

# Polycrystalline LPCVD 3C-SiC Thin Films on SiO<sub>2</sub> Using Alternating Supply Deposition

Philipp Moll<sup>1</sup>, Georg Pfusterschmied, Sabine Schwarz, Werner Artner, and Ulrich Schmid

**Abstract**—In this paper, we demonstrate the deposition of 3C-SiC thin films on SiO<sub>2</sub> using the alternating supply deposition (ASD) technique in a low-pressure chemical vapor deposition (LPCVD) furnace. We provide data of the thin film properties showing strong dependencies on the process gas flow rates of silane, propane and hydrogen. For comparative reasons all gas flow compositions were performed on <100> silicon and SiO<sub>2</sub>. A decreased rate of growth per cycle of ~37 % was discovered on SiO<sub>2</sub>. X-ray photoelectron spectroscopy (XPS) depth profiling revealed an oxygen content of 7.5 % ±2.5 % throughout the entire thin film when grown on SiO<sub>2</sub>. High resolution transmission electron microscopy (HRTEM) showed a 15 nm amorphous carbon layer at the 3C-SiC/Si interface. Conversely, on SiO<sub>2</sub> a 10 nm graphite layer was determined as intermediate layer leading to prominent <111> 3C-SiC X-ray diffraction (XRD) peaks. Independent of the substrate type a similar microstructure is observed in cross-sectional analyses. Atomic force microscopy (AFM) surface roughness measurements showed for all SiO<sub>2</sub> thin films lower values with a minimum of 4.9 nm (RMS), compared to 7 nm on Si. The electrical film resistivity was determined on SiO<sub>2</sub> with CTLM analysis, depending on the process gas composition. The gained knowledge is beneficial for MEMS applications, where tailored 3C-SiC-on-SiO<sub>2</sub> structures are desired. [2024-0114]

**Index Terms**—3C-SiC, alternating supply deposition, LPCVD, silicon carbide, SiO<sub>2</sub>.

## I. INTRODUCTION

**D**URING the early stages of CMOS scaling, the importance of engineered substrates gained significant interest. One key parameter was understanding the interface interaction between sublayers [1]. Without this innovation the rapid progress of doubling the possible numbers of transistors in an integrated circuit every two years, described by Moore's law, would not have been possible [2], [3]. In addition to the field of microelectronics, engineered substrates such as *silicon-on-insulator* (SOI) wafers are of utmost importance for microelectromechanical system (MEMS) devices, such

as resonators [4], [5], [6], [7], [8], [9], [10], [11], [12] or membrane-based architectures [13], [14], [15], [16]. However, silicon (Si) suffers from electrical instability at ambient temperatures of 250 °C for electronic devices [17], [18], [19]. SiC is known as a wide band gap material with 2.29 to 3.3 eV [20], [21] and features a high breakdown electric field strength of 3 – 4·10<sup>6</sup> V/cm compared to Si (2·10<sup>5</sup> V/cm) [17], making it suitable for high-power applications at elevated temperatures. Additionally, above temperatures of 500 °C Si loses its mechanical reliability for structural applications [22], [23]. Therefore, the approach of *silicon carbide on insulator* (SiCOI) [24], [25], [26], [27], [28] has been developed already 25 years ago. The improvement of the temperature related mechanical stability over 350 °C [19], [22], hardness of 20.9 GPa [29], [30], mechanical strength of 185 GPa [31], [32], [33] compared to Si and combined with the excellent chemical inertness [34], makes SiC an outstanding candidate for high-temperature applications. Although over 280 polytypes are known for SiC, only the cubic 3C-SiC can be synthesized at process temperatures below 1400 °C making it suitable for depositions on Si or SiO<sub>2</sub> composite substrates [35], [36]. As many MEMS devices are not restricted to single crystalline materials, polycrystalline 3C-SiC is regarded as a suitable material for structural applications.

In this paper we present results on polycrystalline 3C-SiC thin films deposited with LPCVD on thermally grown silicon dioxide (SiO<sub>2</sub>) on Si. Studies of 3C-SiC-on-SiO<sub>2</sub> can already be found in the literature, where traditional simultaneous supply depositions (SSD) were employed [37], [38], [39], [40], [41]. Nevertheless, we employed for the first time the alternating supply deposition (ASD) technique [42], [43], [44], [45], [46] on SiO<sub>2</sub> as a special type of pulsed CVD. We already showed in a previous study, that with ASD it is possible to specifically tailor the coefficient of thermal expansion (CTE) of 3C-SiC thin films in a range from 4.37 ppm/K to 13.96 ppm/K [47]. Additionally, ASD claims to yield 3C-SiC thin films featuring both lower full width at half maximum (FWHM) when evaluating X-ray diffraction (XRD) rocking curves and smoother surfaces [45]. This study will show, that we are able to fabricate 3C-SiC thin films on SiO<sub>2</sub> with tailored characteristics such as growth rate, surface roughness, grain size and electrical conductivity solely controlled with the process gas flow rates. We will provide knowledge on the growth mechanism on SiO<sub>2</sub> and discuss substantial differences and similarities of 3C-SiC films deposited in parallel on Si. To the best of the authors' knowledge, the deposition of

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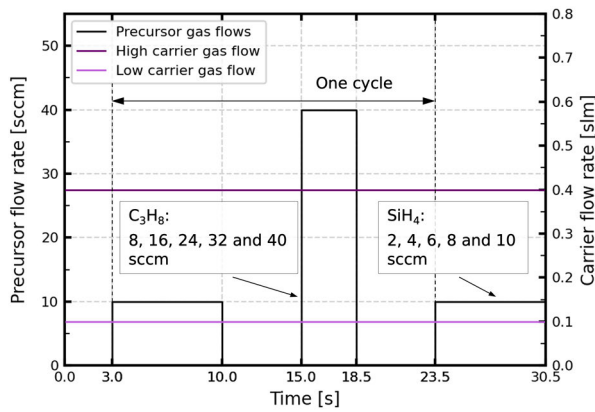


Fig. 1. Schematic representation of one full ASD deposition cycle. The flow rates of the precursors were set 2, 4, 6, 8 and 10 sccm for  $\text{SiH}_4$  and 8, 16, 24, 32 and 40 sccm for  $\text{C}_3\text{H}_8$ , whereas the ratio of  $\text{SiH}_4$  to  $\text{C}_3\text{H}_8$  was always set to 1/4.

3C-SiC thin films on  $\text{SiO}_2$  using ASD is not reported in the literature so far.

## II. EXPERIMENTAL DETAILS

As  $\text{SiO}_2$  substrates we used 4 “Si <100> wafers with a 250 nm LPCVD grown oxide on top. For reasons of comparison, we also used pure, n-doped 4 “Si <100> and <111> wafers with a bulk conductivity of  $50 \Omega\cdot\text{cm}$ . As LPCVD furnace we utilized an EasyTube 3000EXT industrial quartz oven from FirstNano. More details on the LPCVD equipment can be found elsewhere [48]. The basic deposition process consists of three parts. First, all samples were treated with a hydrogen ( $\text{H}_2$ ) *in situ* cleaning step at  $1000^\circ\text{C}$  with 5 slm for 10 minutes at atmospheric pressure. Subsequently, a tailored fast carbonization step was applied [49], [50], [51], [52] to overcome the lattice mismatch of  $\sim 20\%$  [35], [53] and passivate the Si surface against thermal hydrogen etching [54]. Hereby, the furnace ramps up from  $900^\circ\text{C}$  to  $1085^\circ\text{C}$  in 20 minutes with 1 slm of  $\text{H}_2$  and 100 sccm of propane ( $\text{C}_3\text{H}_8$ ). After 5.5 minutes at an ambient pressure of 20 torr the formation of the intermediate buffer layer was finished. The third step is the implementation of the ASD synthetization of polycrystalline 3C-SiC. To do so, the precursor gases were introduced successively into the reaction chamber separated by pump-out steps. For the silicon precursor we used pure silane ( $\text{SiH}_4$ ) and for the carbon precursor propane. To achieve a supportive flow regime  $\text{H}_2$  was implemented as carrier gas [34], [55]. Figure 1 shows a schematic of one representative ASD cycle. The precursor gases are introduced into the reaction chamber with a ratio of 1/4 for  $\text{SiH}_4/\text{C}_3\text{H}_8$  with flow rates of 2, 4, 6, 8, 10 sccm and 8, 16, 24, 32, 40 sccm for  $\text{SiH}_4$  and  $\text{C}_3\text{H}_8$ , respectively. For the carrier gas, flow rates of  $\Psi_{\text{low}} = 0.1$  slm and  $\Psi_{\text{high}} = 0.4$  slm of  $\text{H}_2$  were investigated.

During all depositions the temperature of  $1000^\circ\text{C}$  and the pressure of 0.32 torr were constant. In a first step we determined the growth per cycle on all substrates to achieve in a second step comparable layer thicknesses of  $400 \text{ nm} \pm 5\%$  for thin film characterization. To investigate the growth behavior, we deposited on each substrate type thin films with 10, 30, 60, 90 and 120 cycles and process gas flows of

10 sccm for  $\text{SiH}_4$  and 40 sccm for  $\text{C}_3\text{H}_8$  with  $\Psi_{\text{high}}$ . For better readability and due to the frequent occurrence of this gas flow combination, the parameter  $\Psi_A$  is introduced. The layer thickness was optically measured utilizing a Filmetrics F20-UVX thin film analyzer within a wavelength range from 600 nm to 1600 nm. For cross-sectional scanning electron measurements (SEM) recordings a Hitachi SU8030 scanning electron microscope with an acceleration voltage of 4 kV was used. For chemical analysis the samples were investigated with X-ray photoelectron spectroscopy (XPS) depth profiling, utilizing a SPECS XP-spectrometer equipped with a monochromatic Al-K $\alpha$  X-ray source and a hemispherical WAL-150 analyzer. The crystallographic orientation of the thin films was determined with X-ray diffraction (XRD) measurements with a Malvern PANalytical X’pert PRO X-ray powder diffractometer. The measurements were taken at 45 kV and 40 mA with a copper K $\alpha$  radiation source. Diffraction patterns in Bragg-Brentano configuration were performed from  $2\theta = 15^\circ$  to  $90^\circ$  and rocking curves were taken in the range from  $34^\circ$  to  $38^\circ$ , since the most prominent 3C-SiC peak was expected at around  $36^\circ$  [39]. To quantitatively assess the preferred crystallographic growth orientation, we employed the texture coefficient (TC) calculation using the Harris method outlined in [56]

$$TC_{(hkl)} = \frac{I_m(hkl)/I_0(hkl)}{\left(\frac{1}{n}\right) \sum [I_m(hkl)/I_0(hkl)]}. \quad (1)$$

Here, the term  $(hkl)$  refers to the specific growth plane under examination,  $I_m$  denotes the normalized measured intensity with the background intensity subtracted,  $I_0$  represents the intensity of a randomly oriented polycrystalline sample sourced from a Powder Diffraction File (PDF) card and  $n$  gives the number of planes analyzed. The PDF card utilized in this study for 3C-SiC was #00-029-1129 and the crystallographic planes investigated for TC calculation were (111), (200), (220) and (311) resulting in  $n = 4$ .

We performed AFM measurements on sample areas of  $3 \times 3 \mu\text{m}^2$  to determine the surface roughness with a Bruker Dimension Edge atomic force microscope in tapping mode, with a NHCV-tip (spring constant  $\sim 40$  N/m). HRTEM and electron diffraction investigation on the Si/3C-SiC and  $\text{SiO}_2/3\text{C-SiC}$  interfaces were performed with a TECNAI F20 transmission electron microscope operating with an acceleration voltage of 200 kV.

For the electrical characterization of the unintentionally doped 3C-SiC thin films circular transmission line method (CTLM) analysis were performed to determine the film resistivity  $\rho$  and the specific contact resistance  $\rho_C$ , according to the model of Klootwijk and Timmering [57]. I/V (current/voltage)-curves for each gas flow composition were recorded, whereas the  $\text{SiO}_2$  served as an electrical insulating layer to the Si substrate. The contact pads were fabricated with a reverse sputter lift off process consisting of 15 nm of titanium (Ti) and 100 nm of tantalum (Ta) with no post deposition annealing. Following [57], ten contact pads per sample were designed with the shape of two concentric circles forming a center electrode with a radius of  $200 \mu\text{m}$  and an accompanying outer electrode with an inner radius ranging from 220 to  $400 \mu\text{m}$ .

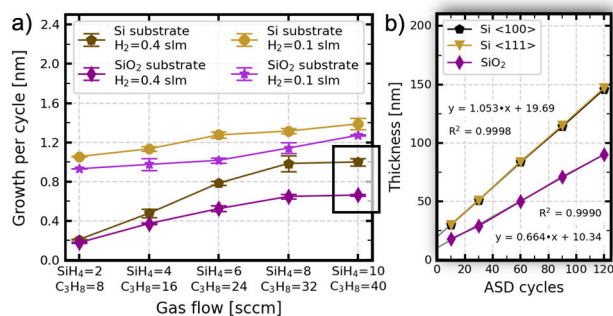


Fig. 2. Growth per cycle of polycrystalline 3C-SiC thin films deposited with increasing precursor gas supply and two different carrier gas flows on Si and SiO<sub>2</sub> substrates can be seen in a). The layer thickness was measured in the middle of the sample. Depositions of five  $\Psi_A$  thin films with 10, 30, 60, 90 and 120 cycles on <100> Si, <111> Si and SiO<sub>2</sub> are presented in b). A linear increase in layer thickness on all substrates can be seen, whereas the growth per cycle on both silicon substrates was almost identical, while the growth rate for SiO<sub>2</sub>, however, was significant lower.

The I/V curves for the CTLM analysis were taken with a Süss PM8 waferprober and an Agilent B2911A Precision Source/Measure Unit from  $-3$  V to  $3$  V in 61 steps.

### III. RESULTS

To start, the growth per cycle for different process gas flow rates is presented in Figure 2 a). Regardless of the substrate type the growth per cycle was lower for  $\Psi_{\text{high}}$ . This finding originates from the nature of the ASD deposition scheme, where at each cycle first an ultra-thin silicon layer is deposited, followed by a carbon-assisted silicon redistribution, thus forming 3C-SiC [48]. During deposition this thin silicon layer is attacked by the hydrogen gas, hence reducing the growth per cycle, when using higher H<sub>2</sub> flow rates. This influence is well-known in silicon technology as hydrogen thermal etching [54]. Simultaneously, hydrogen passivation [58] and inhibition-effect [59], [60], [61] takes place, both contributing in a decreased growth per cycle. This entails the saturation of surface defects or Si dangling bonds with hydrogen atoms and interactions with the Si surface, thus reducing the reactivity of the Si and inhibit the adsorption and incorporation of C precursors. This stands in contrast with SSD depositions where increased growth rates are typically observed with higher H<sub>2</sub> flow rates. For  $\Psi_{\text{low}}$  depositions a linear dependency as a function of the precursor flow rates was measured. For  $\Psi_{\text{high}}$  depositions, however, when exceeding a threshold of 8 sccm for SiH<sub>4</sub> and 32 sccm for C<sub>3</sub>H<sub>8</sub> the growth per cycle saturated. Since the growth per cycle did not further increase by providing even more precursor gases this saturation effect is explained by a reaction limited process [62].

The main difference visible in Figure 2 a) is that for the same deposition conditions the growth per cycle was always lower on SiO<sub>2</sub>, than on Si. Very little can be found in the literature, when it comes to LPCVD growth rates on SiO<sub>2</sub> compared to Si [37]. To clarify the origin of the difference in growth per cycle we present the thickness of different depositions with increasing number of ASD cycles on <100> Si, <111> Si and SiO<sub>2</sub>, which is illustrated in Figure 2 b). The two Si samples showed identical results and a highly linear

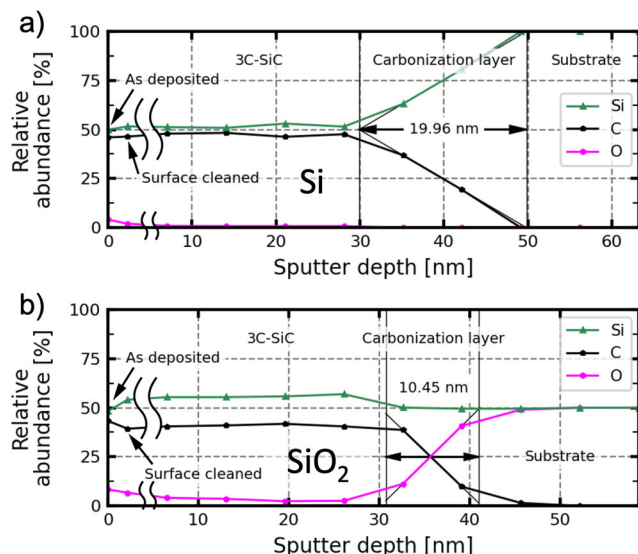


Fig. 3. XPS depth profiling showing the relative presence of Si, C and O through a 3C-SiC cross section. In a) the element analysis of the 3C-SiC/Si interface can be seen. With a depth resolution of  $\sim 7$  nm, a carbonization layer thickness of  $\sim 20$  nm was measured. The depth profile of b) gives the chemical composition around the 3C-SiC/SiO<sub>2</sub> interface. The main differences between both measurements are a smaller carbonization layer thickness of  $\sim 10$  nm (with  $\sim 6.5$  nm resolution) and an oxygen proportion of 7.5 % through the whole thin film.

growth behavior with a measured growth per cycle rate of 1.05 nm. We also found a highly linear growth behavior for SiO<sub>2</sub>, however, with a lower growth per cycle of 0.66 nm. Both rates correspond to what is presented in Figure 2 a) for the  $\Psi_A$  deposition (black box). Additionally, the two calculated intercepts with the y-axis represents the thickness of the carbonization layer resulting in 19.69 nm and 10.34 nm for Si and SiO<sub>2</sub>, respectively. Two hypotheses seem reasonable for an inhibited growth behavior on SiO<sub>2</sub>: first, the amorphous seed layer is disadvantageous for 3C-SiC crystal growth. This can be motivated, since amorphous substrates lack of a crystalline structure, which can hinder the nucleation and growth of 3C-SiC compared to single crystal Si, thus resulting in reduced growth rates and lower thin film quality. Second, oxygen effusion from the SiO<sub>2</sub> surface prevents the crystal growth compared to Si.

To further investigate this hypothesis, we performed XPS depth-profiling. Two representative measurements of the same  $\Psi_A$  deposition on Si and SiO<sub>2</sub> are shown in Figure 3. All thin films grown on SiO<sub>2</sub> exhibit an oxygen content of  $7.5 \pm 2.5$  % through the whole thin film, compared to  $<1$  % on Si. The homogeneous distribution of oxygen throughout the entire thin film can be explained by vacancy diffusion [63], [64]. Eickhoff et al. [37] showed, that at elevated temperatures above 1150 °C SiO<sub>2</sub> substrates reach up to 100 % selectivity during 3C-SiC CVD depositions compared to pure Si surfaces. It is also shown that for lower temperatures the growth rate is decreased compared to Si.

This indicates that the oxygen inhibits the growth of 3C-SiC. Even more, Nagasawa and Yagi state, that in their studies even the thinnest SiO<sub>2</sub> layers prevent the surface from chemically adsorbing the carbon precursor [65]. We propose, that the

oxygen present within the thin films partially forms  $\text{SiO}_2$  during each  $\text{SiH}_4$  ASD sequence, thus inhibiting the carbon to react with the silicon and hence decreasing the growth per cycle. This assumption is also in line with the lower thickness of the carbonization layer on  $\text{SiO}_2$  since the oxygen lowers the adsorption of the carbon during the carbonization step.

Thus, only a  $\sim 10$  nm thin carbon layer is deposited on the surface. Furthermore, XPS measurements revealed, that all thin films deposited on Si are confirmed to be stoichiometric, whereas thin films deposited on  $\text{SiO}_2$  turned out to be non-stoichiometric with a higher percentage of Si content.

Figure 4 shows the XRD spectra of the different 3C-SiC films on both substrates, whereas the most prominent peak was measured at  $35.78^\circ$ , which can be attributed to the  $\langle 111 \rangle$  direction. With the other peaks detected at  $41.48^\circ$  and  $60.1^\circ$ , which can be assigned to the  $\langle 200 \rangle$  and  $\langle 220 \rangle$  directions, we conclude that the thin films consist only of poly 3C-SiC. One representative rocking curve for each substrate is shown in the inserts of Figure 4, with the lowest full width at half maximum (FWHM) values at  $1.112^\circ$  and  $0.446^\circ$  for Si and  $\text{SiO}_2$ , respectively. For both substrates the  $\Psi_A$  gas flow configuration resulted in the lowest FWHM results. Typical values for comparable poly 3C-SiC thin films on Si are in the range of  $1.4^\circ$  to  $1^\circ$  [43]. For  $\text{SiO}_2$ , however, no similar study could be found to compare our results.

The texture coefficient, presented in Figure 4 b) and d), revealed only for  $\Psi_{\text{high}}$  depositions on  $\text{SiO}_2$  a preferred growth orientation in the (111) plane. Other gas flow combinations did not exhibit a preferred growth orientation. The reason for the prominent formation of  $\langle 111 \rangle$  3C-SiC on the  $\text{SiO}_2$  is because of the low free surface energy of  $\langle 111 \rangle$  facets, which tend to form  $\langle 111 \rangle$  3C-SiC even on amorphous surfaces [39]. This phenomenon is well-documented in the literature for 3C-SiC depositions [40], [66], [67], [68], despite the fact that the  $\langle 110 \rangle$  direction exhibits an even lower surface energy of  $3.4 \text{ J/m}^2$  compared to the  $4.2 \text{ J/m}^2$  of the  $\langle 111 \rangle$  direction [69]. On single crystal substrates the deposited thin films tend to grow in the same direction as the seed crystal [67]. It is also reported, that if the growth direction of the thin film deviates from the substrate orientation, two explanations can be given. First, because of a higher density of aligned atoms per unit, forming a more stable thin film [70]. For example, the density of aligned atoms per unit of  $\langle 111 \rangle$  3C-SiC on  $\langle 110 \rangle$  Si is 4 times higher than  $\langle 110 \rangle$  3C-SiC on  $\langle 110 \rangle$  Si, thus forming a thin film, where the  $\langle 111 \rangle$  direction is preferred [67], [70]. The second opinion proposed by Zheng et al. is the impact of the carbonization layer [67].

For better understanding of the XRD results we took HRTEM recordings of the 3C-SiC/Si and 3C-SiC/ $\text{SiO}_2$  interfaces, presented in Figure 5 a) and b). Electron diffraction patterns of the two substrates and the corresponding 3C-SiC thin films are presented in the section e) – h) of Figure 5. The Si diffraction pattern shows a single crystal pattern with a lattice constant of  $d_{111} = 3.123 \text{ \AA}$  almost matching the lattice constant of a stress-free Si crystal with  $d_{111} = 3.136 \text{ \AA}$  [53]. As expected, an amorphous pattern was recorded for the  $\text{SiO}_2$  substrate. The diffraction patterns for the thin films confirm a polycrystalline microstructure in the investigated regions. Thin

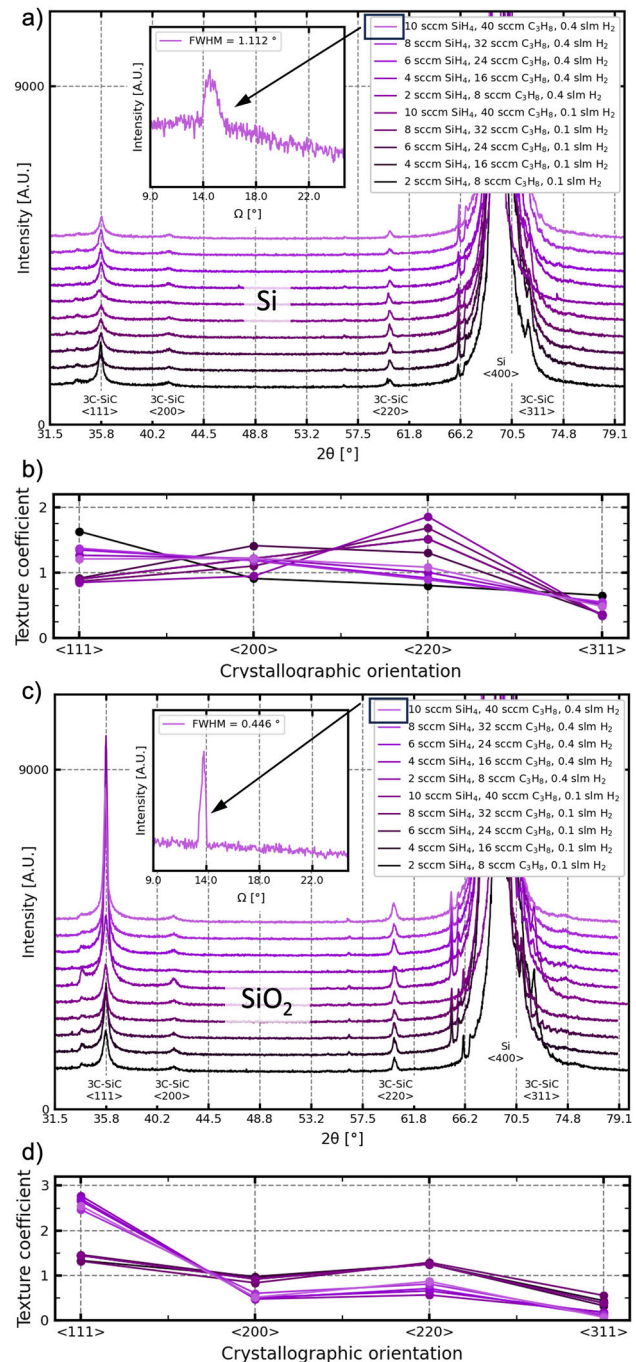


Fig. 4. XRD measurements of the 3C-SiC thin films investigated in this study, while a) gives the results on Si [48], and c) those on  $\text{SiO}_2$ . Most prominent in both diagrams is the  $\langle 111 \rangle$  of 3C-SiC at  $35.79^\circ$ . Both, in a) and c) the inserts show the rocking curves with the lowest FWHM values. The TC was calculated for each crystallographic plane, illustrated for Si in b) and  $\text{SiO}_2$  in d).

films grown on Si, however, show streaks at the diffraction spots, indicating higher strain levels, compared to those grown on  $\text{SiO}_2$ .

Electron energy loss spectroscopy (EELS) measurements [48] of the 3C-SiC/Si interface revealed that the carbonization layer consists of a 3 – 6 nm carbonized Si section, followed by a 10 – 15 nm amorphous carbon layer, on top the 3C-SiC thin film. This means that the thin films grow on an amorphous carbon seed layer when using Si as

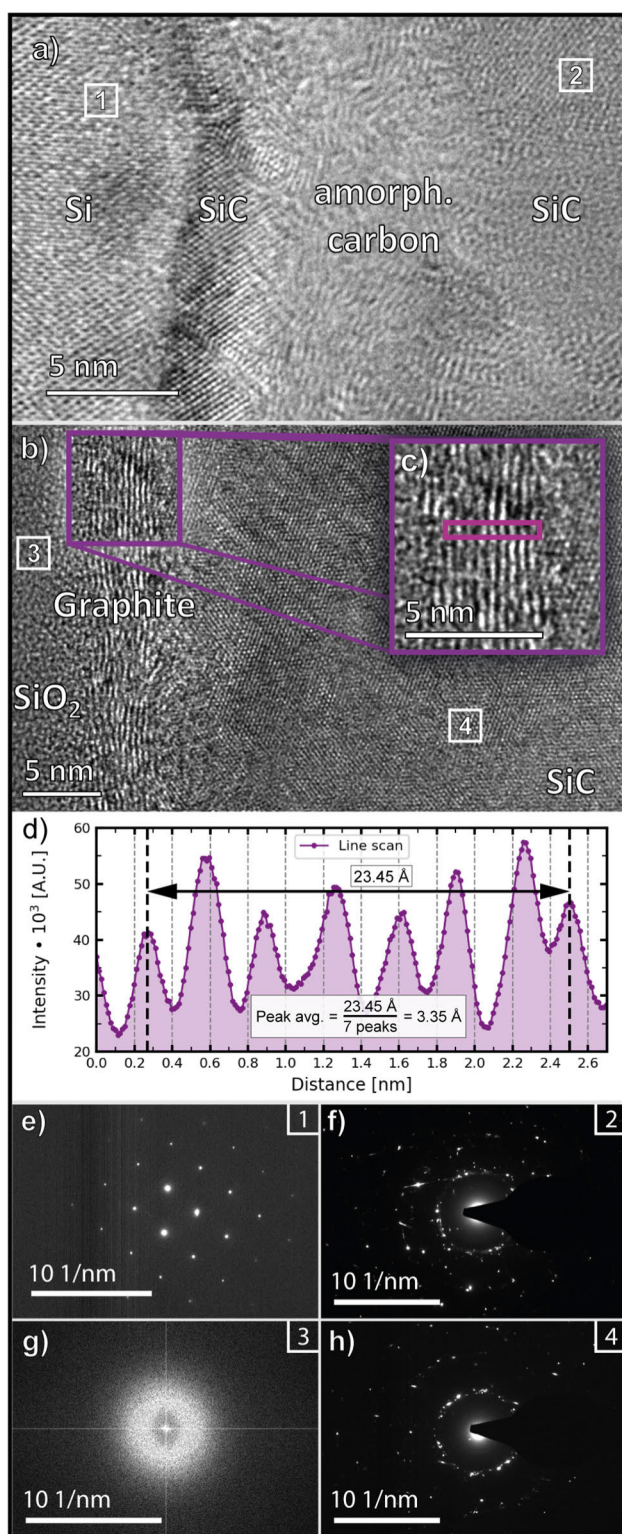


Fig. 5. In a) a HRTEM of the carbonization layer on Si can be seen. For comparison the carbonized interface of the SiO<sub>2</sub> substrate with 3C-SiC on top is presented in b) and with higher magnification in c). A 2.7 nm long intensity profile of the graphitic layer is depicted in d). The position, where the profile was conducted is highlighted in Figure 5 c). The electron diffraction patterns from d) – g) depict the regions of pure Si, 3C-SiC grown on Si, pure SiO<sub>2</sub> and 3C-SiC grown on SiO<sub>2</sub>, respectively.

substrate, explaining the most prominently  $\langle 111 \rangle$  crystallographic orientation of the poly 3C-SiC. The resemblance of the XRD results on both Si and SiO<sub>2</sub> substrates is explained

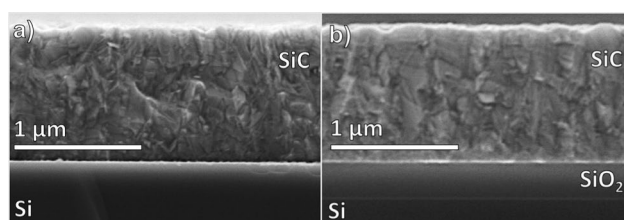


Fig. 6. Similar 3C-SiC thin film microstructures on Si and SiO<sub>2</sub> visualized by cross-sectional SEM recordings as illustrated in a) and b). For reasons of comparison a desired layer thickness of 1 μm was achieved with 1050 and 1500 cycles for Si and SiO<sub>2</sub>, respectively.

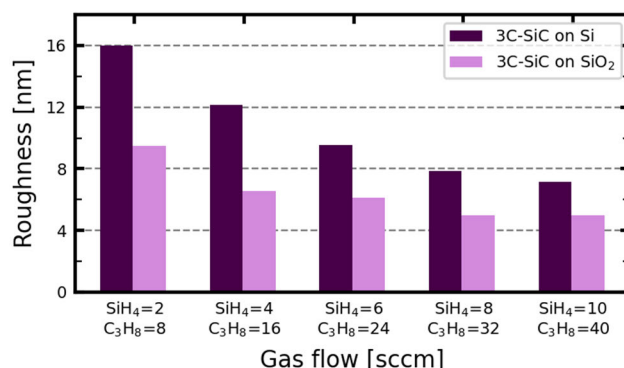


Fig. 7. Comparison of the surface roughness of representative  $\Psi_{\text{high}}$  thin films with a thickness of 400 nm ± 5% deposited on Si and SiO<sub>2</sub> as a function of the precursor flow rate.

by the excess carbon at the interface originated from the carbonization step. The smaller FWHM results from a crystalline graphite carbonization layer. This graphitic layer was confirmed by an intensity profile of the HRTEM, as highlighted in Figure 5 c). The lattice plane distance  $d_{200}$  was determined with 3.35 Å ± 0.03 Å, as shown in Figure 5 d). This result is in perfect agreement with similar studies [71] and to what is provided for graphite lattice distances in XRD studies [72], [73]. On Si, however, the carbonization step resulted in an amorphous carbon layer, thus providing less vital conditions for the deposition of high-quality 3C-SiC thin films. It can be concluded that the choice of an individual carbonization step for each substrate is crucial. The values for the carbonization layer thickness determined from the growth rate (see Figure 2 b) and the measured thickness of the XPS depth-profiling are in good agreement with the corresponding value from the HRTEM analysis. Even more, it is in good agreement to data in the literature [45], [49], [74].

The SEM analysis of the 3C-SiC thin films on Si and SiO<sub>2</sub>, given in Figure 6 a) and b), show comparable cross-sections. This can be explained by the excess carbon from the carbonization step, resulting in similar microstructures for all 3C-SiC thin films independent of the substrate type.

The results from AFM surface roughness measurements were set in relation with the corresponding layer thickness and are illustrated in Figure 7. At a comparable film thickness, a smoother surface topography is measured on all SiO<sub>2</sub> samples compared to Si. For both substrates a decreasing trend was measured for the RMS roughness, whereas independent from the carrier flow rate, smoothest surfaces can be contributed to

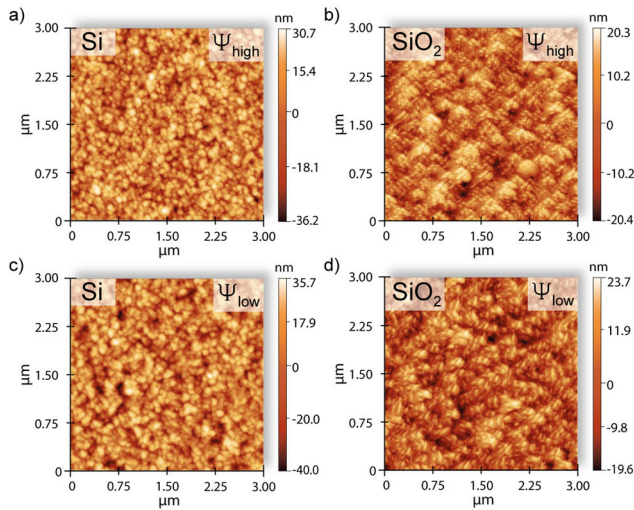


Fig. 8. Representative AFM pictures taken from depositions with 10 sccm for SiH<sub>4</sub> and 40 sccm for C<sub>3</sub>H<sub>8</sub> with  $\Psi_{\text{high}}$  can be seen in a) and b) and with  $\Psi_{\text{low}}$  in c) and d). The left column shows typical spherical shaped grain patterns on Si, whereas in the right column different results on SiO<sub>2</sub> can be seen. (Note: the pictures for Si and SiO<sub>2</sub> were taken consecutively for every precursor and carrier gas flow to avoid any doubts in the correctness of the measurements.)

the highest precursor flow rates. This can be explained by the inhibiting hydrogen effects, influencing not only the growth rate, but also affecting the surface roughness by extensive hydrogen exposure. Lowest values for the RMS roughness of 4.9 nm were measured on SiO<sub>2</sub> at a 3C-SiC film thickness of about 400 nm. In literature, most values for SiO<sub>2</sub> can be found around 20 nm for the RMS roughness for thicknesses  $>2 \mu\text{m}$ , whereby also mirror-like surfaces with 3 – 5 nm can be achieved with special post-deposition treatments (*e.g.* polishing) [39], [40], [41]. In addition, we achieved for thin films deposited on Si, mirror-like surfaces [68] with an RMS roughness below 7 nm. For 3C-SiC thin films on Si substrates a broad range is reported for this parameter in the literature ranging from 3 nm at  $\sim 400 \text{ nm}$  [43] up to 60 nm, strongly dependent on the layer thickness, the process gases, pressures or deposition regimes [43], [68], [75]. Not only smoother surfaces were measured on SiO<sub>2</sub>, also different patterns for the surface topography were recorded by AFM measurements, as shown in Figure 8. For the Si depositions we achieved typical patterns for the grain structure, comparable to what can be found in the literature [39], [68], [76]. On SiO<sub>2</sub>, however, the ASD thin films revealed a completely different topology. The 3C-SiC-on-SiO<sub>2</sub> topology of this study cannot be compared to SSD depositions on SiO<sub>2</sub>, where the grains exhibit a spherical shape and are indistinguishable from those on Si substrates [39], [77], [78], [79]. We propose that the variance between the two topographies of Si and SiO<sub>2</sub> depositions is also a result of the oxygen within the SiO<sub>2</sub> thin films.

Typical I/V measurements, as illustrated in Figure 9, provide an overview on the electrical properties of the investigated undoped 3C-SiC thin films as a function of the precursor flow rate. The Ti/Ta contact pads show perfect ohmic behavior within the investigated voltage range without the need for an annealing step.

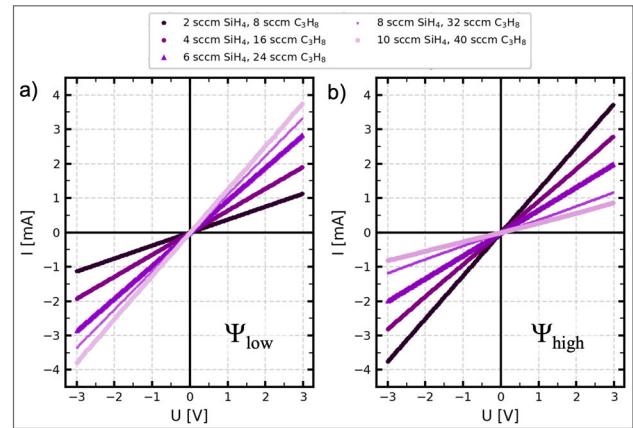


Fig. 9. I/V curves from undoped poly 3C-SiC thin films depending on the process flow rates. The I/V curves were measured on the contact pad with a spacing of 200  $\mu\text{m}$  between the electrodes.

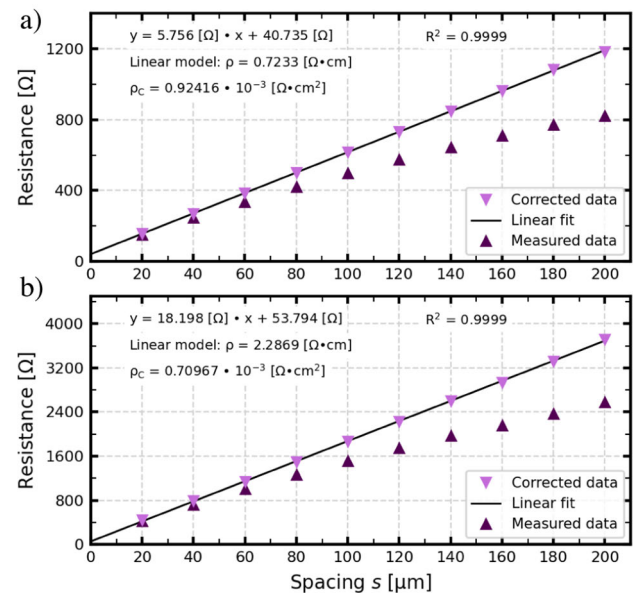


Fig. 10. Representative CTLM measurements for the lowest and highest total resistance, presented in a) and b) corresponding to the flow rates of 2 sccm SiH<sub>4</sub> and 8 sccm C<sub>3</sub>H<sub>8</sub>, and 10 sccm SiH<sub>4</sub> and 40 sccm C<sub>3</sub>H<sub>8</sub>, both deposited with  $\Psi_{\text{high}}$ .

Two opposing trends can be seen for each of the carrier flow rate. While for  $\Psi_{\text{low}}$  the total resistance  $R$  decreases with increasing precursor flow rate, while for  $\Psi_{\text{high}}$   $R$  decreases with decreasing precursor flow rate. Two examples, representing the highest and lowest resistance, are employed to illustrate the evaluation of the CTLM analysis, presented in Figure 10. The total resistances measured for each CTLM contact pad are plotted as a function of the corresponding spacings and a linear approximation is introduced, using a geometrical correction factor [57]. From the slope of the curve, the dimensions of the contacts and the film thickness, the film resistivity can be determined. The contact resistance  $R_c$  and the specific contact resistance  $\rho_c$  are determined from the intersections of the linear approximation with the x- and y-axes, respectively.

The results from the CTLM analysis are illustrated in Figure 11. For  $\Psi_{\text{low}}$  depositions decreasing resistivities with increasing precursor flow rates from 2.11 to 0.74  $\Omega \cdot \text{cm}$  were measured. The resistivities for  $\Psi_{\text{high}}$  were in the same range

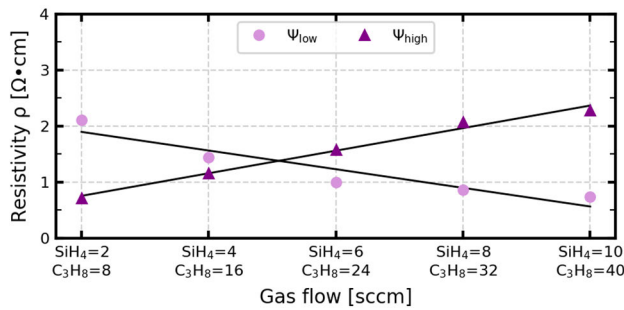


Fig. 11. Film resistivity of undoped 3C-SiC thin films as a function of the process gas flow rates, derived from CTLM analysis.

from 2.28 to 0.72  $\Omega\cdot\text{cm}$ , however, as a result of decreasing precursor flow rates. Strongly depending on the carbonization step, resistivities for undoped poly 3C-SiC thin films in the range of about 1 up to 6600  $\Omega\cdot\text{cm}$  can be achieved [1]. The contact resistances  $R_C$  showed no dependency on the gas flow rates and were in the range of 15 to 27  $\Omega$ , resulting in specific contact resistances  $\rho_C$  ranging from  $7.1\cdot 10^{-4}$  to  $1.27\cdot 10^{-3}$   $\Omega\cdot\text{cm}^2$ .

Comparable undoped 3C-SiC thin films exhibit values for  $\rho_C$  in the range of  $5\cdot 10^{-4}$  to  $1.2\cdot 10^{-5}$   $\Omega\cdot\text{cm}^2$  [80], [81], [82]. In general, electron scattering at grain boundaries can be named as the main phenomenon influencing the resistivity in unintentionally doped polycrystalline thin films [83], together with the incorporation of nitrogen atoms [84]. Taking this into account and the characteristic of an ideal SiC bulk crystal being electrically insulating at room temperature, it can be concluded that the two thin films showing the highest film resistivity ( $\Psi_{\text{low}}$ : 2 sccm for SiH<sub>4</sub>, 8 sccm for C<sub>3</sub>H<sub>8</sub> and  $\Psi_{\text{high}}$ : 10 sccm for SiH<sub>4</sub>, 40 sccm for C<sub>3</sub>H<sub>8</sub>), exhibit lowest impurity levels and lowest grain boundary density. This is in good agreement with the XRD results, where the two abovementioned gas flow compositions achieved lowest FWHM values. Both thin films originated from a balanced carrier to precursor gas flow ratio (low carrier flow rate – lowest precursor flow rate and high carrier flow rate – highest precursor flow rate), whereby the resistance as well as the FWHM of the other thin films increased the more unbalanced this ratio becomes. It is worth noting that the grow-inhibiting hydrogen effects dominate the thin film characteristics, thus influencing the film resistivity depending on the process gas flows.

#### IV. CONCLUSION

In this study we demonstrated the deposition of 3C-SiC on SiO<sub>2</sub> using ASD in a LPCVD system. By comparing the growth per cycle of the same depositions on Si we measured an inhibited growth behavior on SiO<sub>2</sub>. XPS depth profiling revealed an increased oxygen content of 7.5 % throughout the entire thin film grown on SiO<sub>2</sub>. The elevated oxygen level could potentially explain the decreased growth per cycle. On both substrates typical poly 3C-SiC XRD patterns were measured, whereas the  $\langle 111 \rangle$  peak appeared as the most prominent crystallographic orientation. Hereby  $\Psi_{\text{high}}$  depositions on SiO<sub>2</sub> exhibited a preferred growth orientation

in the  $\langle 111 \rangle$  direction. The HRTEM analysis revealed an amorphous carbon layer on top of the carbonized Si substrate. On the SiO<sub>2</sub> interface, however, a graphite layer was found, providing an explanation for lower FWHM values for the 3C-SiC thin films. This indicates the outstanding importance of tailored carbonization steps for each substrate. Nevertheless, the excess carbon at both interfaces provides comparable growth conditions, therefore explaining the highly  $\langle 111 \rangle$  textured thin films with similar microstructure verified with cross-section analyses on both substrates. From AFM measurements we found smoother surfaces for SiO<sub>2</sub>, with a minimum of 4.9 nm for the RMS roughness, whereas also different surface topologies were recorded. For both substrates the RMS roughness decreased with increasing precursor flow rate.

CTLM measurements showed the importance of the ratio from carrier to precursor gas flow rates. Hereby,  $\Psi_{\text{low}}$  paired with lowest precursor flow rates of 2 sccm for SiH<sub>4</sub> and 8 sccm for C<sub>3</sub>H<sub>8</sub> as well as  $\Psi_{\text{high}}$  with 10 sccm SiH<sub>4</sub> and 40 sccm C<sub>3</sub>H<sub>8</sub> achieved highest film resistivities of  $\sim 2.3$   $\Omega\cdot\text{cm}$ , giving strong indication of a lower impurity as well as grain boundary density values compared to thin films deposited with other gas flow compositions. This is consistent with the XRD rocking curve results, whereas lowest FWHM values were achieved by the above-mentioned well-balanced ratios of the process gas flow rates. The results in this study provide knowledge to achieve tailored thin film properties, enabling customized fabrication of improved SiCOI structures for MEMS applications.

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