$	+ oscillator terms	$\omega_p$ plasma frequency $\gamma$ relaxation frequency (dissipation)
--	---	--------------------	--

- Recent data suggests:**
- No dissipation for vacuum fluctuations (short distance)!  
Bimonte *et al.*, Phys. Rev. B **93**, 184434 (2016)  
Banishev *et al.*, PRL **110**, 137401 (2013)  
Bimonte *et al.*, Universe **7** 93 (2021)  
Liu *et al.*, Phys. Rev. A **100**, 052511(2019)
  - Situation at large distance ( $\gtrsim 5 \mu\text{m}$ ) unclear

## Physics?

- Neither of the two models are derived from first principles.
- None of these models can be a description for a real material

## But ...

- (DC) Dissipation is a proven concept (e.g. Johnson noise)
- ... captured by the Drude model

## Q: Why is this still a topic worth investigating?

**A1: Maybe a test of the fluctuation dissipation theorem**

**A2: Role of Non-locality**

**A3: Maybe different responses for propagating and evanescent waves?**

# Thermal Casimir forces (in thermal equilibrium).

Include the spectrum of thermal Planck photons  $T \neq 0$

$$\mathcal{F} = -k_B T \ln \sum_{n=0}^{\infty} e^{-\hbar\omega_n/k_B T} \quad \text{free energy}$$

**Formal replacement:**

Matsubara frequencies

$$\frac{\hbar}{2\pi} \int_0^{\infty} d\xi \leftrightarrow k_B T \sum_{\ell=0}^{\infty} \quad \xi_{\ell} = 2\pi \frac{k_B T}{\hbar} \ell$$

Casimir energy at non-zero temperature between parallel plates

$$\frac{\mathcal{F}_{\text{ren}}(a, T)}{A} = \frac{k_B T}{(2\pi)c^2} \int_1^{\infty} dp \sum_{\ell=0}^{\infty} p \xi_{\ell}^2 \left[ \ln \frac{\Delta_{\perp}(i\xi_{\ell}, a)}{\Delta_{\perp, \infty}(i\xi_{\ell})} + \ln \frac{\Delta_{\parallel}(i\xi_{\ell}, a)}{\Delta_{\parallel, \infty}(i\xi_{\ell})} \right],$$

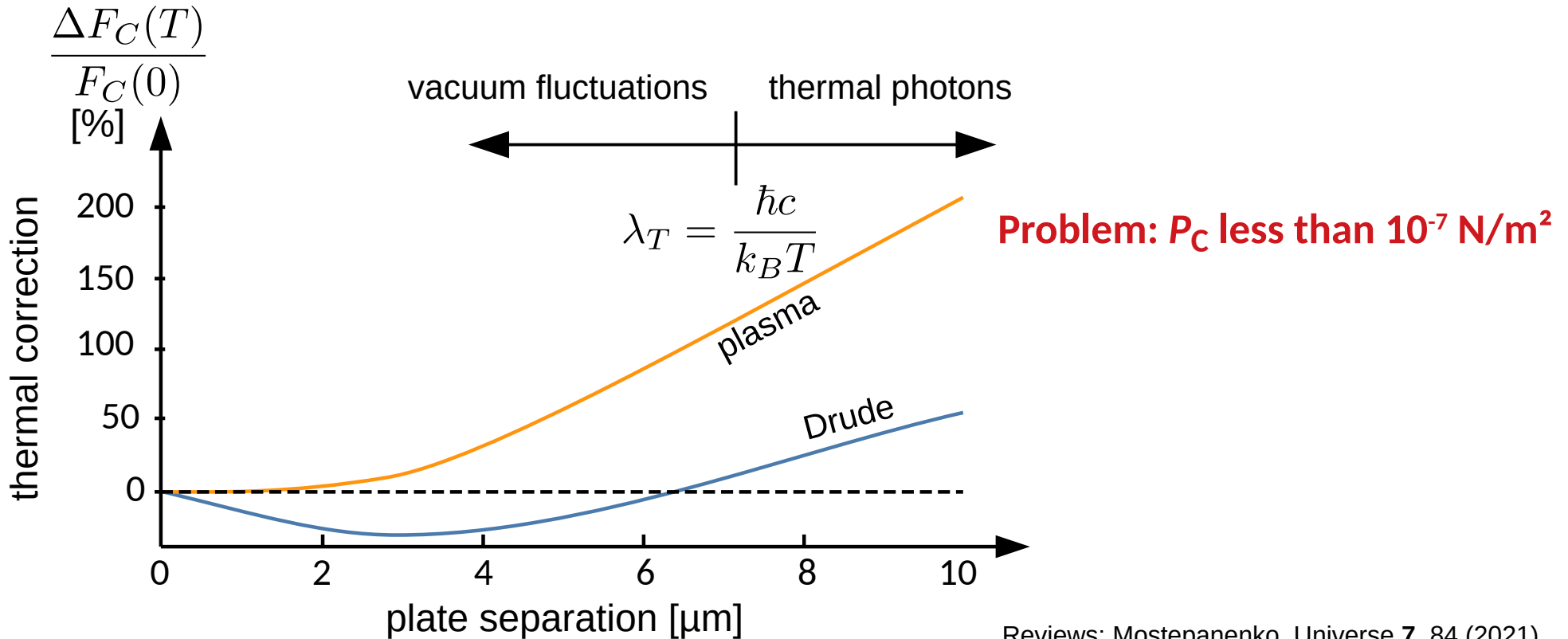
**Split into vacuum and thermal contributions:**

$$\mathcal{F}(a, T) = \underbrace{E_{\text{ren}}(a, T = 0)}_{\text{vacuum}} + \underbrace{\Delta_T \mathcal{F}(a, T)}_{\text{thermal fluctuations}}$$

# Thermal Casimir forces and the Casimir Puzzle. Some ideas on what to measure?

## Idea 1 (the obvious one)

Measure the force (pressure) at a  $\sim 10 \mu\text{m}$ :



Reviews: Mostepanenko, Universe 7, 84 (2021)

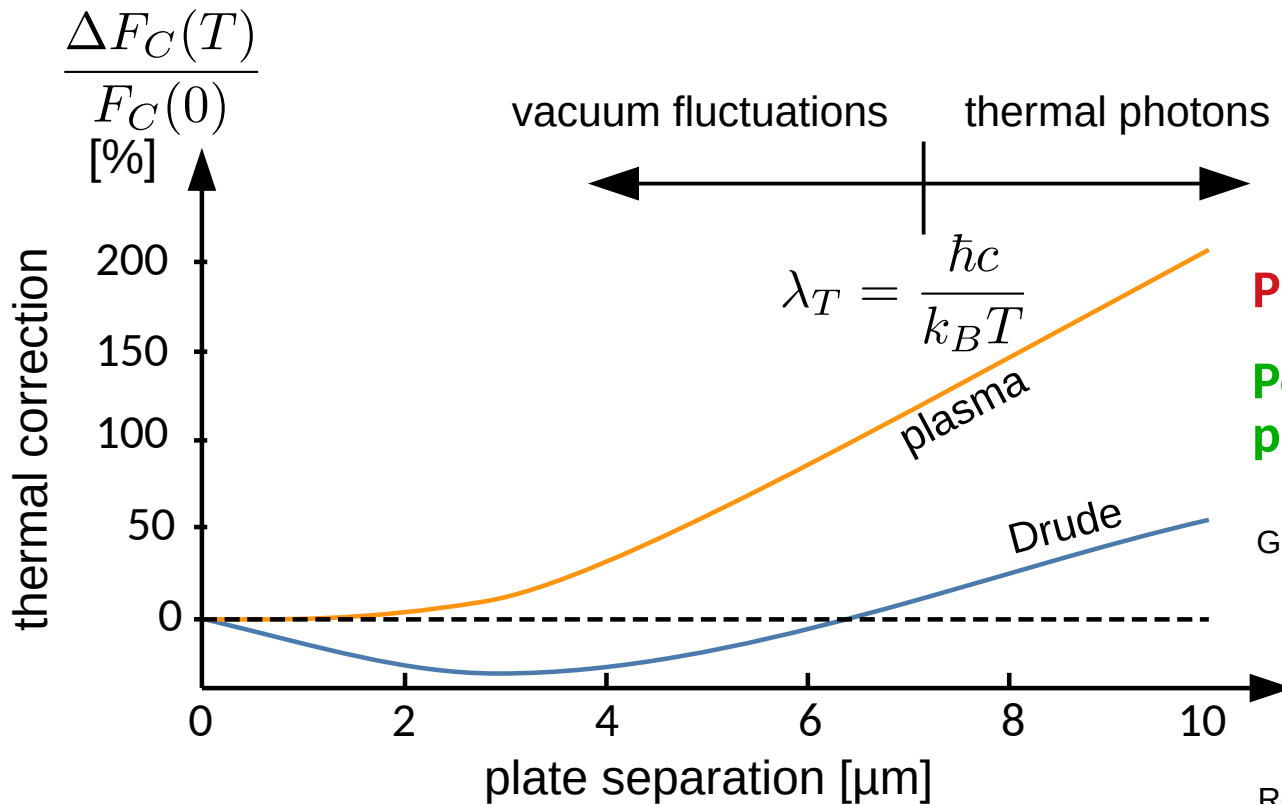
M.. Bordag et al "Advances in the Casimir effect",  
Oxford Science Publications, (2015).



# Thermal Casimir forces and the Casimir Puzzle. Some ideas on what to measure?

## Idea 1 (the obvious one)

Measure the force (pressure) at a  $\sim 10 \mu\text{m}$ :



**Problem:  $P_C$  less than  $10^{-7} \text{ N/m}^2$**

**Possible to measure with parallel plates.**

G. Klimchitskaya et al, *Symmetry* **11**, 407 (2019)

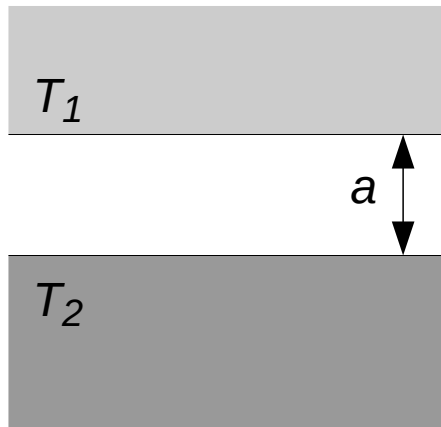
Reviews: Mostepanenko, *Universe* **7**, 84 (2021)

M.. Bordag et al "Advances in the Casimir effect", *Oxford Science Publications*, (2015).

# Casimir forces out of thermal equilibrium

## Idea 2

Two half spaces at **different** temperature



**Half spaces:**

$$P(a, T_1, T_2) = \frac{1}{2} [P_{\text{eq}}(a, T_1) + P_{\text{eq}}(a, T_2)]$$

equilibrium pressures

**pioneering works:**

Antezza *et al*, PRL 95, 113202 (2005) , PRA **77**, 022901 (2008),  
PRL 97, 223203 (2006)

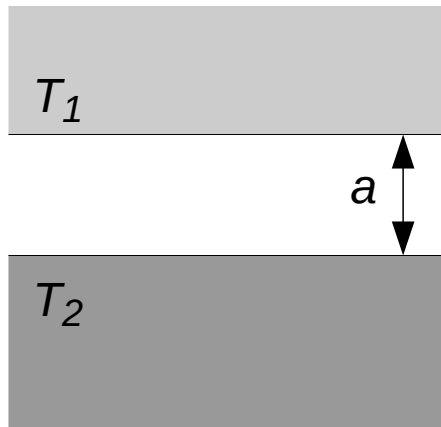
Krüger *et al*. EPL 95, 21002 (2011)

Bimonte *et al*, PRA **84**, 042503 (2011), PRA **92**, 032116 (2015)

# Casimir forces out of thermal equilibrium

## Idea 2

Two half spaces at **different** temperature



**Half spaces:**

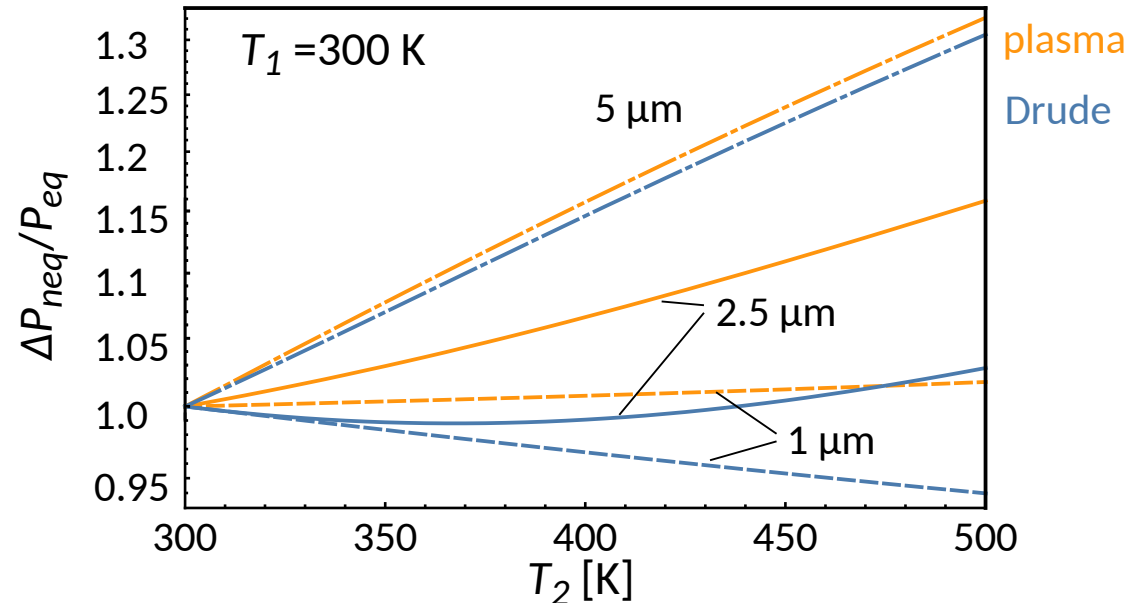
$$P(a, T_1, T_2) =$$

$$\frac{1}{2} [P_{\text{eq}}(a, T_1) + P_{\text{eq}}(a, T_2)] + \Delta P_{\text{neq}}(a, T_1, T_2) + \frac{2\sigma}{3c} (T_2^4 + T_1^4)$$

equilibrium pressures
anti-symmetric term
radiation pressure

**pioneering works:**

- Antezza *et al*, PRL 95, 113202 (2005) , PRA **77**, 022901 (2008), PRL 97, 223203 (2006)
- Krüger *et al*. EPL 95, 21002 (2011)
- Bimonte *et al*, PRA **84**, 042503 (2011), PRA **92**, 032116 (2015)



**graph data:** G. Klimchitskaya, V.M. Mostepanenko, and R.I.P. Sedmik, PRA **100**, 022511 (2019)

identical plates:  $\Delta P_{\text{neq}} \rightarrow 0$

Dorofeyev, J. Phys. A **31**, 4369 (1998)

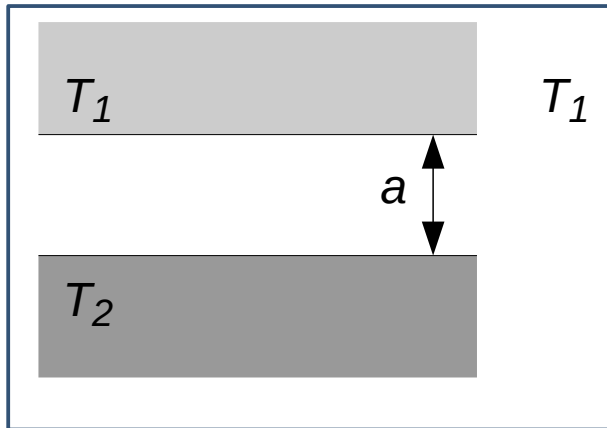
Antezza *et al*, PRA **77**, 022901 (2008)

# Casimir forces out of thermal equilibrium

## Idea 2

Two plates at **different** temperature

One plate in equilibrium with environment



## Finite plates

$$P^{(1)}(a, T_1, T_2) =$$

$$\frac{1}{2} [P_{\text{eq}}(a, T_1) + P_{\text{eq}}(a, T_2)] + \Delta P_{\text{neq}}(a, T_1, T_2) + \frac{2\sigma}{3c} (T_2^4 - T_1^4)$$

equilibrium pressures      anti-symmetric term      radiation pressure

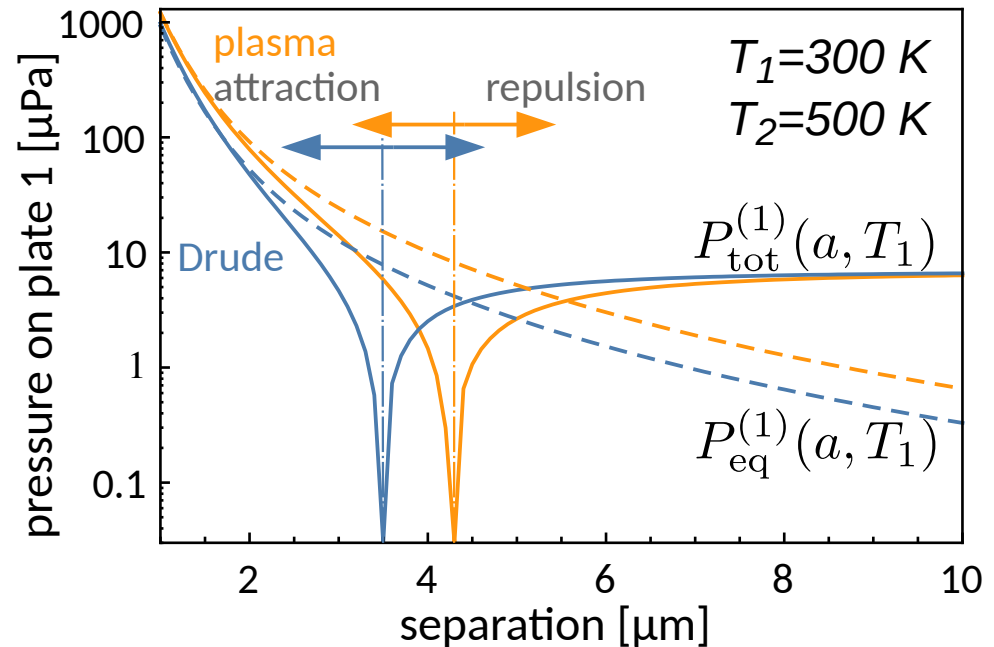
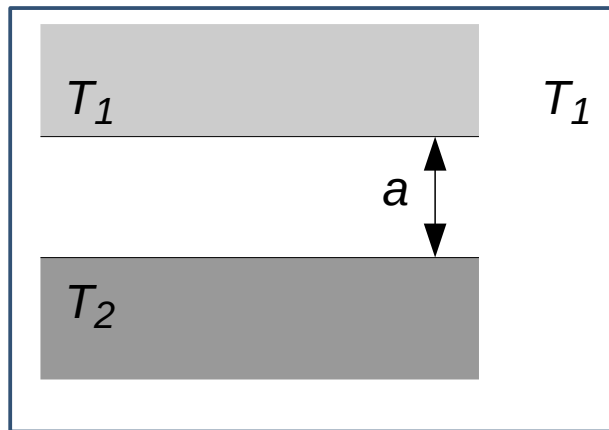
identical plates:  $\Delta P_{\text{neq}} \rightarrow 0$

# Casimir forces out of thermal equilibrium

## Idea 2

Two plates at **different** temperature

One plate in equilibrium with environment



## Finite plates

$$P^{(1)}(a, T_1, T_2) =$$

$$\frac{1}{2} [P_{\text{eq}}(a, T_1) + P_{\text{eq}}(a, T_2)] + \Delta P_{\text{neq}}(a, T_1, T_2) + \frac{2\sigma}{3c} (T_2^4 - T_1^4)$$

equilibrium pressures
anti-symmetric term
radiation pressure

graph data: G. Klimchitskaya, V.M. Mostepanenko, and R.I.P. Sedmik, *PRA* **100**, 022511 (2019)

identical plates:  $\Delta P_{\text{neq}} \rightarrow 0$

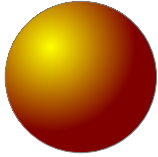
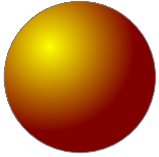
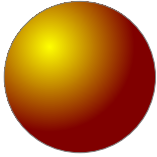
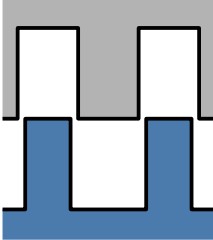

Antezza et al, *PRA* **77**, 022901 (2008)

allows for repulsion, spatially oscillating potentials, bound states, self-propulsion... 22

Can we measure these effects?

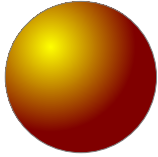
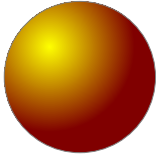
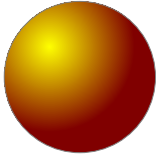
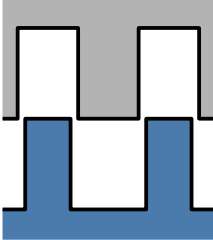

# How to measure these effects?

## A little orders of magnitude comparison

	Torsional balances	AFM-type	Micro-machined oscillators	Micro-beam	Parallel plates
Geometry					
Object size	10 cm	100 $\mu\text{m}$	100 $\mu\text{m}$	1 mm	1 cm
Theory	PFA	PFA	PFA	Full Theory	<b>~exact</b>
Effective Area [ $\text{cm}^2$ ] $A_{eff} \approx \pi R d$	0.01	$10^{-6}$	$10^{-6}$	$10^{-5}$	<b>1</b>
Sensitivity [pN]	1	<b><math>10^{-2}</math></b>	<b><math>10^{-3}</math></b>	1	0.1 (target)

# How to measure these effects?

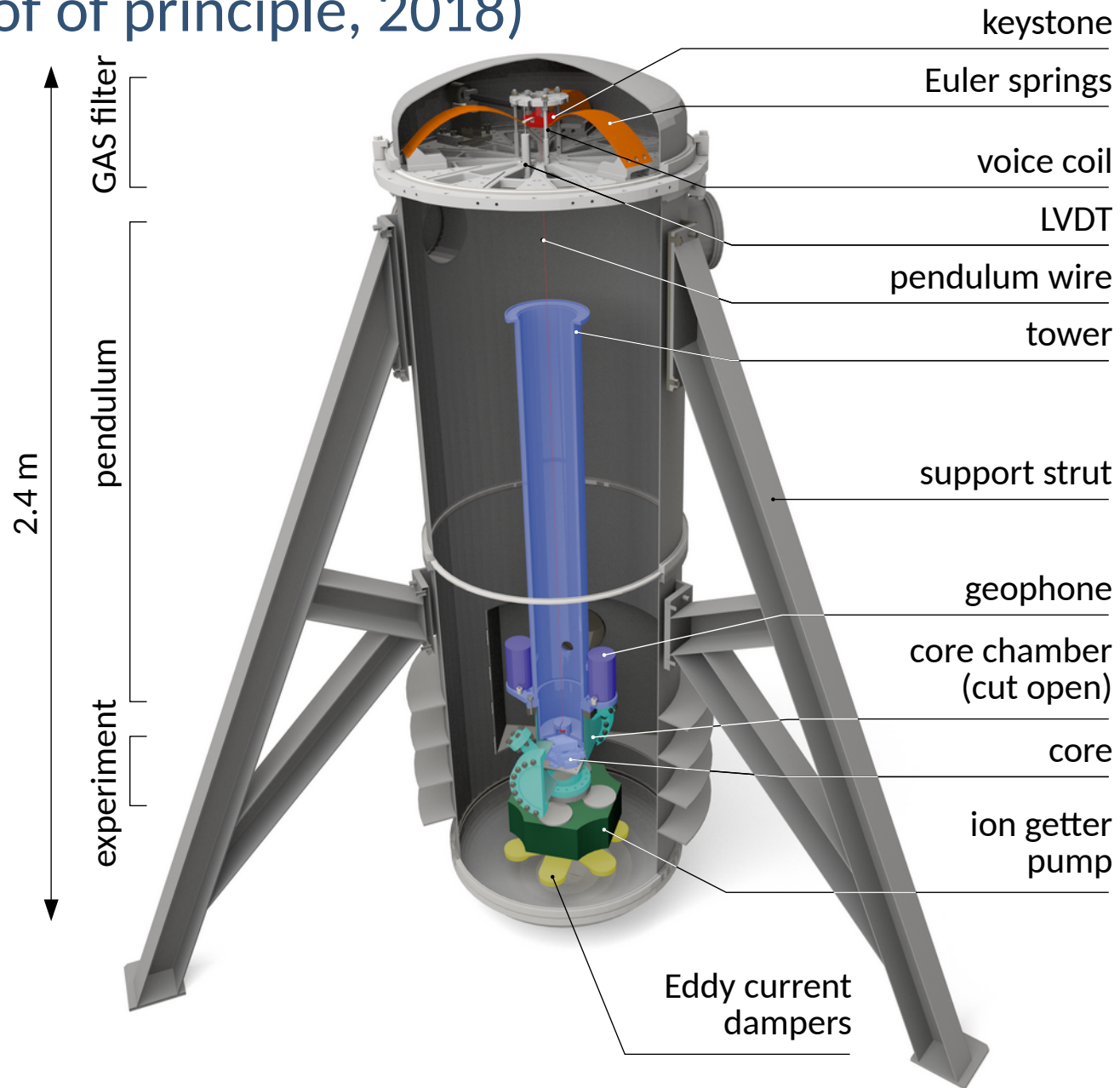
## A little orders of magnitude comparison

	Torsional balances	AFM-type	Micro-machined oscillators	Micro-beam	Parallel plates
Geometry					
Object size	10 cm	100 $\mu\text{m}$	100 $\mu\text{m}$	1 mm	1 cm
Theory	PFA	PFA	PFA	Full Theory	<b>~exact</b>
Effective Area [ $\text{cm}^2$ ] $A_{eff} \approx \pi R d$	0.01	$10^{-6}$	$10^{-6}$	$10^{-5}$	<b>1</b>
Sensitivity [pN]	1	<b><math>10^{-2}</math></b>	<b><math>10^{-3}</math></b>	1	0.1 (target)
[pN/ $\text{cm}^2$ ]	100	$10^4$	1000	$10^5$	<b>0.1</b> (target)
Advantage	Precision	Versatility	Precision	Geometry	<b>Precision</b>



# The CANNEX Setup

# Setup (proof of principle, 2018)

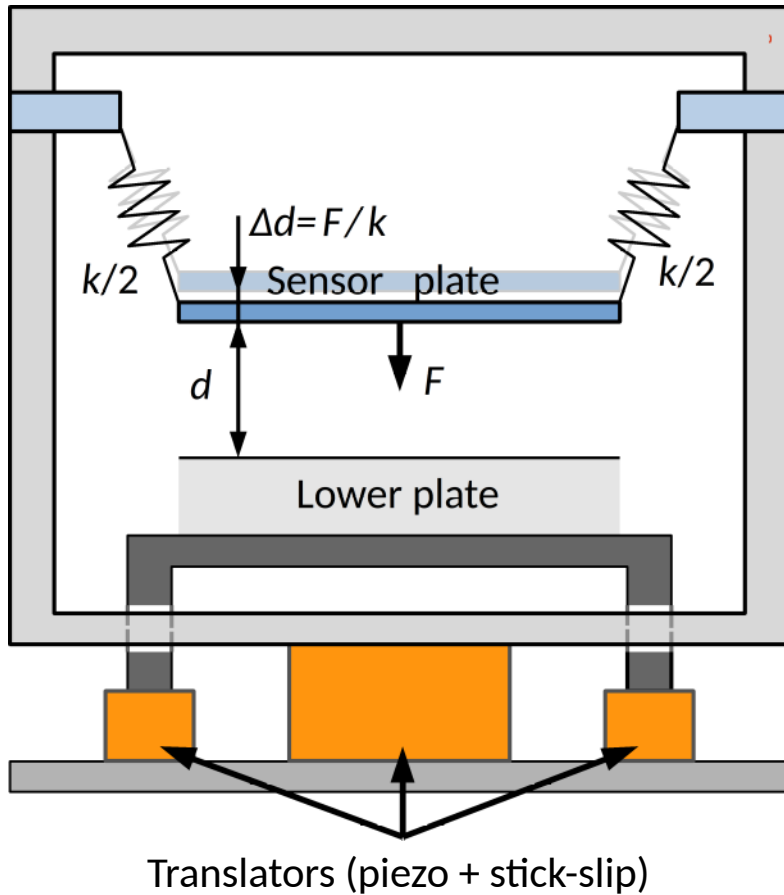


CANNEX overview and isolation system

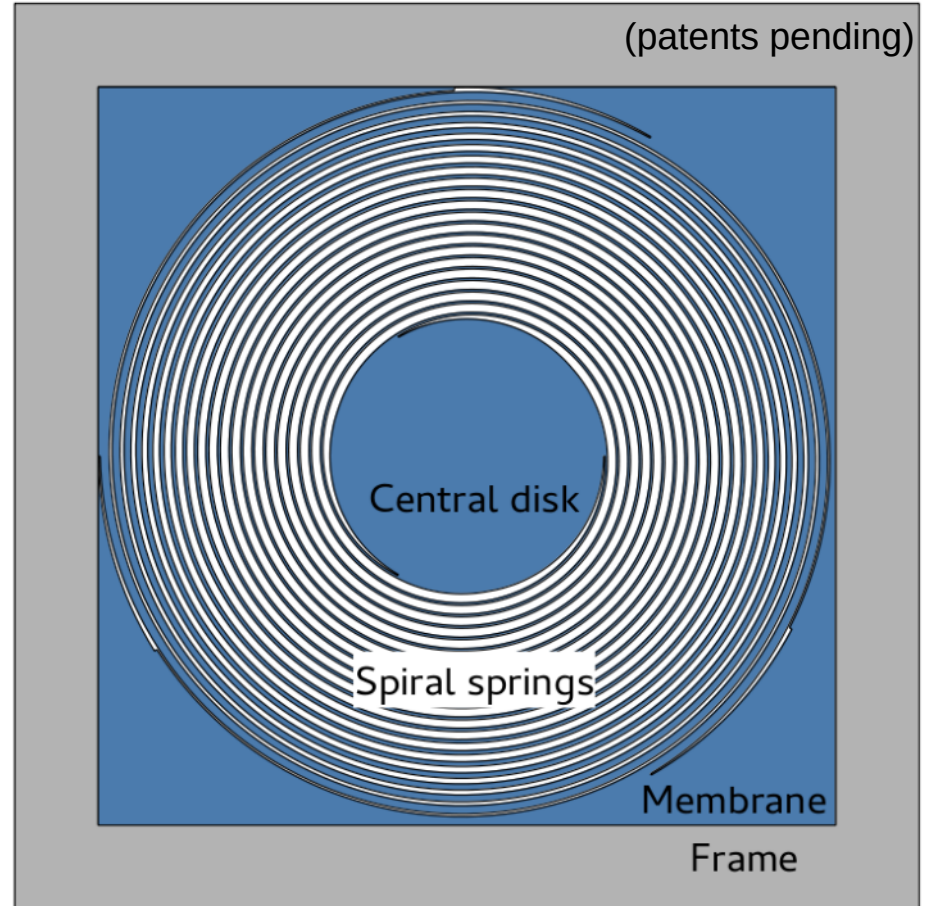
# Force detection

## Principle:

Measure the displacement or resonance freq. of a spring.



## Implementation: upper plate (sensor)



Custom-fabricated Silicon membrane

Force constant:  $0.22 \pm 0.02$  N/m

Disk area :  $1.0834 \pm 0.0005$  cm<sup>2</sup>

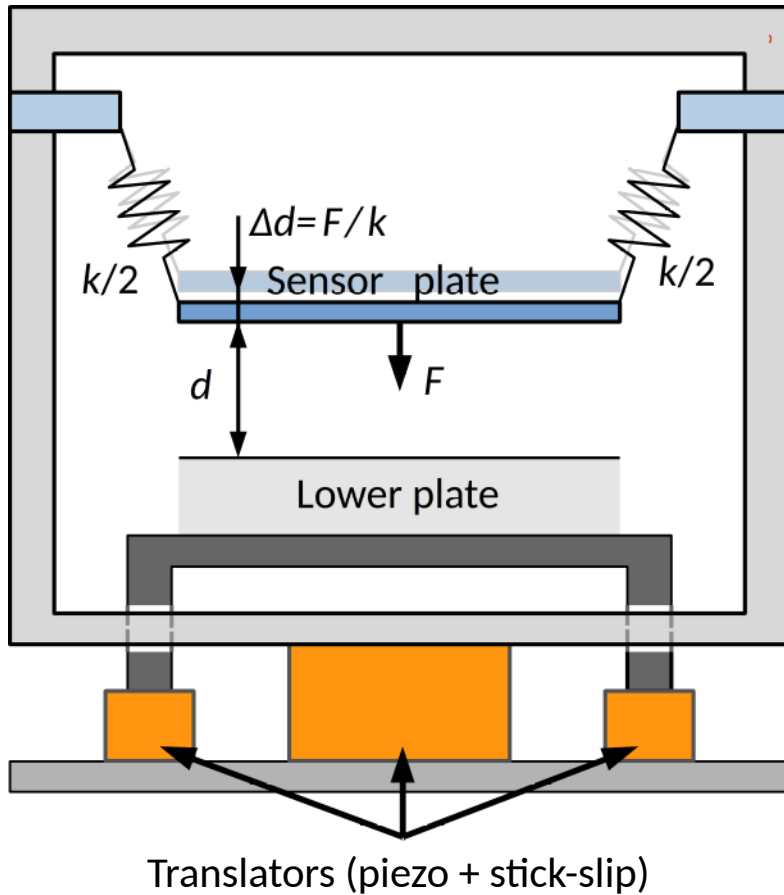
Waviness(disk) < 15 nm (whole area)

Coating: 10 nm TiW + 65 nm Au 28

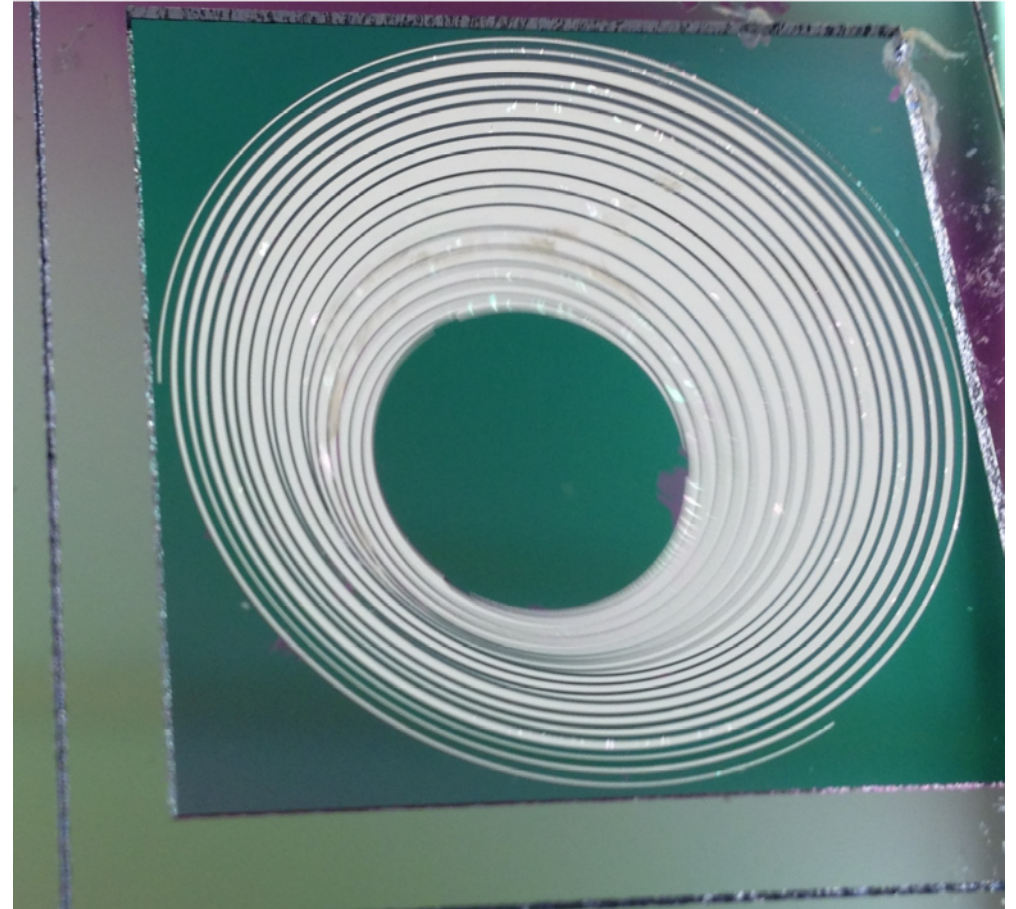
# Force detection

Principle:

Measure the displacement or resonance freq. of a spring.



Implementation: upper plate (sensor)



Custom-fabricated Silicon membrane

Force constant:  $0.22 \pm 0.02$  N/m

Disk area :  $1.0834 \pm 0.0005$  cm<sup>2</sup>

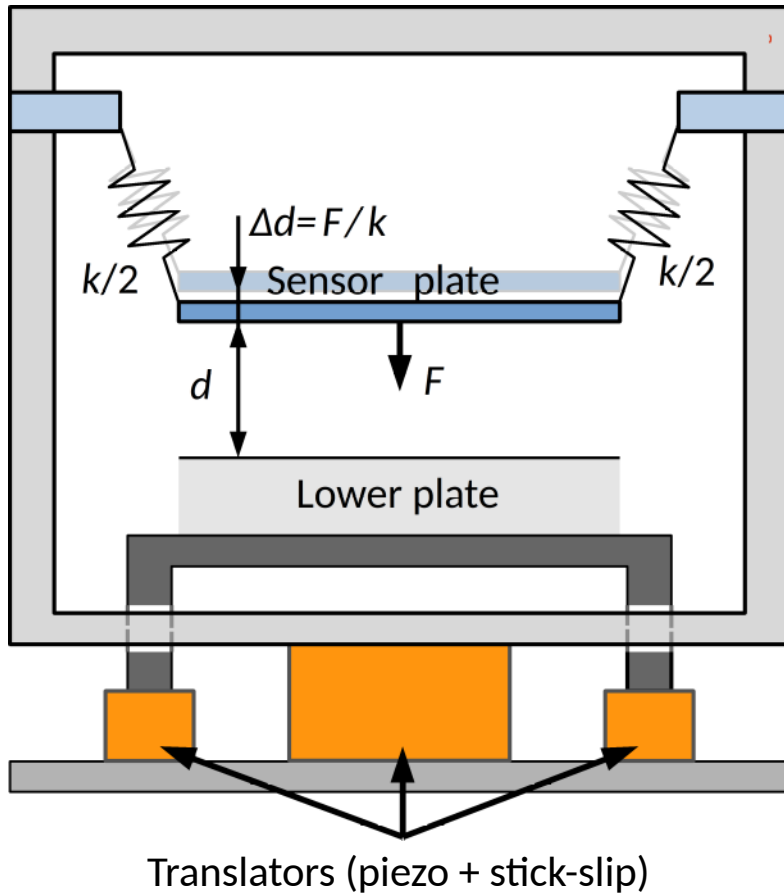
Waviness(disk) < 15 nm (whole area)

Coating: 10 nm TiW + 65 nm Au 29

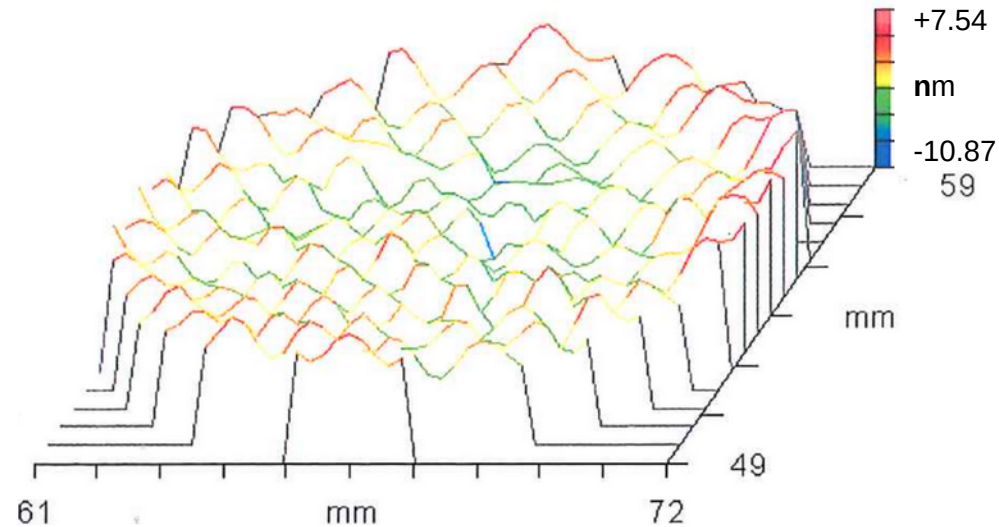
# Force detection

## Principle:

Measure the displacement or resonance frequency of a spring.



## Implementation lower plate:



Custom-fabricated  $\text{SiO}_2$  disk

Thickness 6 mm

Disk area  $1 \text{ cm}^2$

Waviness(disk)  $< 18 \text{ nm}$  (whole area)

Coating 5 nm Cr + 100 nm Au

# Force detection (2018)

Detection Principle:

Forces

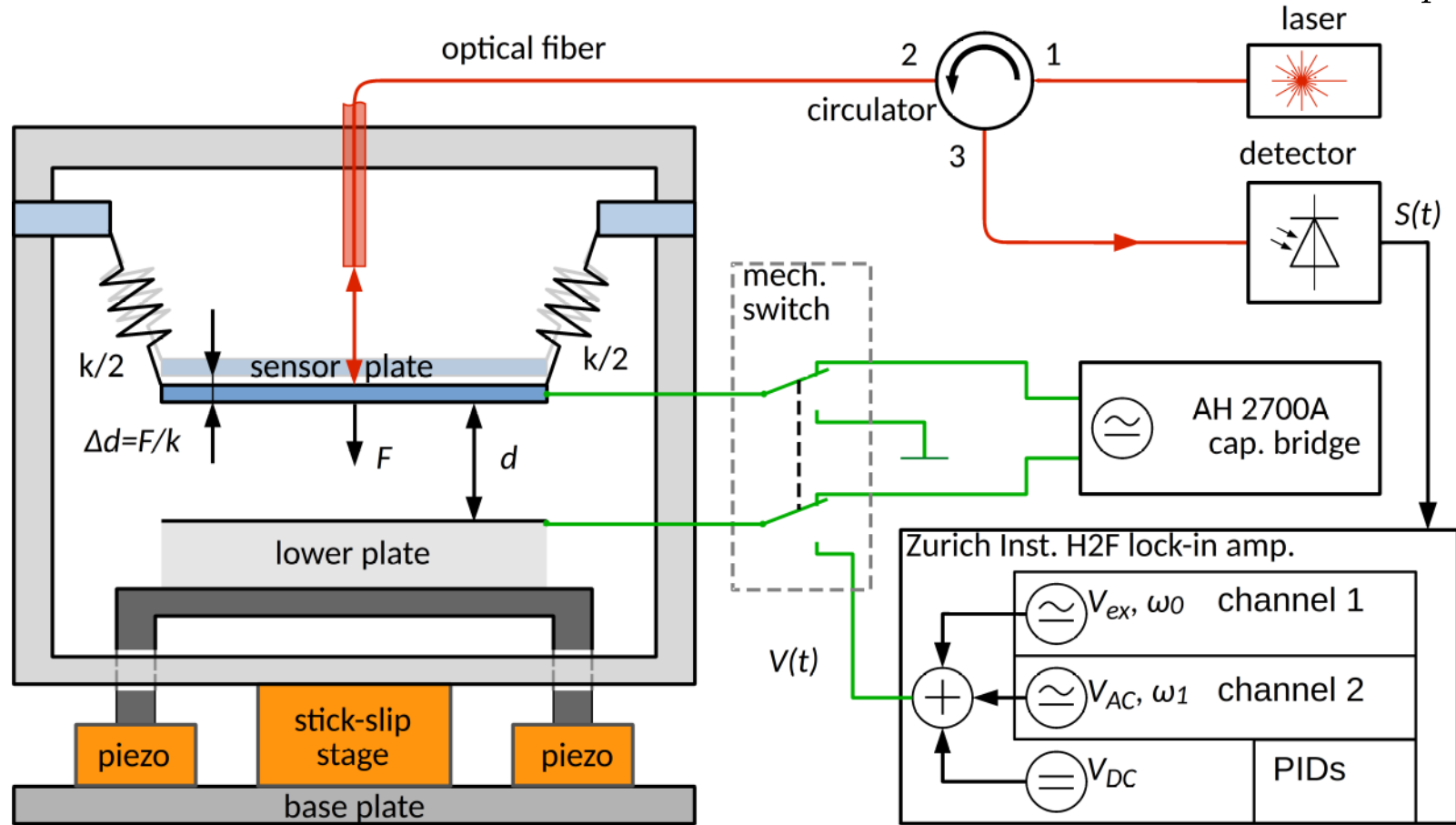
$$F(d) = \Delta d k$$

Force gradients

$$\omega_0 = \sqrt{\frac{k}{m} - \frac{1}{m} \frac{\partial F(d)}{\partial d}}$$

Chameleon: pressure modulation

$$\Delta F_\phi \approx \frac{\partial F_\phi(p)}{\partial p} \Delta p$$



Control: **Parallelism:** tilt modulation + feedback

**Separation:** Electrostatic fit **V<sub>0</sub>:** feedback

# Force detection concepts

Detection  
Principle:

**Forces**

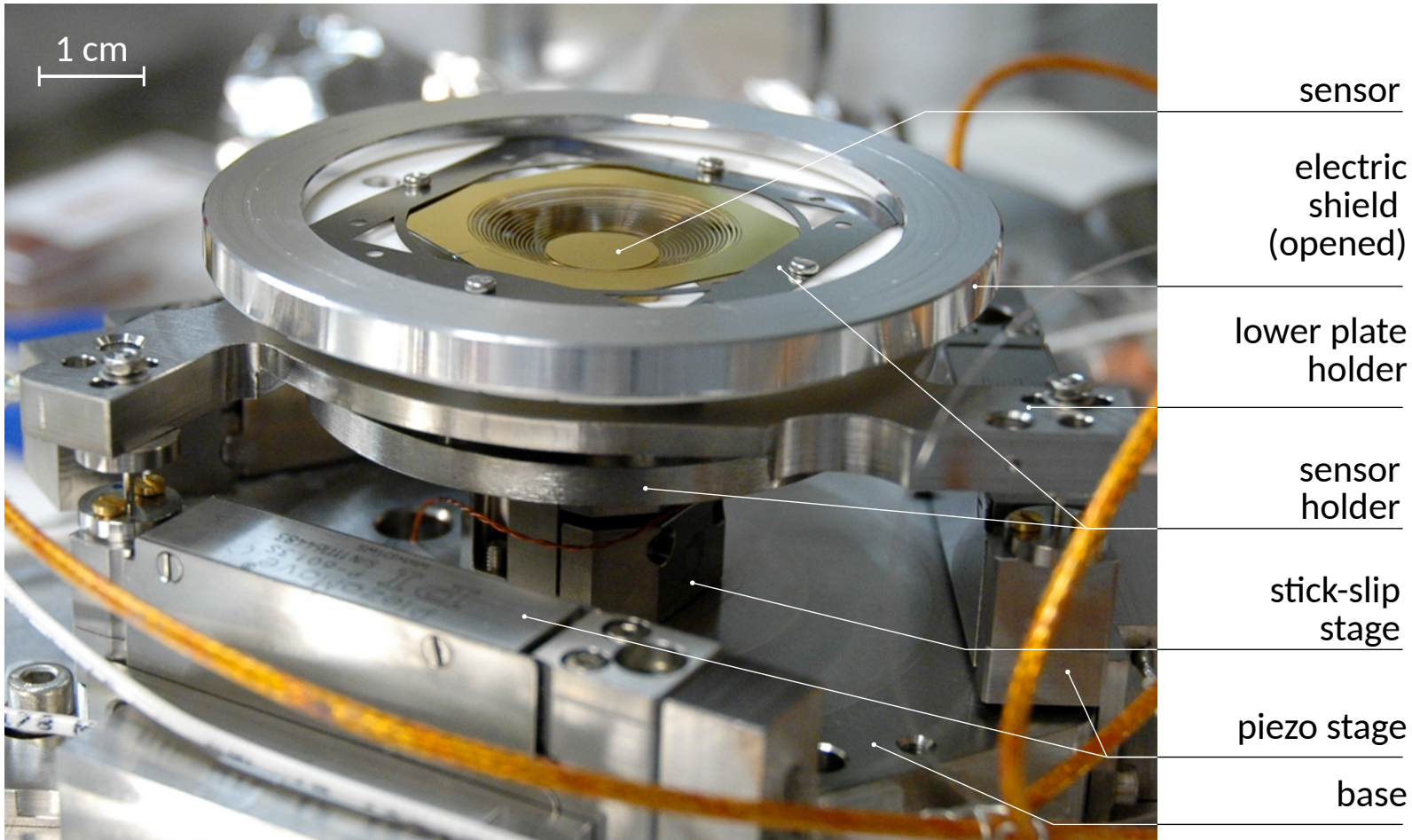
$$F(d) = \Delta d k$$

**Force gradients**

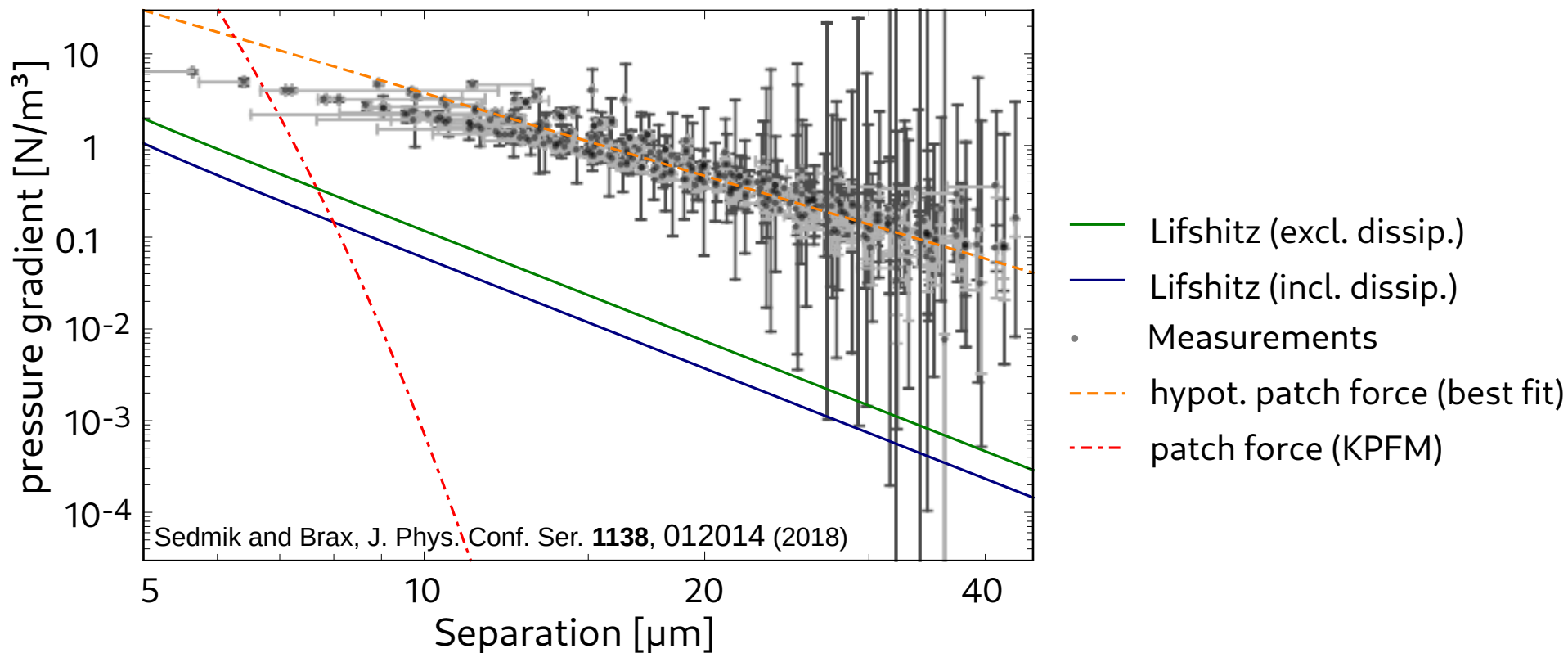
$$\omega_0 = \sqrt{\frac{k}{m} - \frac{1}{m} \frac{\partial F(d)}{\partial d}}$$

**Chameleon: pressure modulation**

$$\Delta F_\phi \approx \frac{\partial F_\phi(p)}{\partial p} \Delta p$$



# First results (under far from ideal conditions)



## Take-home message:

**CANNEX works in principle.  
Proof of concept successful.**

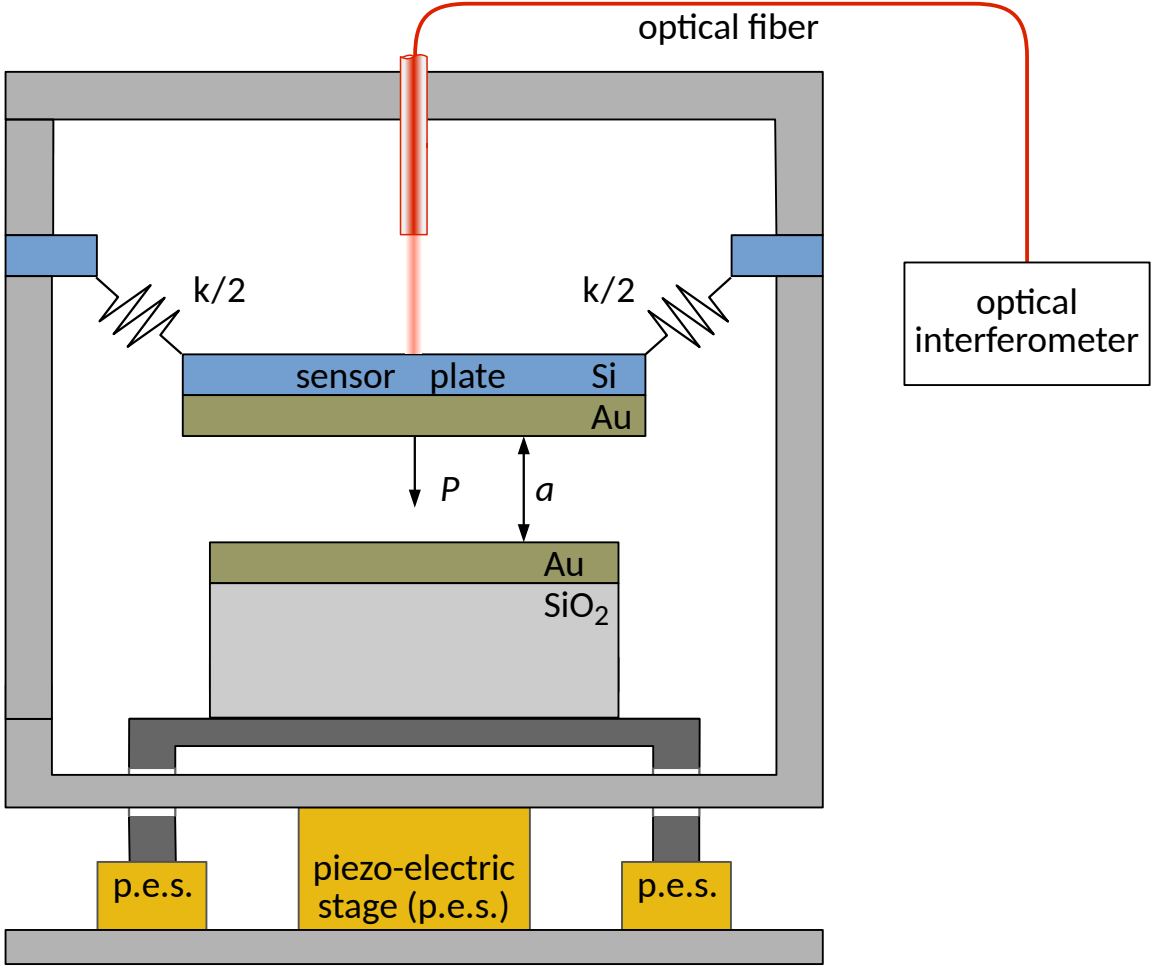
## Calibrated / Measured Parameters

Parallelism (calibrated):	< 200 $\mu\text{rad}$
Residual electrostatic potential:	< 8 $\mu\text{V}$
Drift	< 500 $\text{nm/h}$
Total thermal drift error	< 2.5 $\mu\text{m/run}$



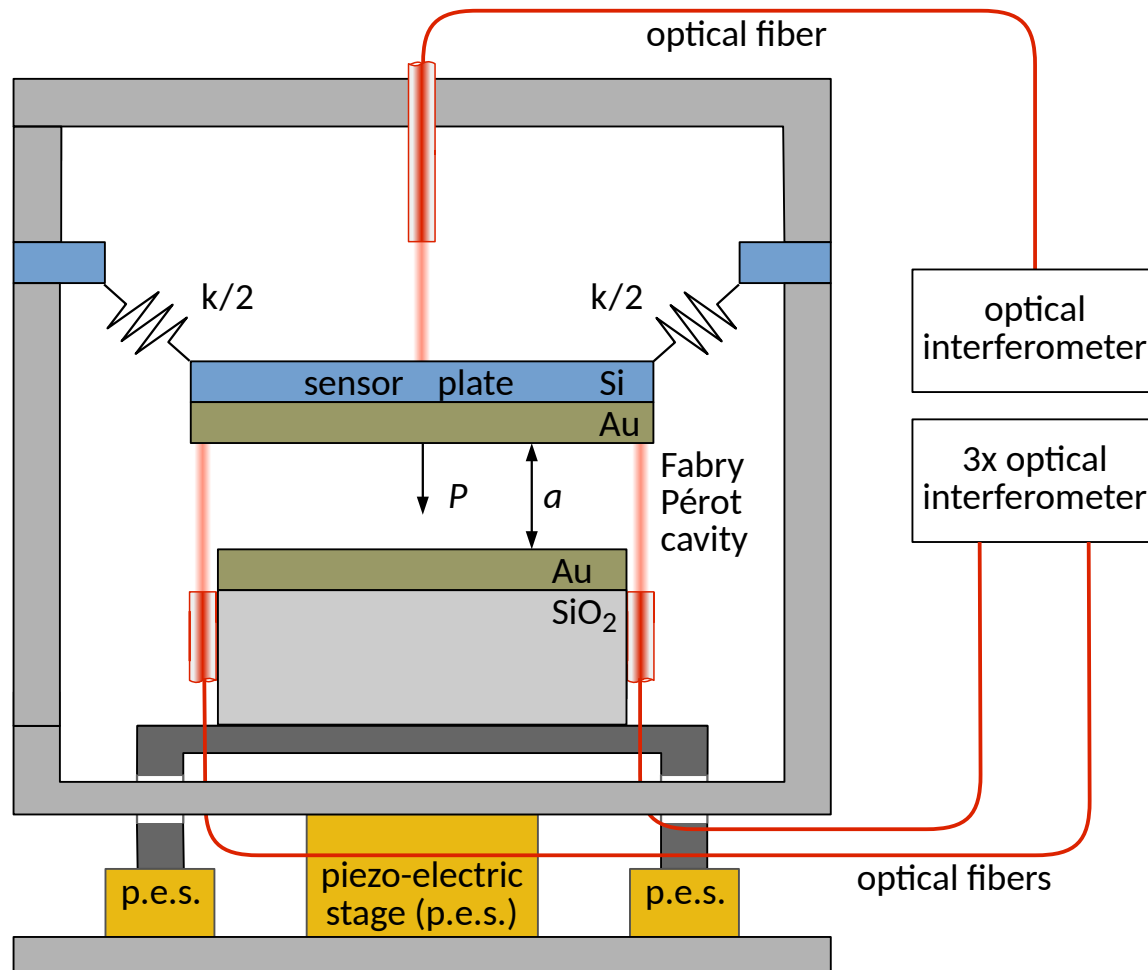
Updates – a new setup

# Planned technical improvements (core)



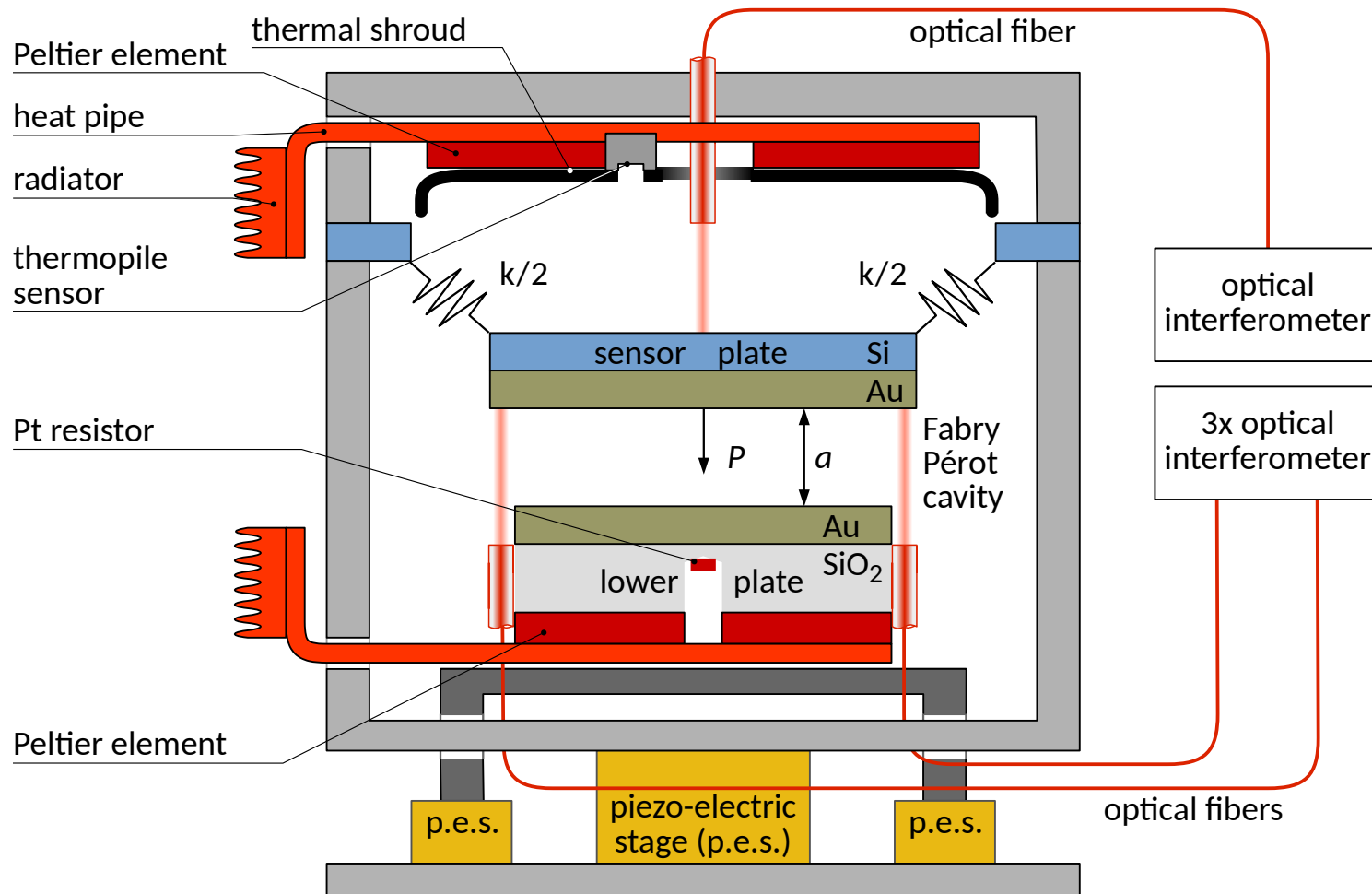
# Planned technical improvements (core)

- All-optical sensing for parallelism and frequency shift → **Less ES background**



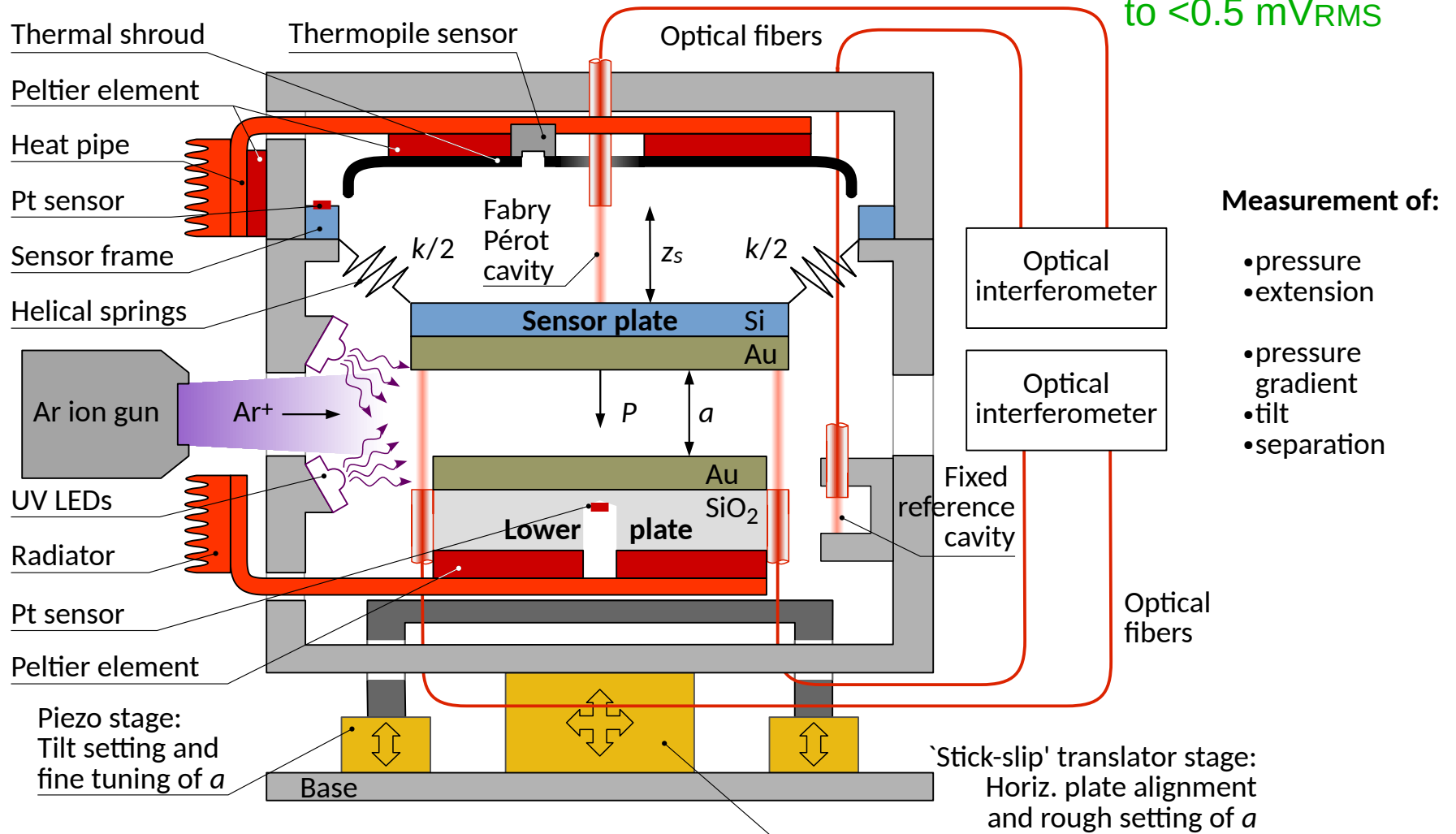
# Planned technical improvements (core)

- All-optical sensing for parallelism and frequency shift → Less ES background
- Independent thermal controls for both plates → Enable  $\Delta T = 10$  K with  $< 1$  mK precision

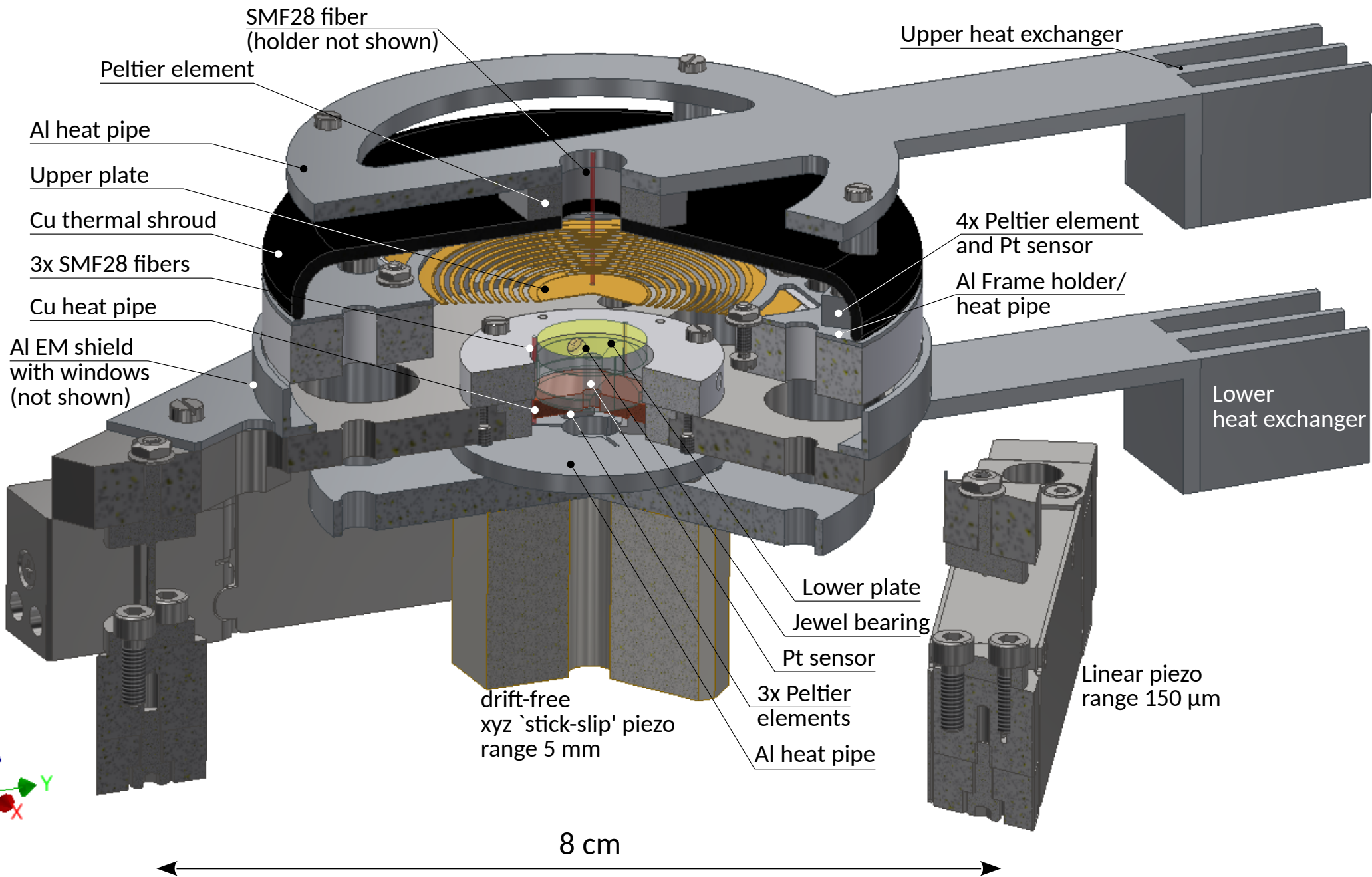


# Planned technical improvements (core)

- All-optical sensing for parallelism and frequency shift → Less ES background
- Independent thermal controls for both plates → Enable  $\Delta T = 10$  K with  $< 1$  mK precision
- In-situ Ar-Ion cleaning and UV irradiation → Reduce surface pot. to  $< 0.5$  mVRMS



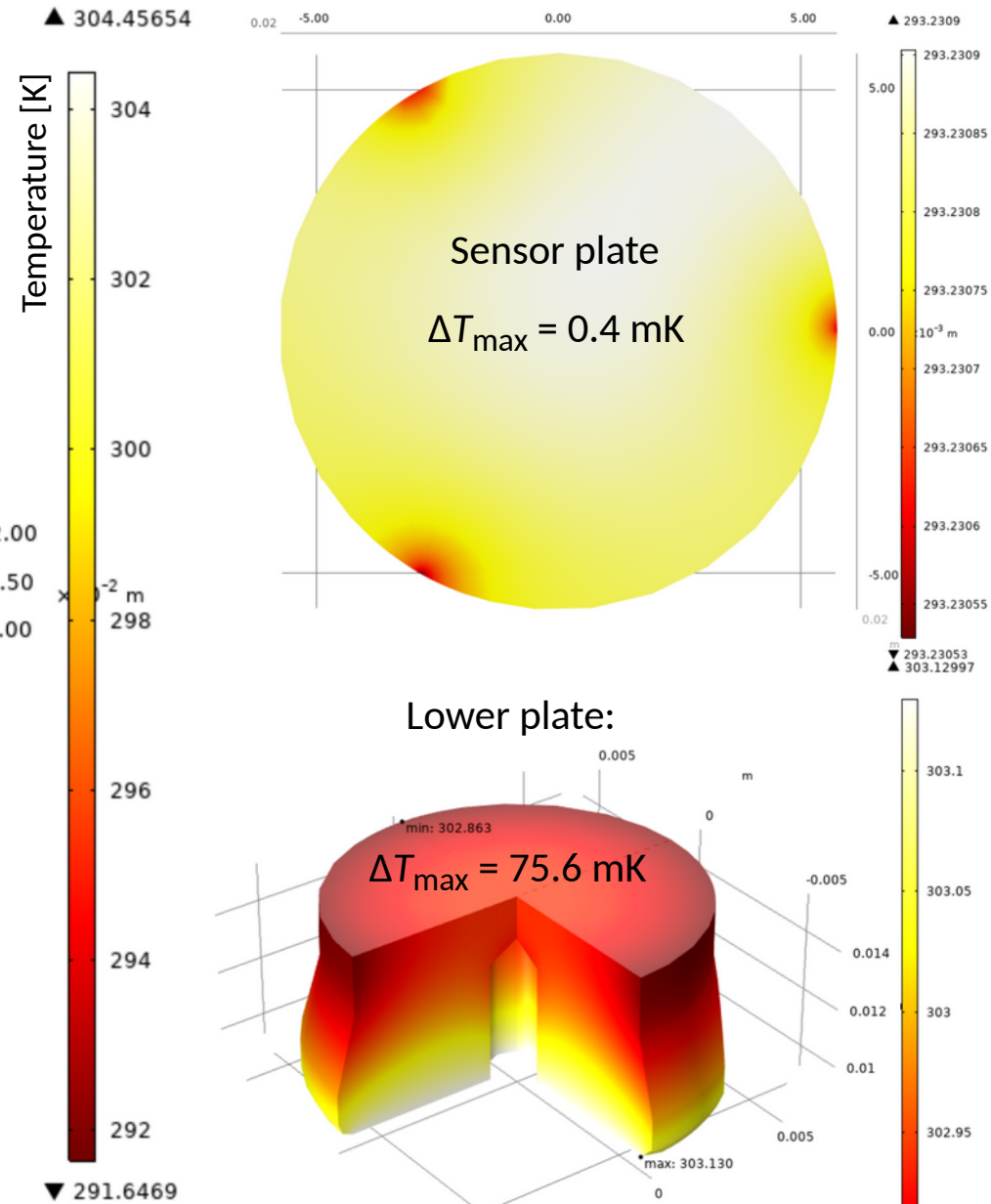
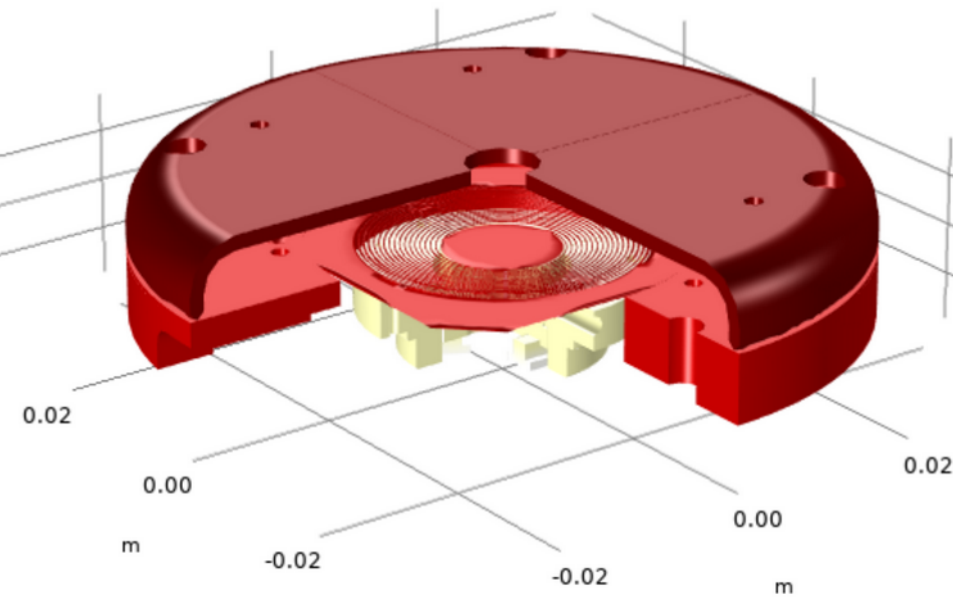
# Planned technical improvements (core) Realization



# Planned technical improvements (core) Realization

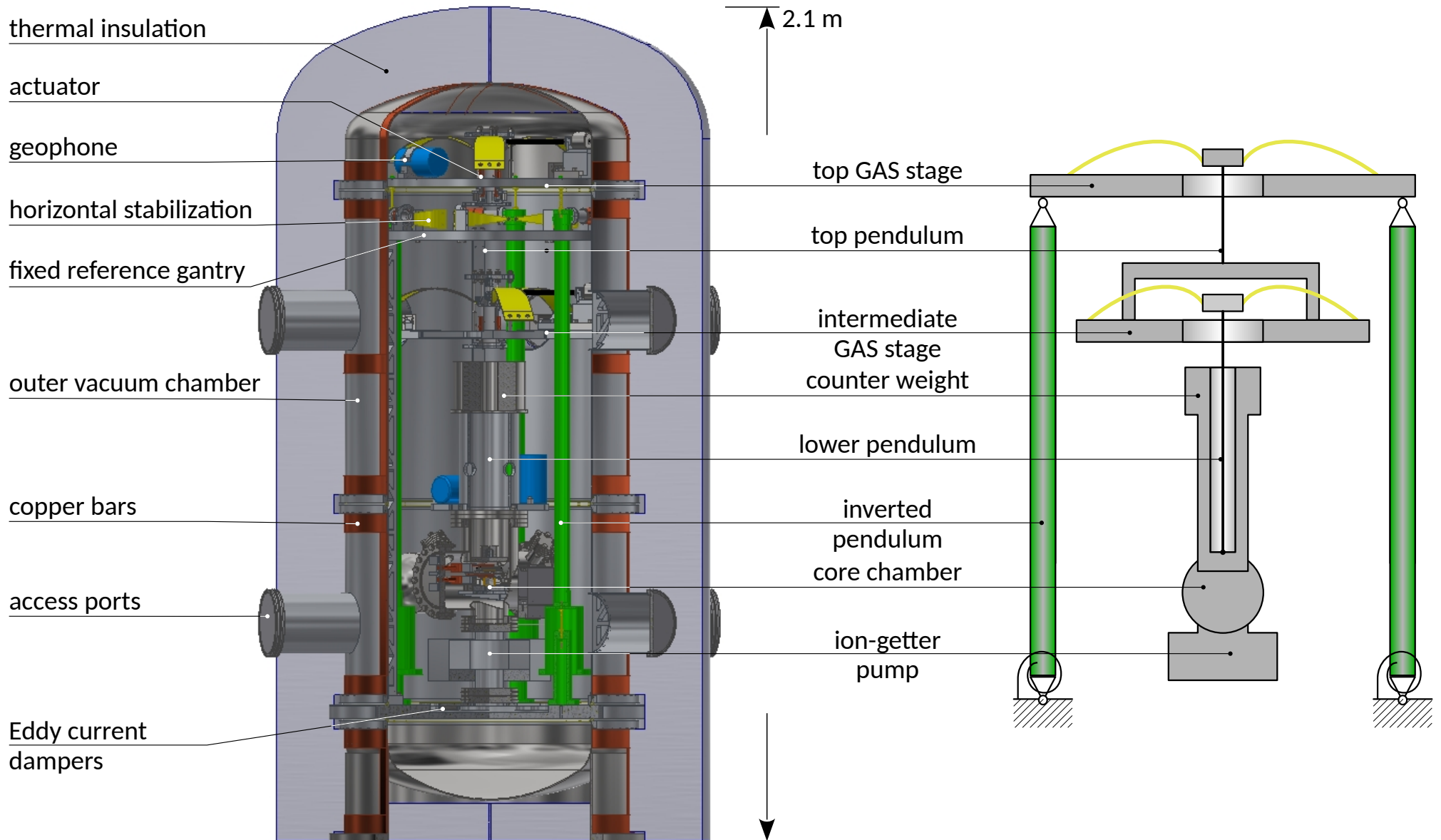
Numerical simulation using actual design and feedback circuits:

**Difference between the plates: 10 K**



**Goal:** <1 mK difference on each surface

# Planned technical improvements: Vibration isolation



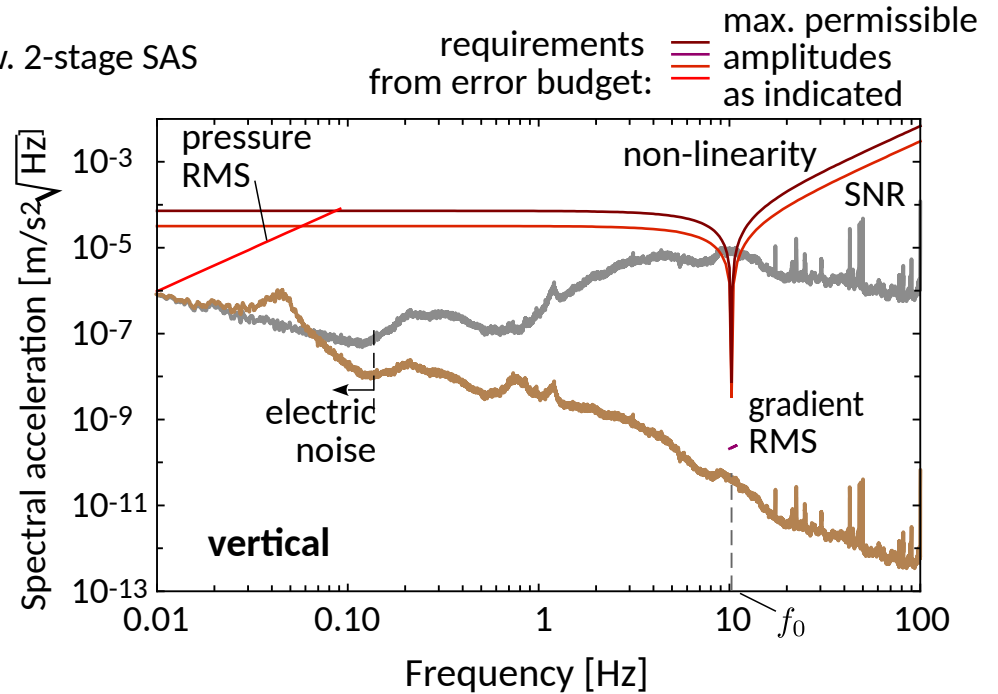
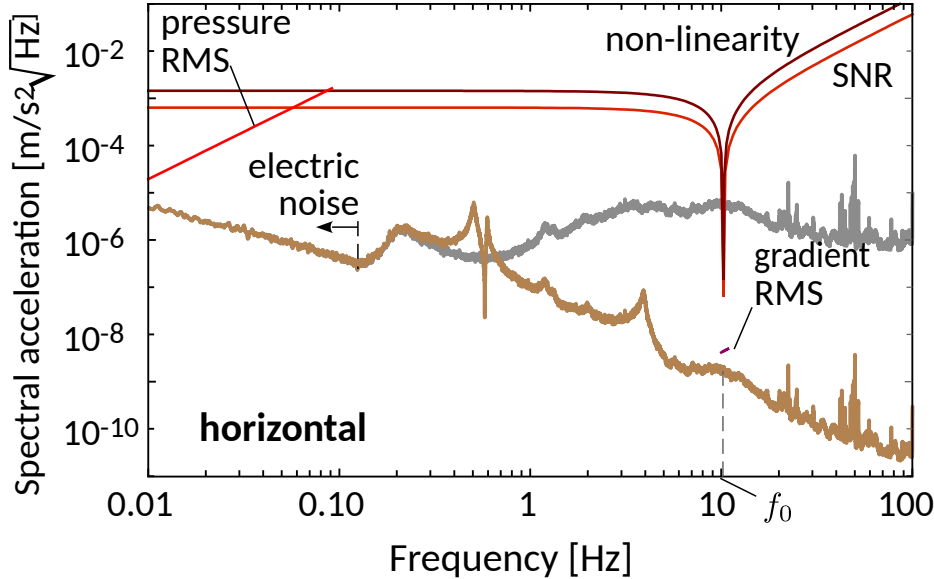


# Planned technical improvements: Vibration isolation (preliminary)



measured background: ● Atominstitut

predicted performance: ● Atominstitut w. 2-stage SAS



requirements from error budget:   
— max. permissible amplitudes as indicated

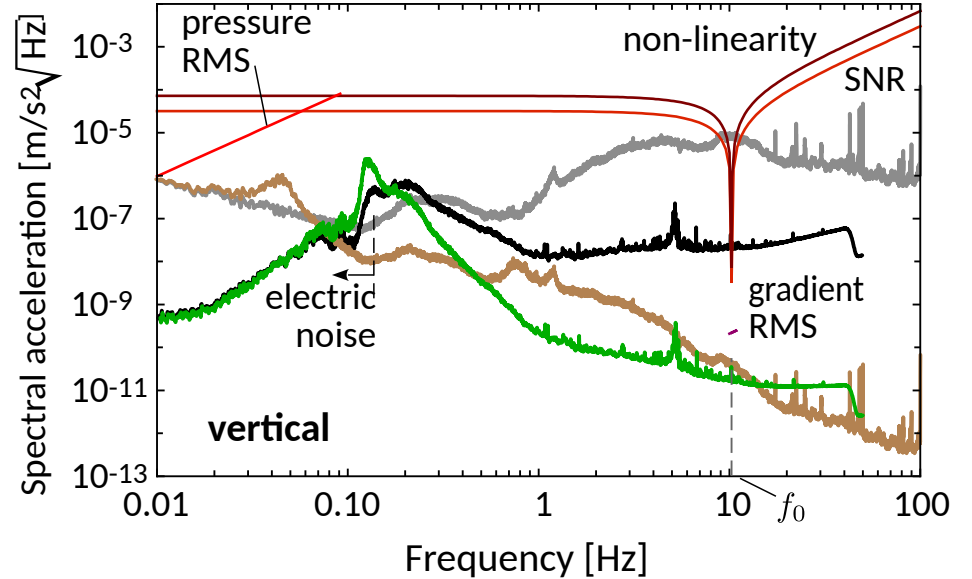
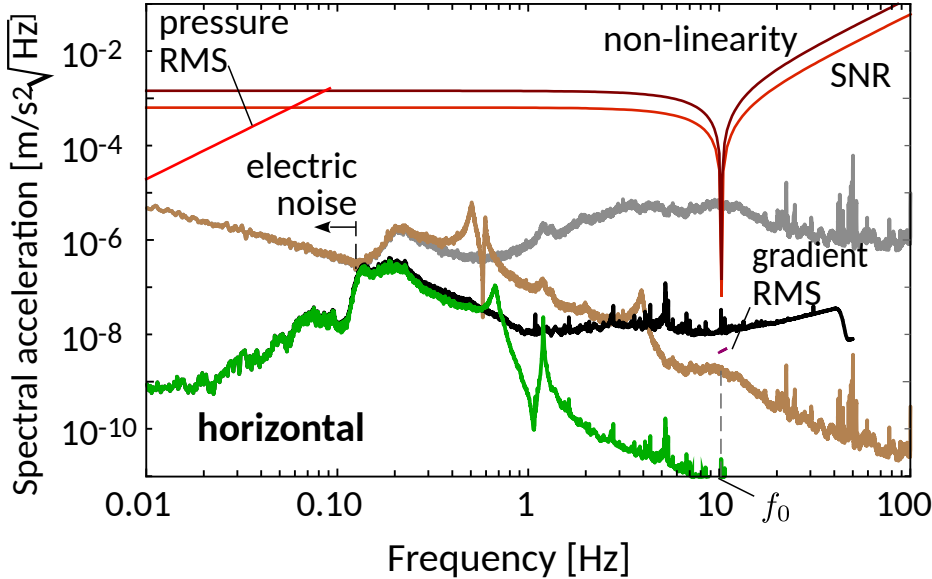
# Planned technical improvements: Vibration isolation (preliminary)



measured background: ● Atominstitut  
● Conrad Observatory

predicted performance: ● Atominstitut w. 2-stage SAS  
● Conrad Observatory w. 1-stage SAS

requirements from error budget: max. permissible amplitudes as indicated



# Metrology

Considered detailed data and models on all relevant systematic and statistical noise sources  
For force measurements and gradient measurements different:

- **Seismic:**
  - Non-linearity (stat.)
  - RMS noise in different bandwidths (stat.)
  - Direct force noise (stat.)
  - Phase noise (stat.)
- **Surface deformations:**
  - Stochastic roughness (stat.)
  - large-scale deformations (sys. + stat.)
- **Residual tilt** (stat.)
- **Thermal fluctuations:**
  - Thermal expansion (stat.)
  - Changes of the Young's modulus (stat.)
- **Mechanical Brownian noise:** (stat.)
- **Detector noise** (stat. + sys.)
- **Phase resolution (PLL)** (sys.)
- **Calibration uncertainties** (stat.)
- **Laser stability/Allen deviation** (stat.)
- **Optical path variations** (sys.)
- **Electrostatic surface potentials** (sys.)
- **Voltage noise** (stat.)

**Detection**

**Total error:**

$$\delta_{\text{tot}} = \left[ \sum (\text{stat.})^2 \right]^{1/2} + \sum (\text{sys.})$$

# Metrology

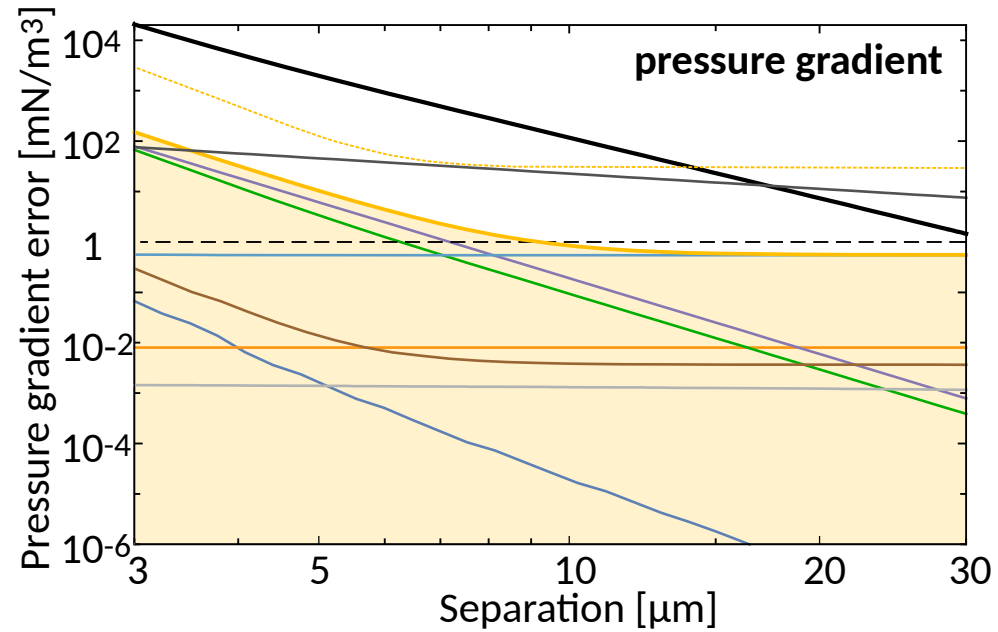
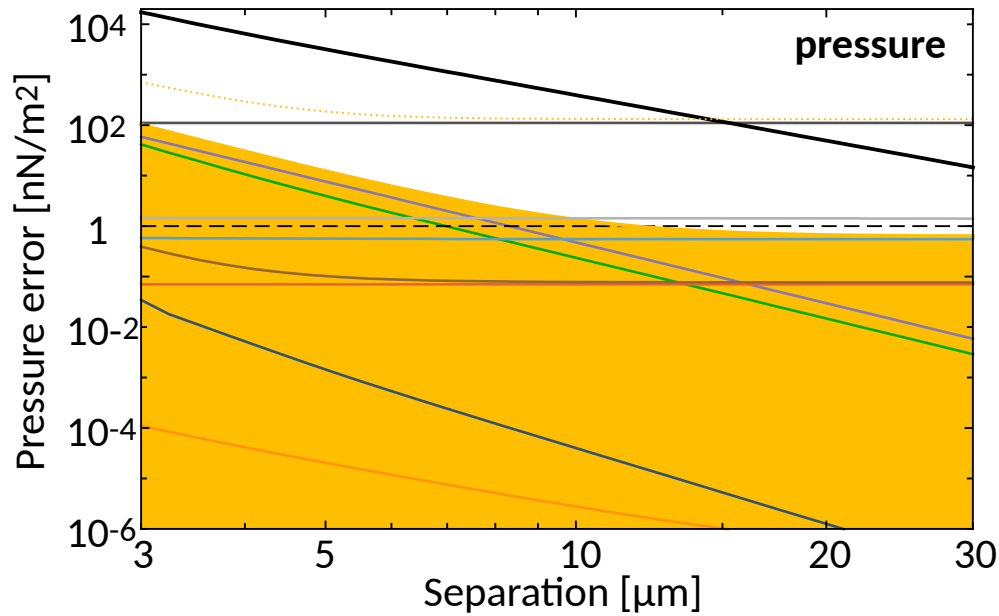
Considered detailed data and models on all systematic and statistical noise sources  
 For force measurements and gradient measurements different:

## Error contributions:

- Electrostat. patches — Corrugations — Detection — Vibrations
- Tilt — Brownian noise — Thermal variations — **Total** — **Prev.**

## Signals:

- Casimir interaction — Electrostatic excitation
- Gravitational interaction — Target sensitivity



Targeted sensitivity 1 nN/m<sup>2</sup> and 1 mN/m<sup>3</sup> achievable at >10 μm

At all separations Measurement of the Casimir force at % level

Measurement options:  $g_S^2, g_P^2, g_S g_P, g_V^2, G(r)$

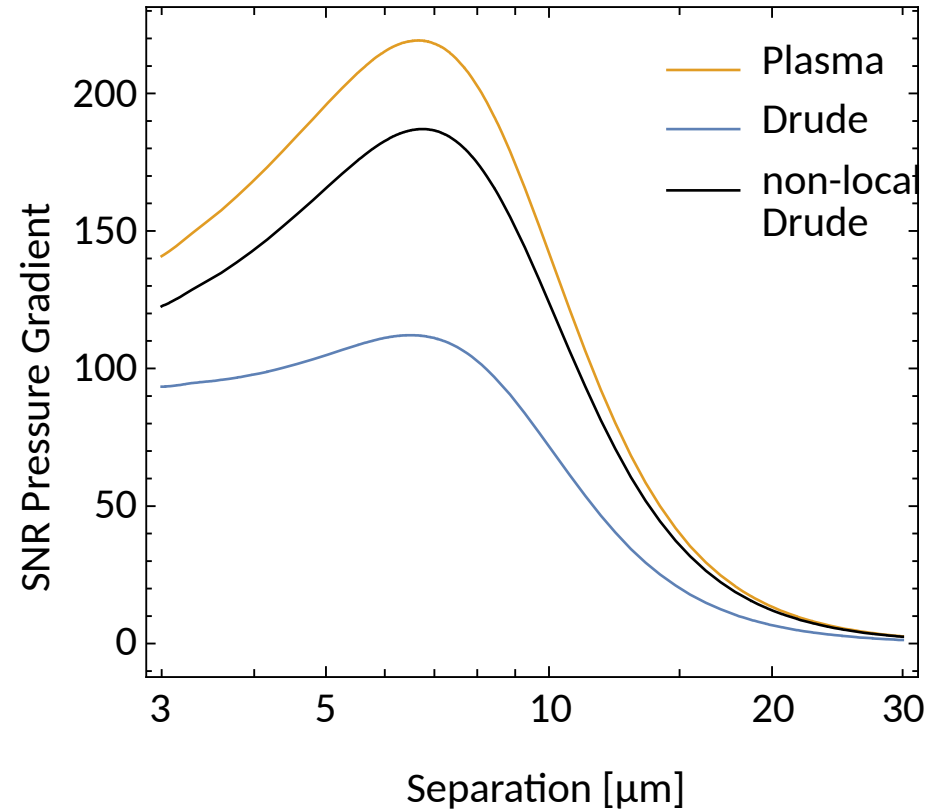
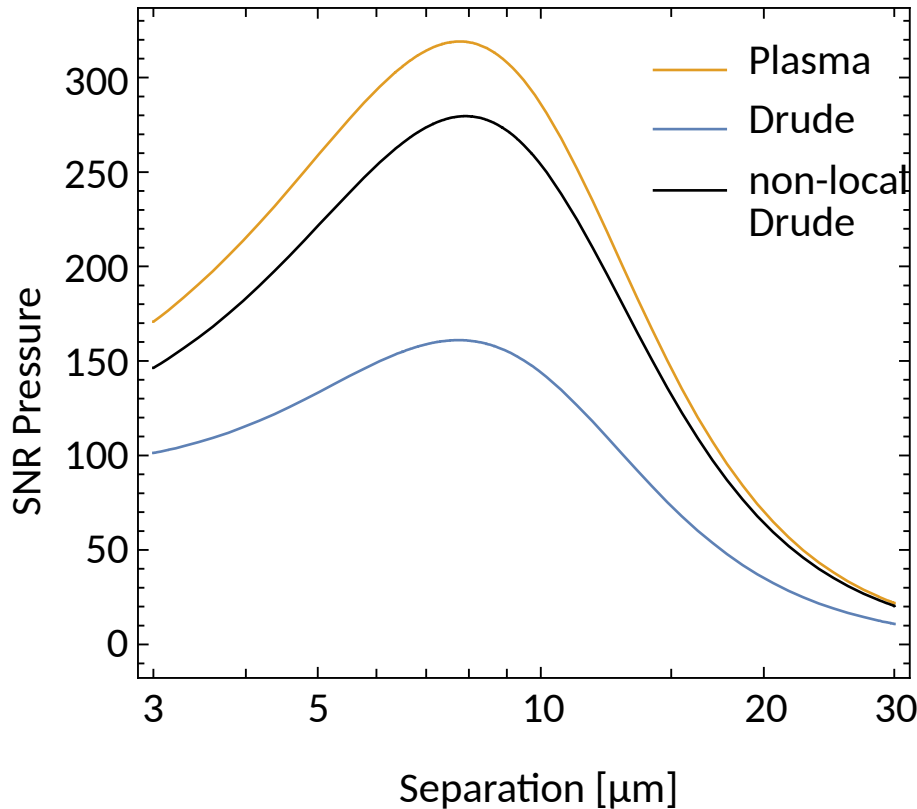
Advantage:  $N = 10^{23}$  amplification factor

# What can be achieved?

First accurate Casimir measurement beyond 10  $\mu\text{m}$  with parallel plates

Cannex is the only experiment that can distinguish all available models

Klimchitskaya and Mostepanenko *EPJ* **80** 900 (2020)  
Sedmik and Pitschmann, *Universe*, **7** 234 (2021)



+ First accurate Casimir measurement out of thermal equilibrium ( $\Delta T = 10$  K)

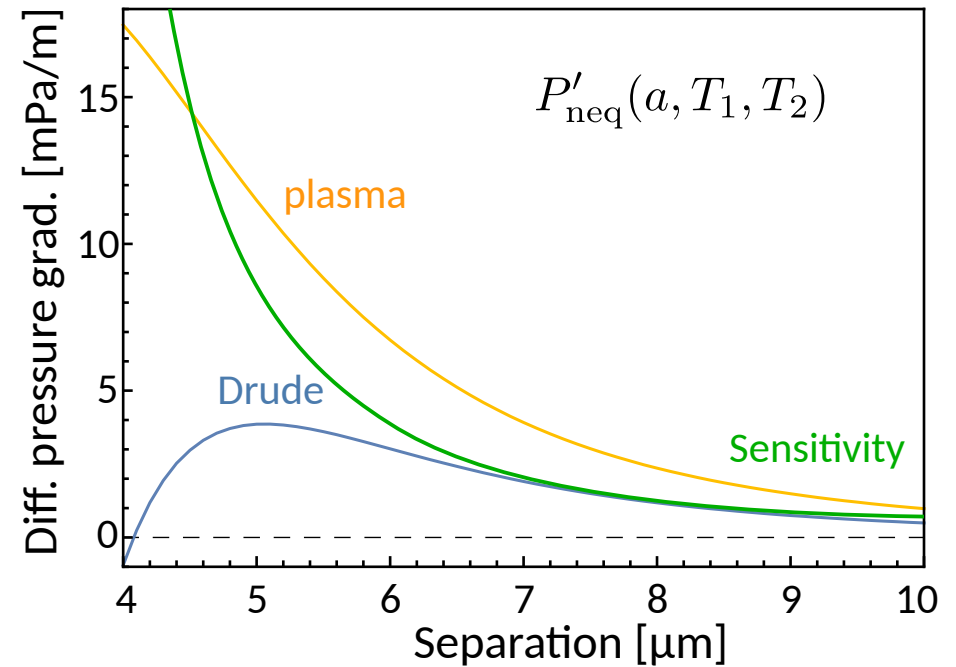
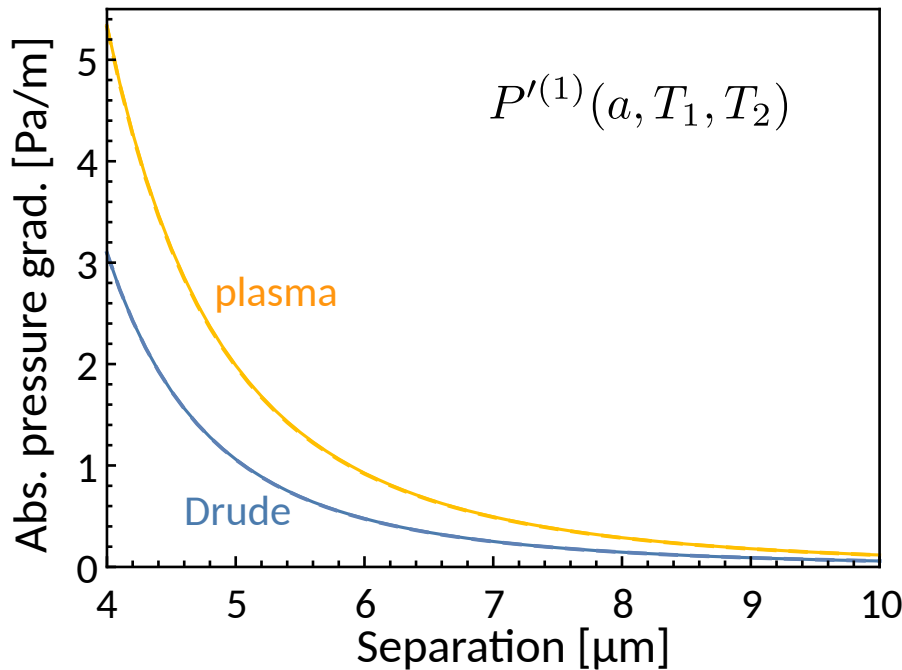
G. Klimchitskaya, V.M. Mostepanenko, R.I.P. Sedmik, and H. Abele, *Symmetry*, **11**(3) 407 (2019)

# What can be achieved?

First measurement of non-equilibrium thermal contributions: **gradient**

$$P'^{(1)}(a, T_1, T_2) = \frac{1}{2}[P'_{\text{eq}}(a, T_1) + P'_{\text{eq}}(a, T_2)] + \boxed{P'_{\text{neq}}(a, T_1, T_2)}$$

$T_1 = 300 \text{ K}$  (sensor),  $T_2 = 310 \text{ K}$  (plate)



graph data: G. Klimchitskaya, V.M. Mostepanenko,  
and R.I.P. Sedmik, *PRA* **100**, 022511 (2019)

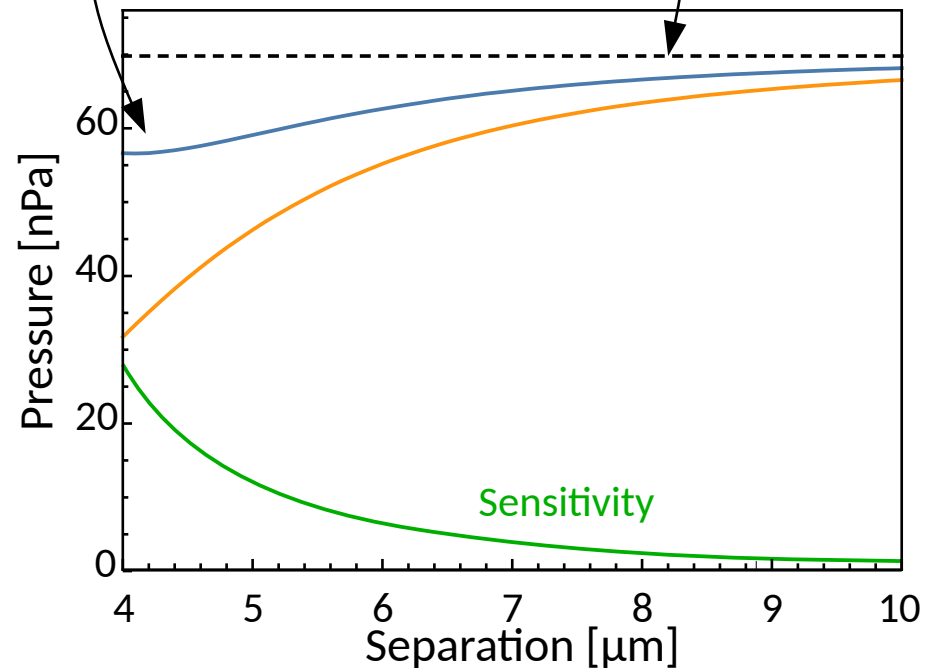
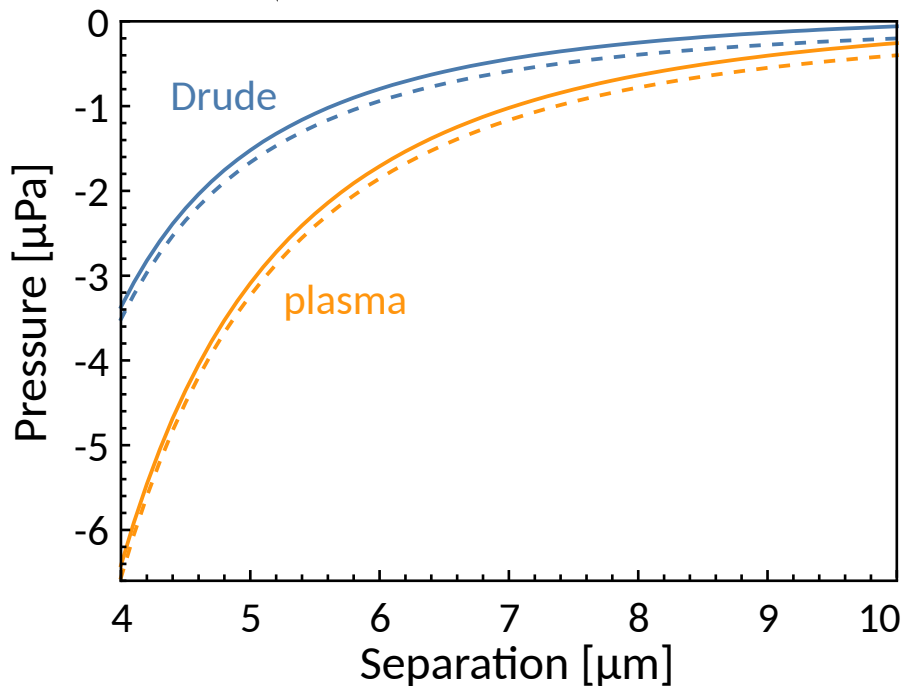
**Difficult but possible with SNR ~ 1.8 (plasma)**

# What can be achieved?

Measurement of separation-independent radiation **pressure**

$$P'^{(1)}(a, T_1, T_2) = \frac{1}{2} [P'_{\text{eq}}(a, T_1) + P'_{\text{eq}}(a, T_2)] + \Delta P'_{\text{neq}}(a, T_1, T_2) + \frac{2\sigma}{3c} (T_2^4 - T_1^4)$$

$T_1 = 300 \text{ K}$  (sensor),  $T_2 = 310 \text{ K}$  (plate)

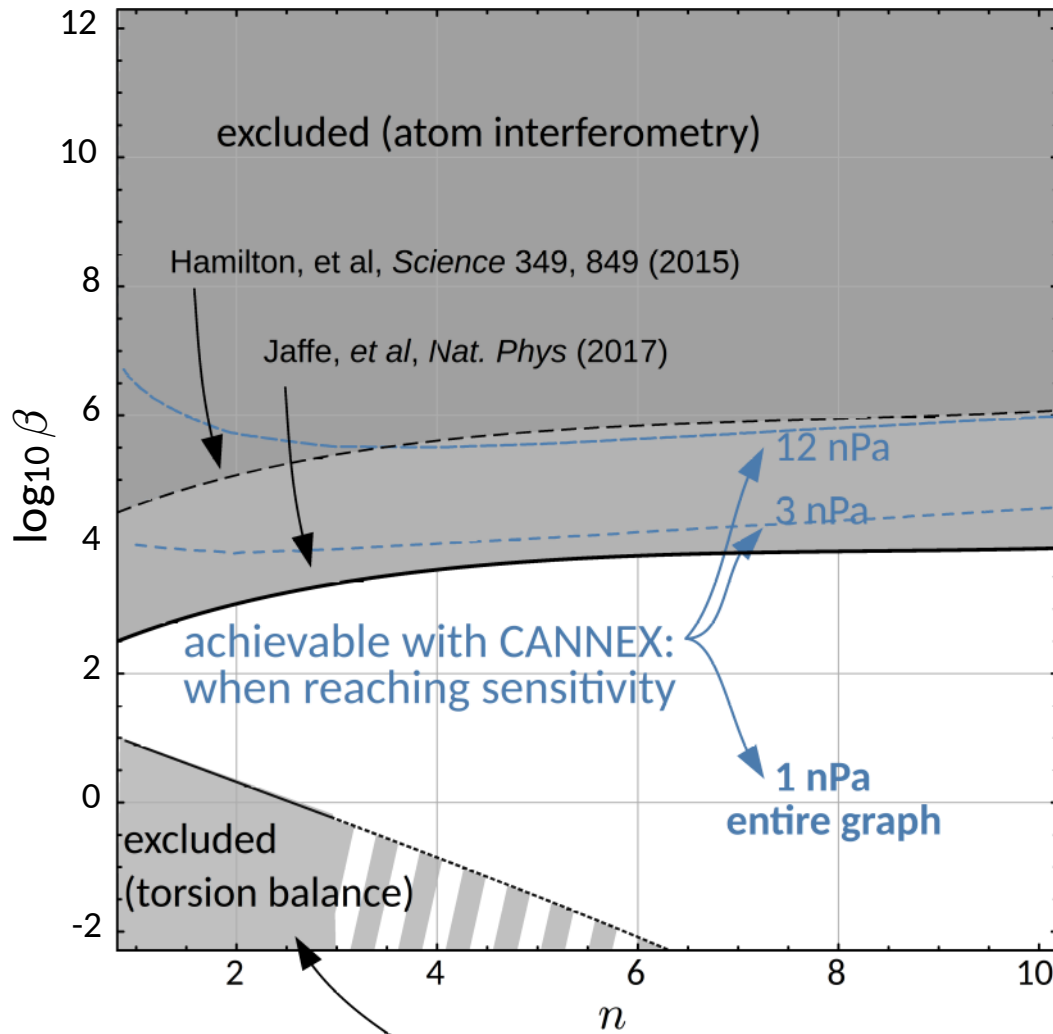


graph data: G. Klimchitskaya, V.M. Mostepanenko, and R.I.P. Sedmik, *PRA* **100**, 022511 (2019)

**Clearly possible to measure (SNR ~ 5)**

# What can be achieved?

## Exclusion of the Chameleon as dark energy candidate



Adelberger, et al, Prog. Part. Nucl. Phys. **62**, 102 (2009)

Dynamic scalar field:

$$\mathcal{L} = \sqrt{-g} \left[ \frac{R}{16\pi G} - \frac{(\partial\phi)^2}{2} + V(\phi) \right] + \mathcal{L}_{\text{sm}}(\psi, \tilde{g}_{\mu\nu})$$

$$\tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}$$

Coupling:  $A(\phi) = e^{\beta\phi/m_{\text{Pl}}}$

Potential:  $V(\phi) = \Lambda^4 \left( 1 + \frac{\Lambda^n}{\phi^n} \right)$

$$\Lambda = \Lambda_0 \approx 2.4 \text{ meV}$$

Effective potential:

$$V_{\text{eff}} = V(\phi) + \rho e^{\frac{\beta}{M_{\text{Pl}}}\phi}$$

Measurement via vacuum pressure modulation

Complete exclusion possible

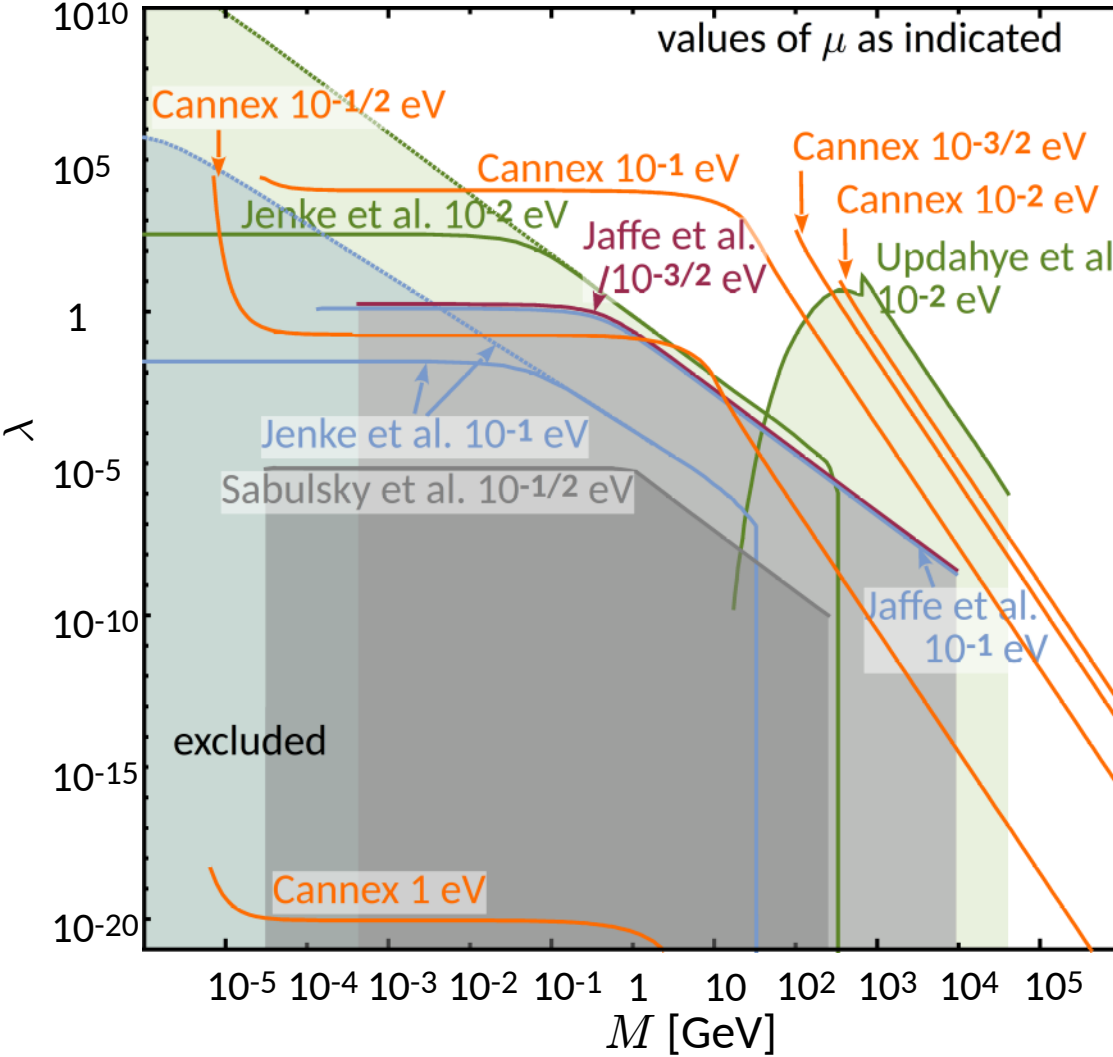


# What can be achieved?

## Limits on the Symmetron as dark energy candidate

Sedmik and Pitschmann, *Universe* 7, 234 (2021)

values of  $\mu$  as indicated



Dynamic scalar field:

$$\mathcal{L} = \sqrt{-g} \left[ \frac{R}{16\pi G} - \frac{(\partial\phi)^2}{2} + V(\phi) \right] + \mathcal{L}_{\text{sm}}(\psi, \tilde{g}_{\mu\nu})$$

$$\tilde{g}_{\mu\nu} = A^2(\phi) g_{\mu\nu}$$

Coupling:  $A(\phi) = 1 + \frac{\phi^2}{2M^2}$

Potential:  $V(\phi) = -\frac{\mu^2}{2}\phi^2 + \frac{\lambda}{4}\phi^4$

At  $\rho < M^2\mu^2$  symmetry spontaneously broken and acquires  $\langle\phi\rangle_{\text{vac.}} = \sqrt{\mu^2/\lambda}$

At  $\rho \gg M^2\mu^2$  symmetry restored,  $\langle\phi\rangle_{\text{vac.}} = 0$  no coupling to matter

Effective potential:

$$V_{\text{eff}} = \frac{1}{2} \left( \frac{\rho}{M^2} - \mu^2 \right) \phi^2 - \frac{\lambda}{4} \phi^4$$

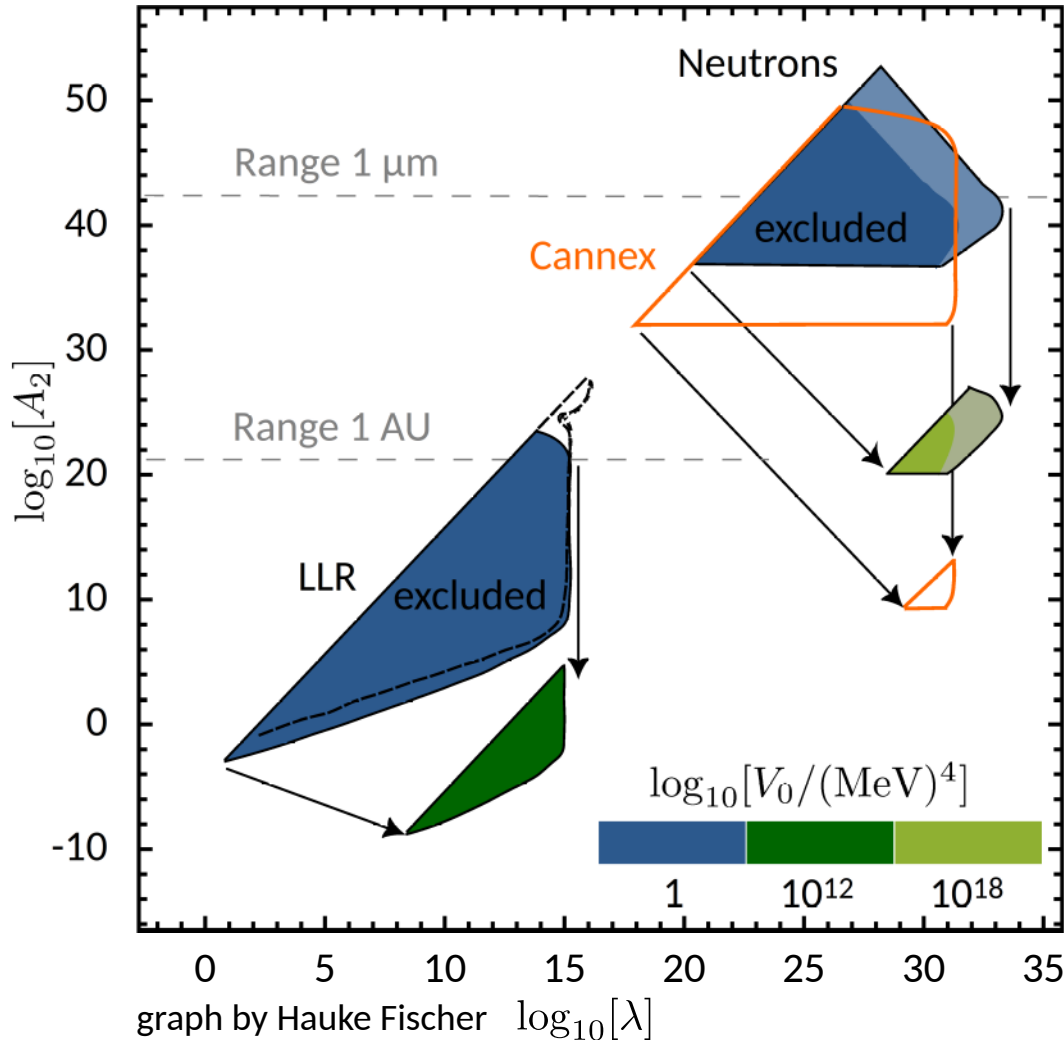
Measurement via vacuum pressure modulation

New limits possible

# What can be achieved?

## Limits on the Dilaton as dark energy candidate

Fischer *et al.* submitted.



Dynamic scalar field:

$$\mathcal{L} = \sqrt{-g} \left[ \frac{R}{16\pi G} - \frac{(\partial\phi)^2}{2} + V(\phi) \right] + \mathcal{L}_{\text{sm}}(\psi, \tilde{g}_{\mu\nu})$$

$$\tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu}$$

Coupling:  $A(\phi) = 1 + \frac{A_2\phi^2}{2m_{\text{pl}}^2}$

Potential:  $V(\phi) = V_0 e^{-\lambda\phi/m_{\text{pl}}}$

Effective potential:

$$V_{\text{eff}} = V_0 e^{-\lambda\phi/m_{\text{pl}}} + A_2 \frac{\phi^2}{2m_{\text{pl}}^2} \rho$$

Measurement via vacuum pressure modulation

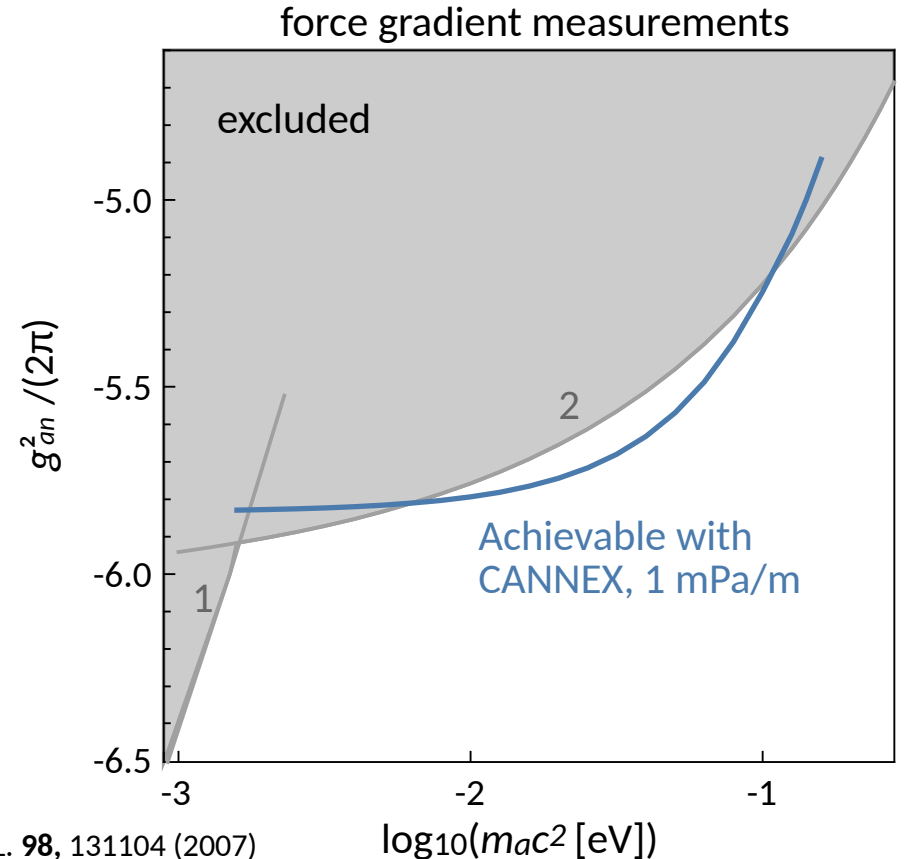
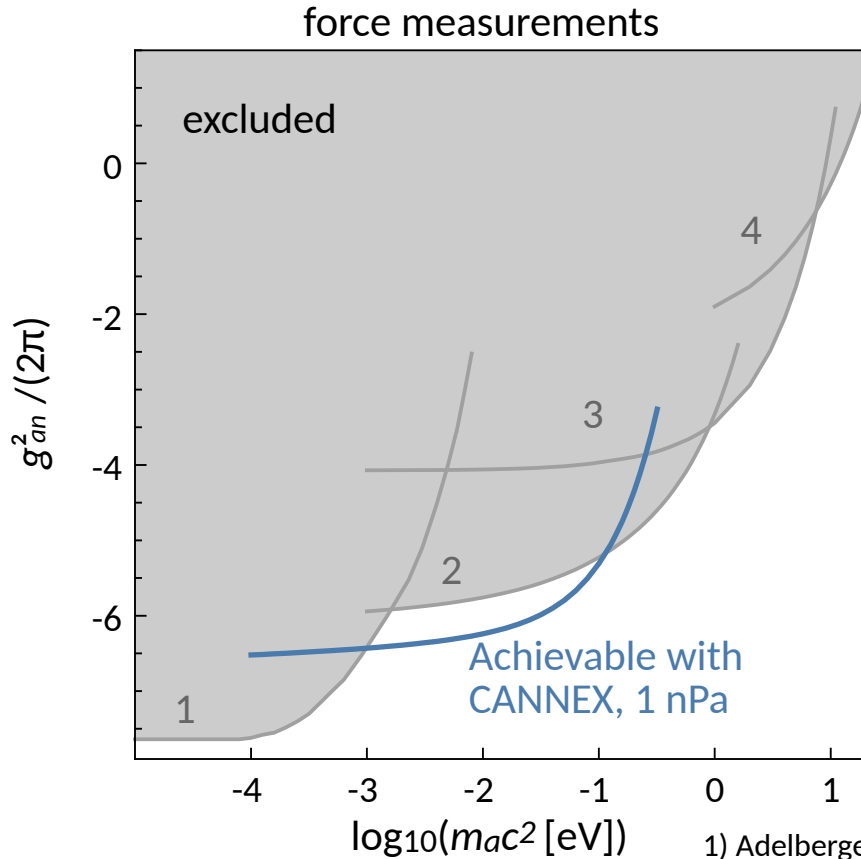
New limits possible

# What can be achieved?

## Improved limits on Axion-nucleon (scalar-pseudoscalar) interactions

G. Klimchitskaya, V.M. Mostepanenko, R.I.P. Sedmik, and H. Abele, *Symmetry*, 11(3) 407 (2019)

$$V_{an}(r) = -\frac{g_{an}^4}{32\pi^3} \frac{\hbar^2 m_a}{m^2} (\mathbf{r}_1 - \mathbf{r}_2)^{-2} K_1\left(\frac{2m_a c |\mathbf{r}_1 - \mathbf{r}_2|}{\hbar}\right).$$



- 1) Adelberger, *et al*, PRL. **98**, 131104 (2007)
- 2) Chen, *et al*, PRL. **116**, 221102 (2016)
- 3) Bezerra, *et al*, EPJC **74**, 2859 (2014)
- 4) Bezerra, *et al*, PRD **90**, 055013 (2014)

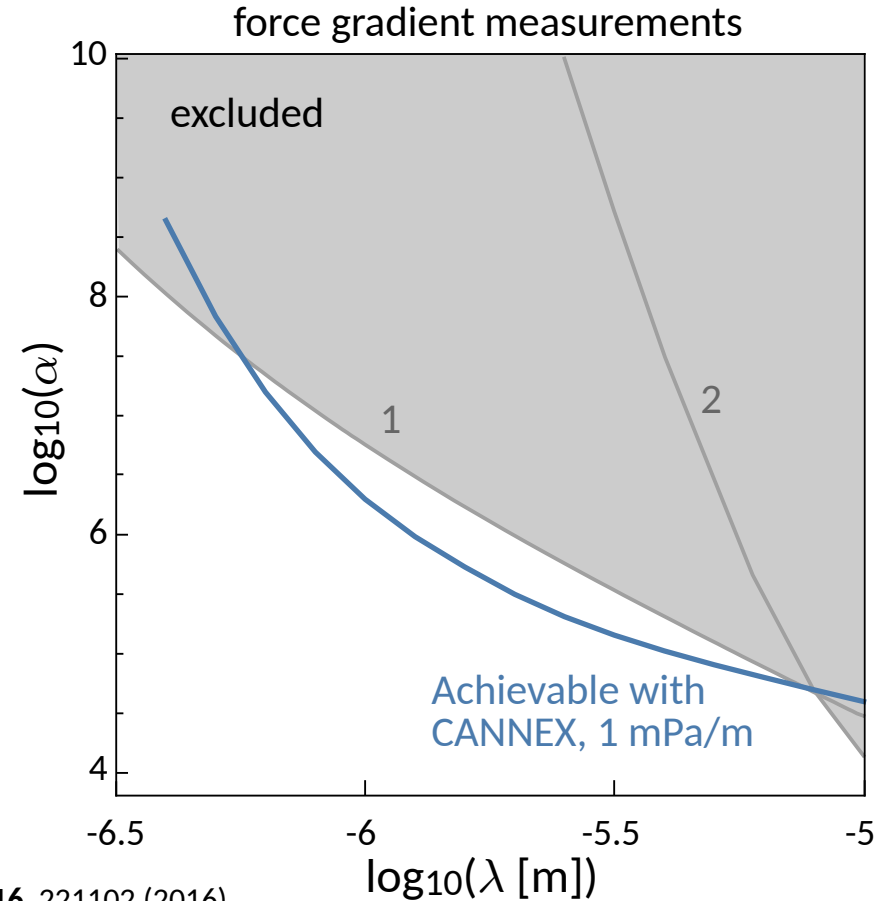
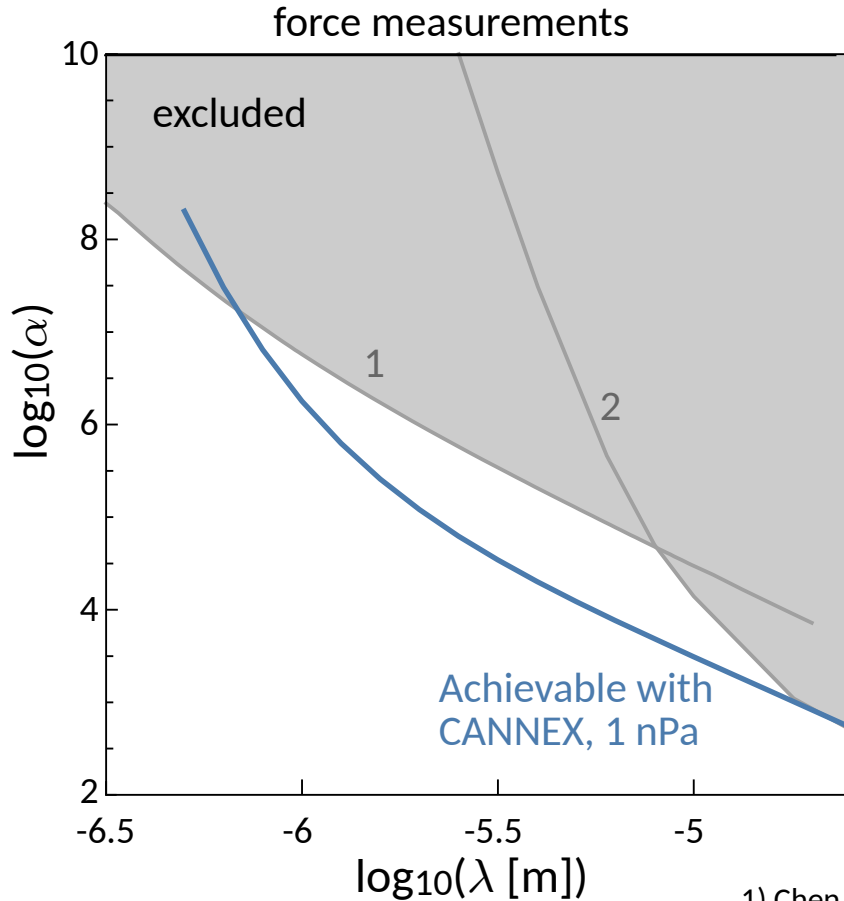
New limits achievable

# What can be achieved?

## Improved limits on Yukawa interactions

G. Klimchitskaya, V.M. Mostepanenko, R.I.P. Sedmik, and H. Abele, *Symmetry*, 11(3) 407 (2019)

$$V_{\text{Yu}}(r) = -\alpha \frac{Gm_1m_2}{|\mathbf{r}_1 - \mathbf{r}_2|} \exp\left(-\frac{(|\mathbf{r}_1 - \mathbf{r}_2|)}{\lambda}\right)$$



- 1) Chen, *et al*, PRL. **116**, 221102 (2016)
- 2) Geraci, *et al*, PRD **78**, 022002 (2008)

**New limits achievable**

# Conclusion

- CANNEX is about to become the **first metrological plane parallel plate Casimir experiment**
- The proof of principle was successful (though solvable technical problems encountered)
- Setup **is being rebuilt** with many enhancements
- First setup to **detect force and force gradient in one measurement**  
**Interfacial and Cavendish configuration possible.**
- CANNEX will allow to measure:
  - **Casimir forces at large separations ( $>10\ \mu\text{m}$ )**
  - **Thermal effects**
  - Dark matter and dark energy “fifth” forces
  - Gravity

# Acknowledgments

## Thank you for your your attention!

Present and former team members

ATI:

**René Sedmik**

Hamid Haghmoradi

Mario Pitschmann

Andrzej Pelczar

Ivica Galic

Roman Gergen

Thomas Dokulil

Hartmut Abele

VU (NL):

Attaallah Almasi

Joost C. Rosier

Lex v. d. Gracht

Davide Iannuzzi

Rogier Elsinga

Rob Limburg

Nikhef (NL):

Alessandro Bertolini

Eric Hennes

Arnold Rietmeijer

Johannes v. d. Brand

Berlin:

Francesco Intravaia

Paris Saclay:

Philippe Brax

St. Petersburg

Galina Klimchitskaya

Vladimir M. Mostepanenko



partially crowd-funded via **Flintwave** and

