

Strongly Interacting Polaritons in Moiré Quantum Materials

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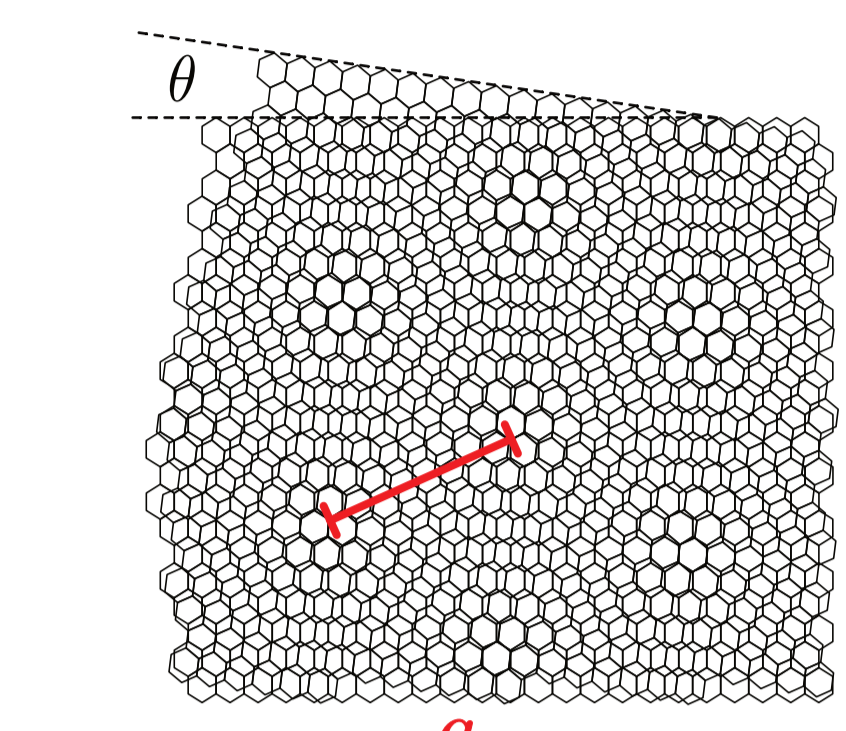
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Introduction - Moiré Excitons in a Nutshell

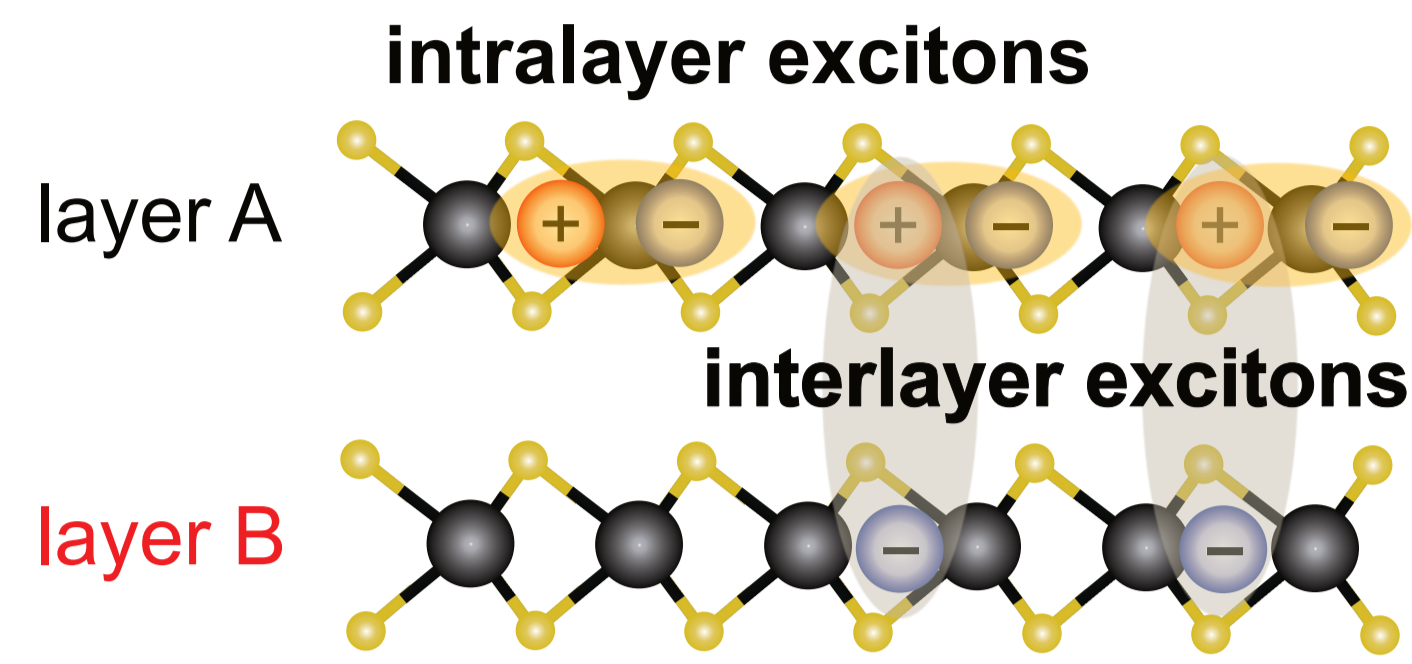
Quantum Materials with a Twist

- Twisting two layers of a two-dimensional quantum material leads to an effective moiré potential where excitons localise at the minima



$$a_M \approx \frac{a}{\sqrt{\delta^2 + \theta^2}} \gg a$$

lattice constant mismatch twist angle



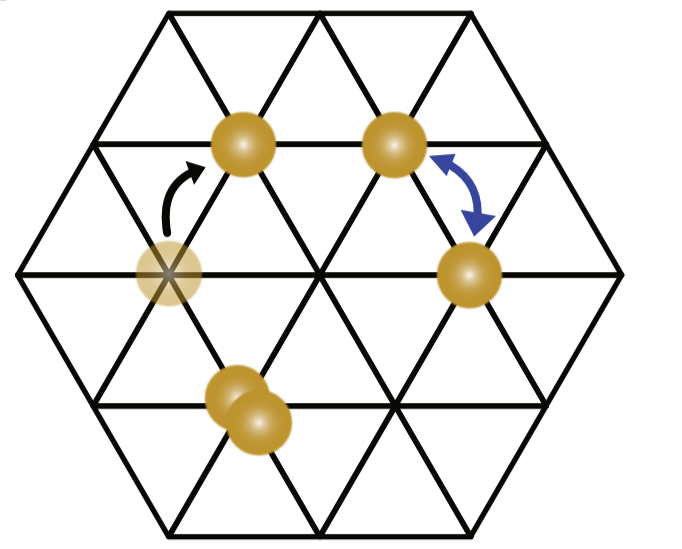
- intralayer excitons:** strong light-matter coupling, short lifetimes (\sim ps)
- interlayer excitons:** indirect excitons (long lifetime \sim ns), weak light-matter coupling, strong dipole-dipole interactions (\sim meV)

$$U\tau_{IX} \sim 10^3 - 10^4$$

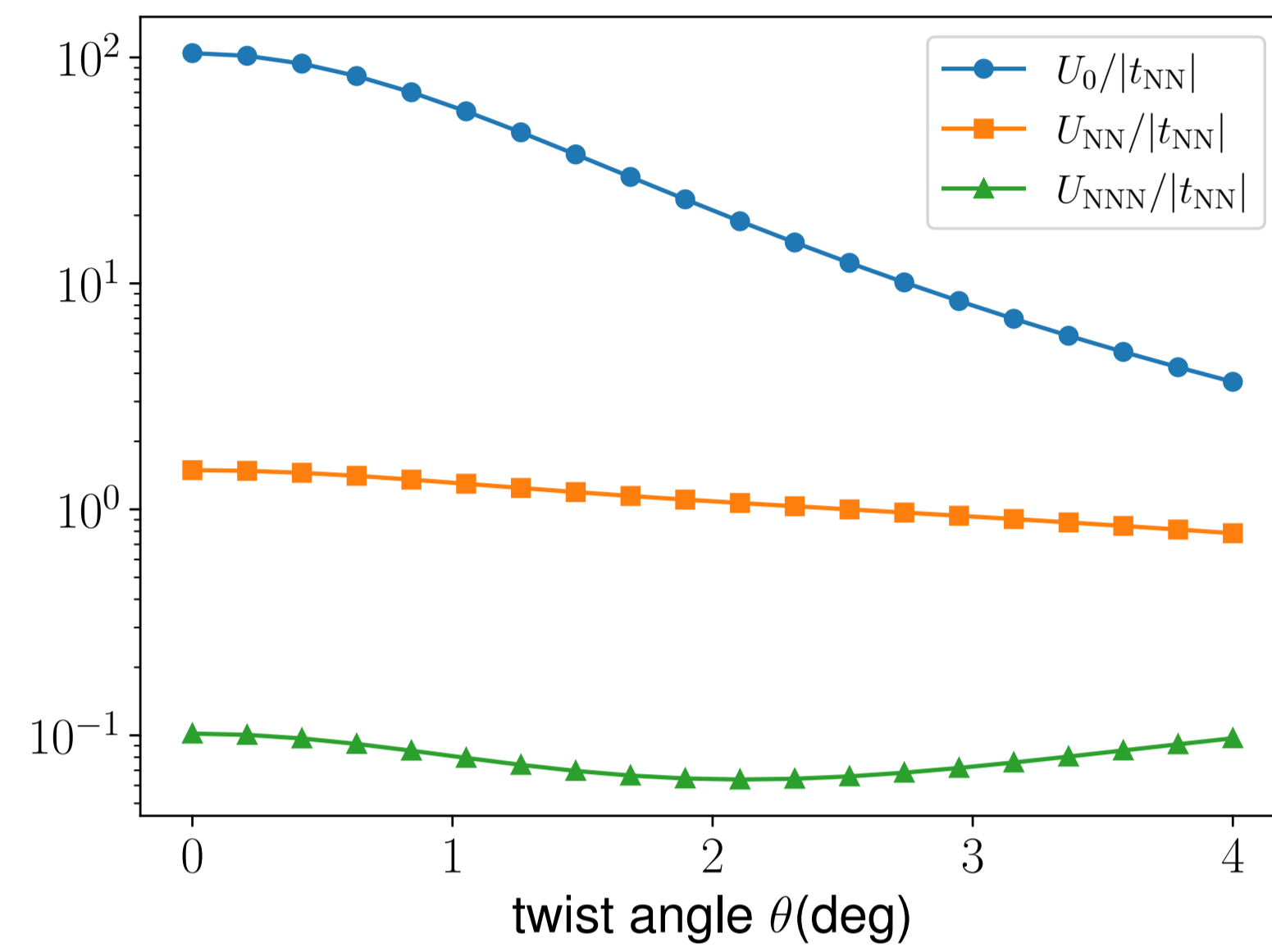
Two-dimensional bilayer materials combine deep subwavelength lattice spacing ($a_M/\lambda \sim 0.01$) with strong interactions and offer an interesting platform for (many-body) quantum optics!

Effective Bose-Hubbard Model

- The lowest band of hybridised moiré excitons can be described in terms of effective (long-range) Bose-Hubbard model on a triangular lattice [1]
- The hopping and interaction strength are determined by twist angle



$$H_{ex} = \sum_{\langle i,j \rangle} t_{NN} x_i^\dagger x_j + \sum_j U_0 x_j^\dagger x_j^\dagger x_j x_j + \sum_{i,j} U_{ij} x_i^\dagger x_j^\dagger x_j x_i$$



- Nearest-neighbour hopping is dominant hopping contribution
- On-site and nearest-neighbour interaction are relevant
- $|t_{NN}| \approx 0.1 - 2.6$ meV
- $U_0 \approx 8.9 - 9.6$ meV

Exciton-Polariton Scattering and Nonlinearity

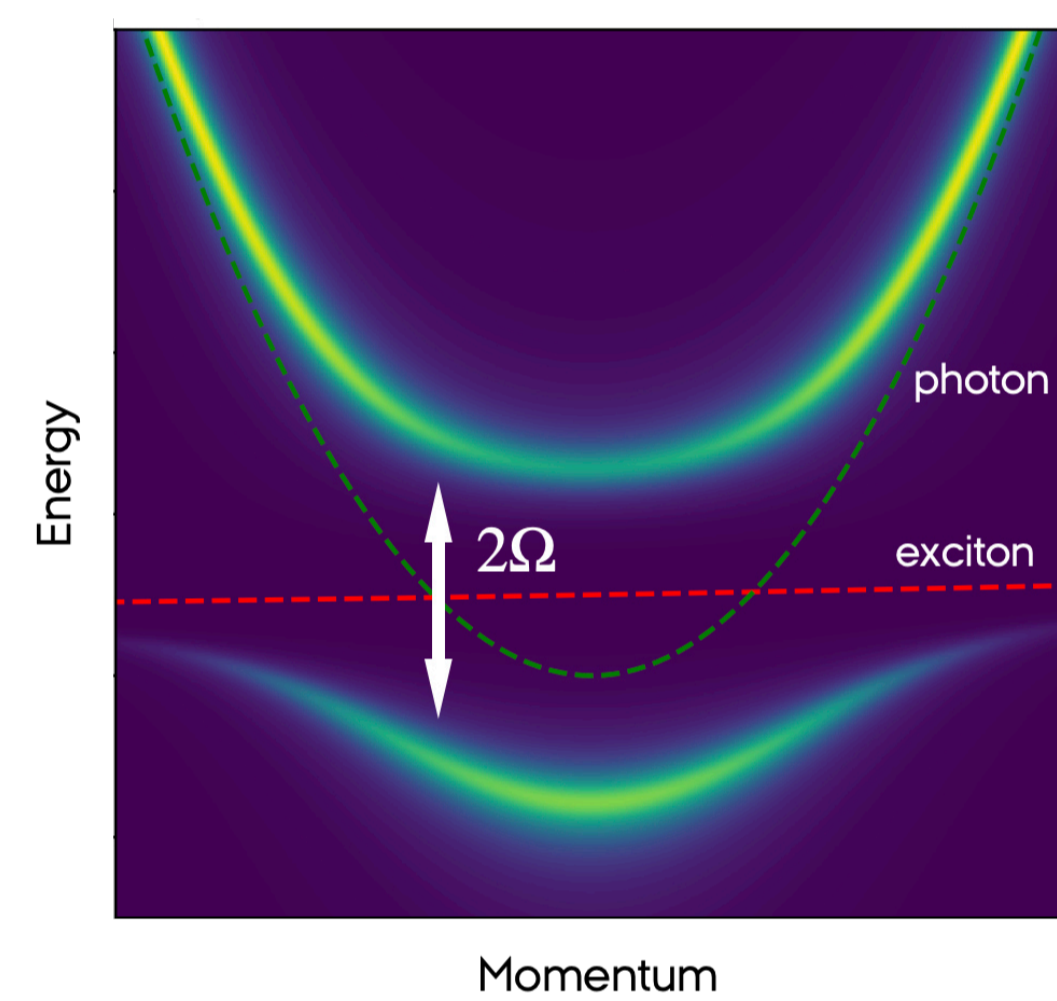
Exciton-Polaritons

$$H = H_{ex} + \sum_{\mathbf{k} \in \text{mBZ}} \left(\frac{\hbar^2 |\mathbf{k}|^2}{2m_c} + \delta \right) a_{\mathbf{k}}^\dagger a_{\mathbf{k}}$$

photon dynamics in 2D plane

$$+ \Omega (a_{\mathbf{k}}^\dagger x_{\mathbf{k}} + x_{\mathbf{k}}^\dagger a_{\mathbf{k}})$$

exciton-photon coupling



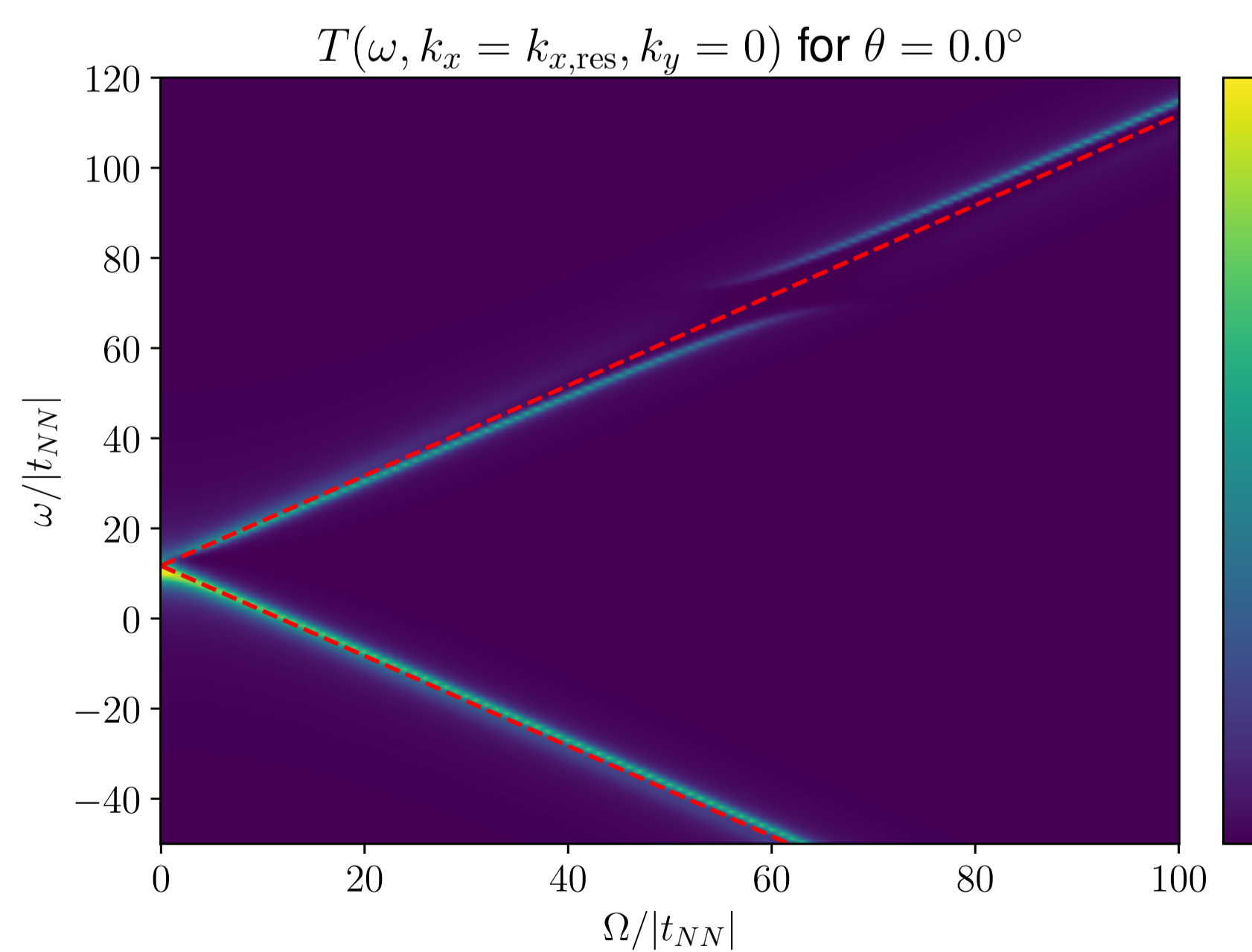
- Use exciton interaction to induce effective photon-photon interaction

Photon Nonlinearity

- Photon transmission is proportional to photonic part of the (many-body) Green's function $T(\omega, \mathbf{k}) \propto |G_c(\mathbf{k}, \omega)|^2$

$$G_c^{-1}(\mathbf{k}, \omega) = \omega - \left(\varepsilon_c(\mathbf{k}) - i\frac{\kappa}{2} \right) - \frac{\Omega^2}{\omega - (\varepsilon_X(\mathbf{k}) + n_X \Gamma_{XX} - i\frac{\gamma}{2})}$$

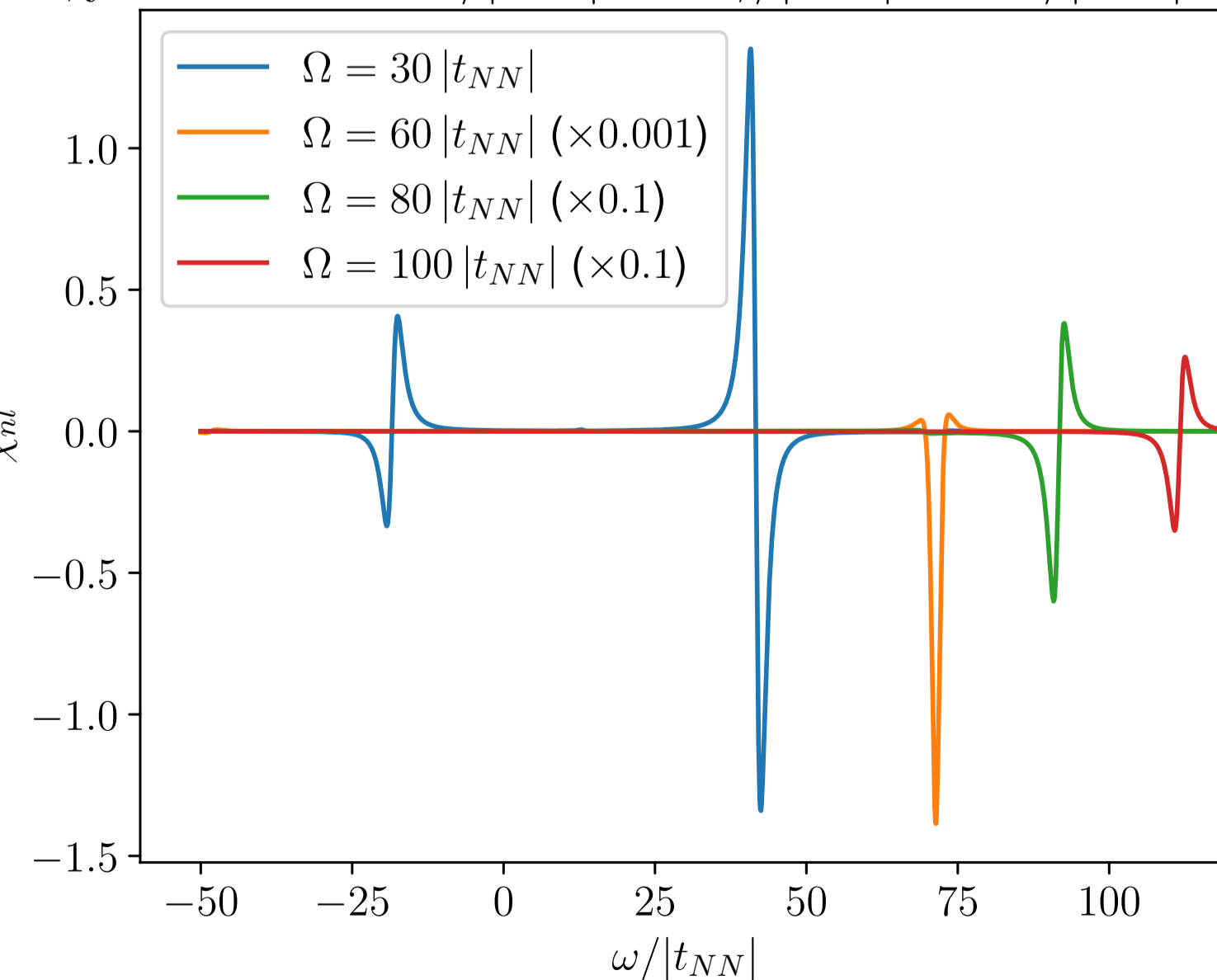
- Scattering resonances lead to strong interaction-induced shifts of the polariton transmission line



nonlinear transmission coefficient

$$T = T_0 (1 + \chi_{nl} I)$$

χ_{nl} for $\theta = 0.0^\circ$, $\kappa/|t_{NN}| = 5$, $\gamma/|t_{NN}| = 1$, $\delta/|t_{NN}| = 10$

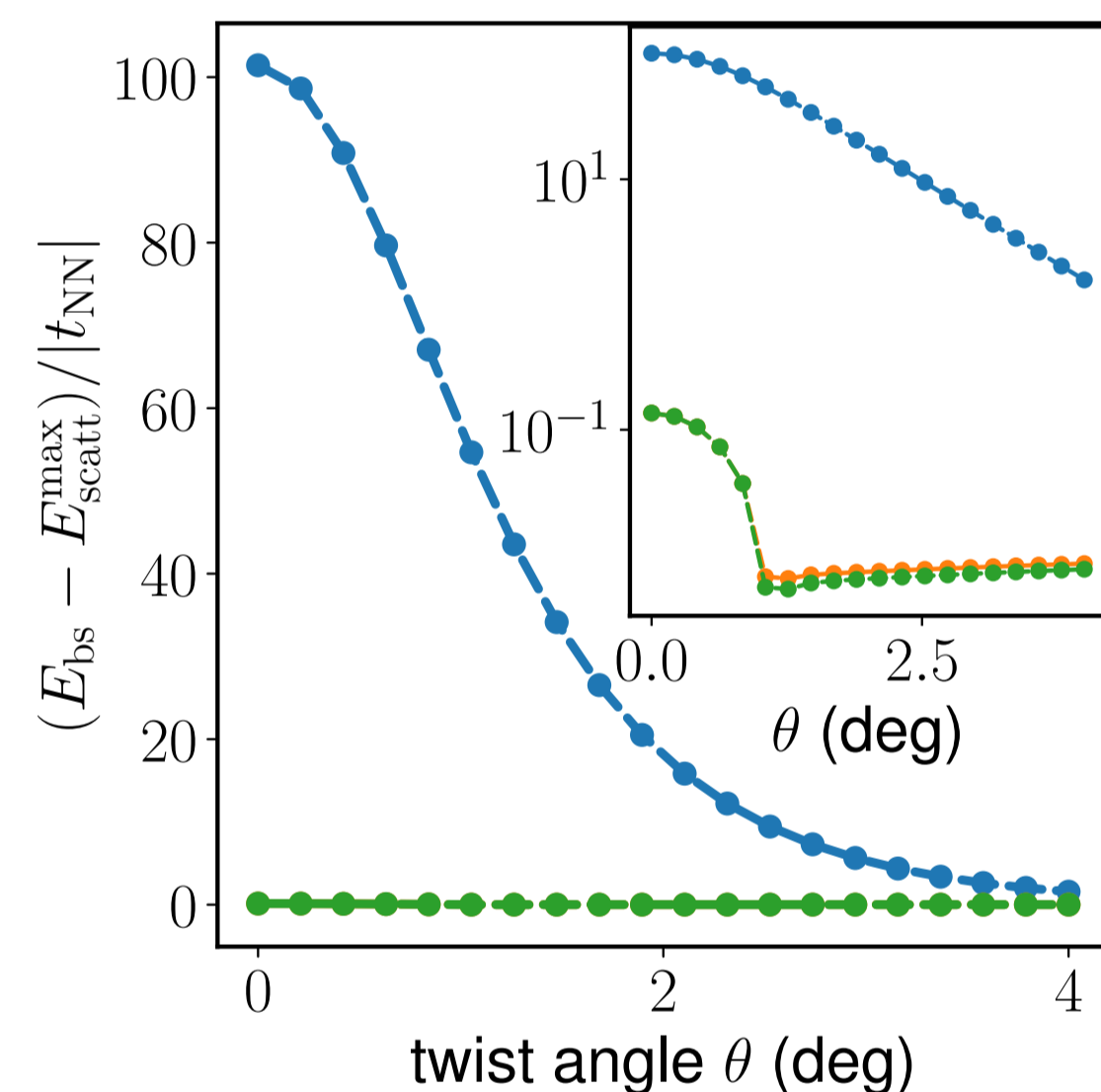
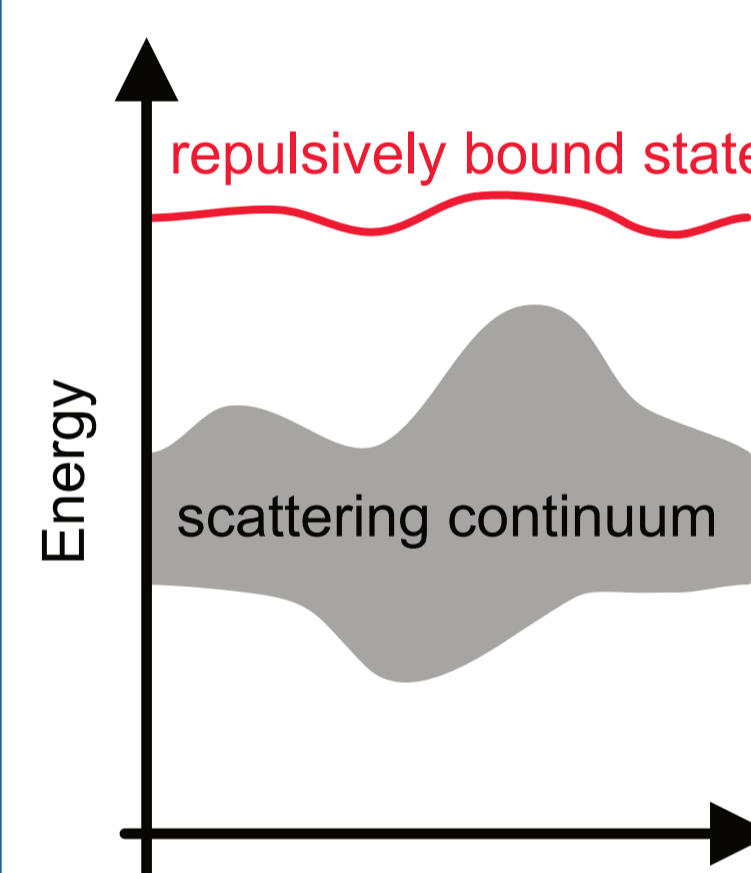


- The upper polariton experiences splitting on resonance due to coupling to repulsively bound state, the lower polariton does not experience this splitting
- Below (above) resonance, the shift is attractive (repulsive)

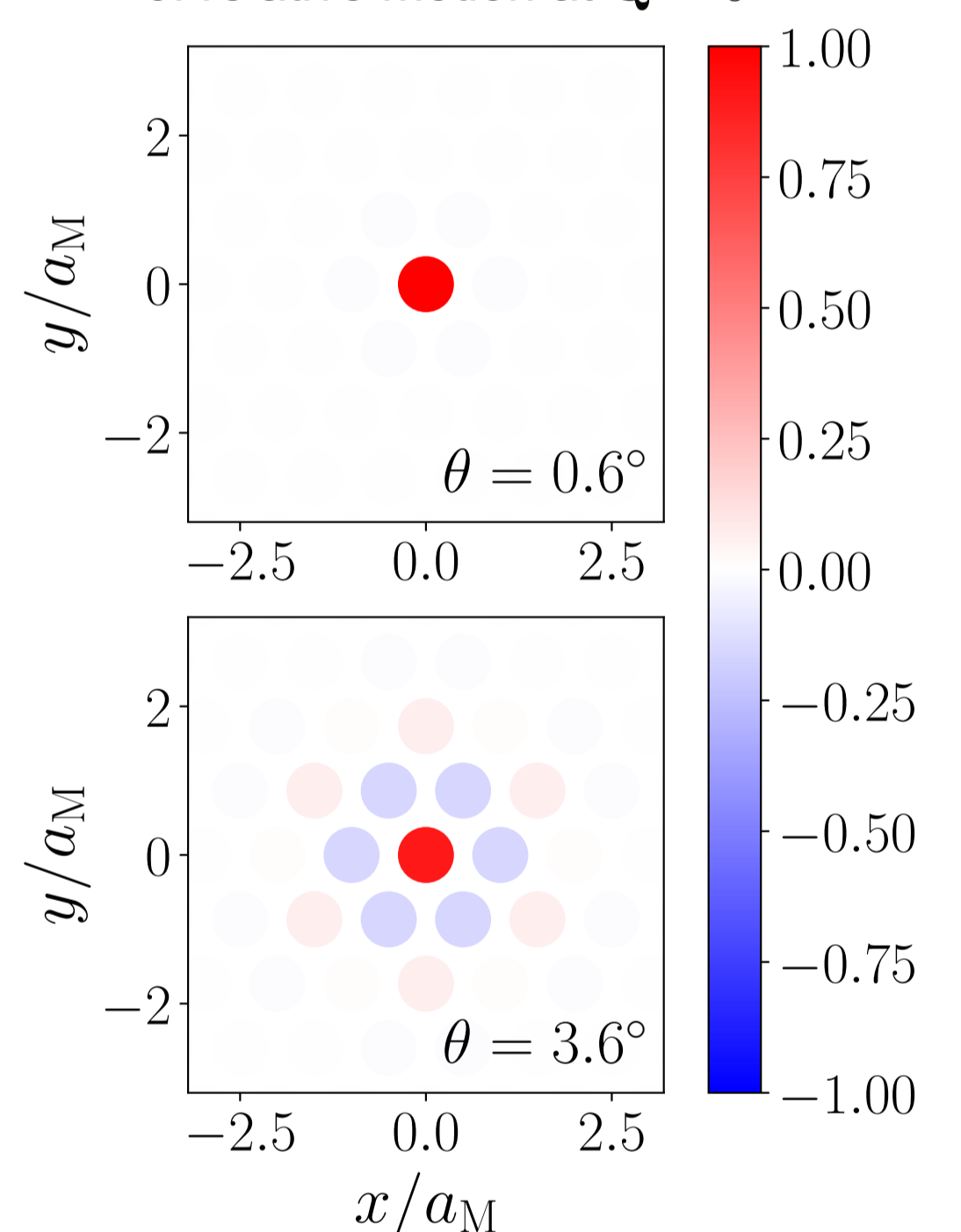
Exciton-Exciton Scattering - Bound States

Repulsively Bound Pairs in a Moiré Lattice

- Repulsively bound pair is unable to convert interaction (binding) energy to kinetic energy due to finite width of Bloch band [2]
- In one and two dimensions, repulsively bound pairs exist for arbitrarily small (on-site) interactions strengths



bound state wave functions of relative motion at $\mathbf{Q} = 0$



- Additional (loosely) bound states appear below a critical twist angle (above critical interaction strength) due to long-range interactions

Modified Exciton-Polariton Scattering

- Exciton scattering is modified due to the light-matter coupling (scattering in the presence of polariton background)
- We restrict the analysis to on-site interactions (still due to dipolar interactions within one moiré cell!)

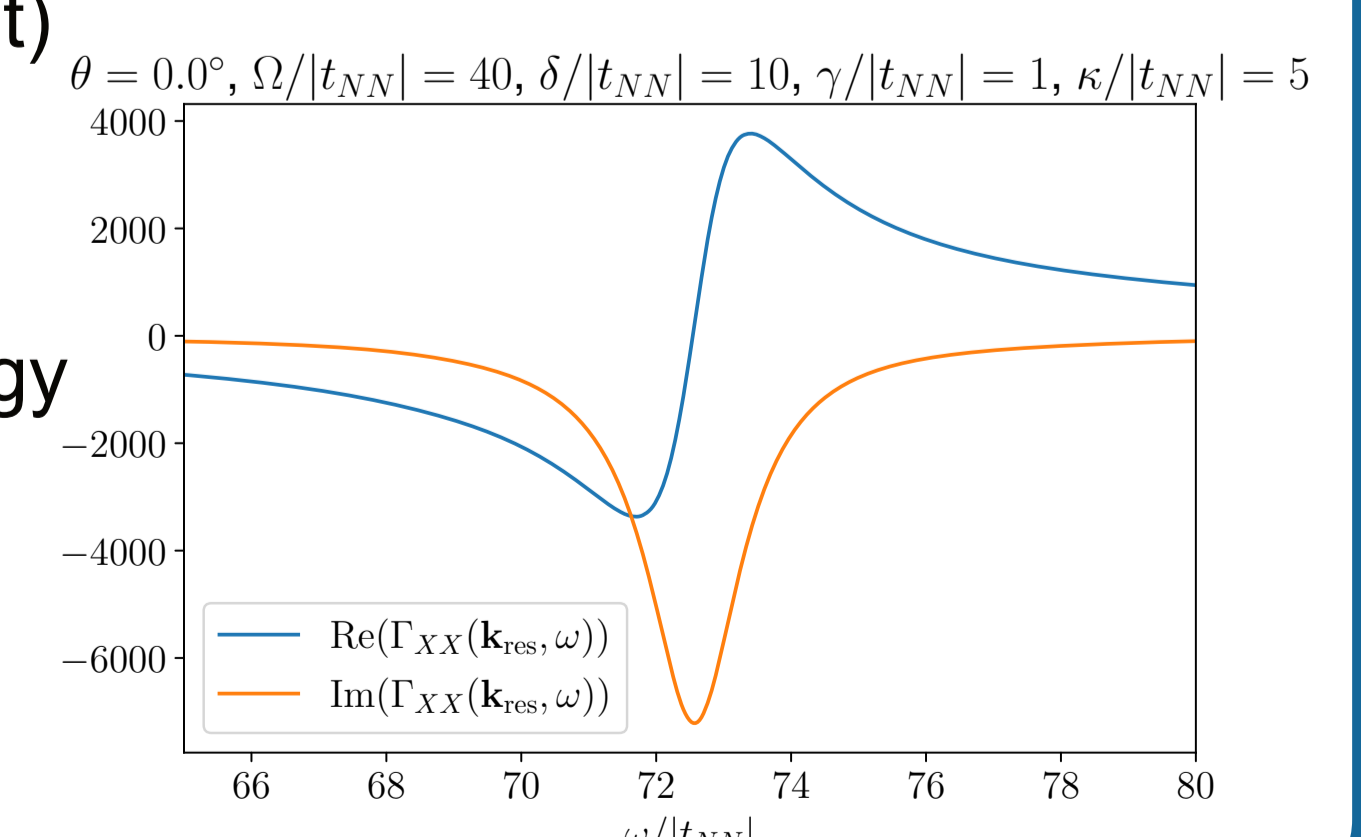
$$\Gamma_{XX}(\mathbf{Q}, \omega) = \frac{U_0}{1 - U_0 \Pi(\mathbf{Q}, \omega)}$$

$$\Pi(\mathbf{Q}, \omega) = \sum_{\mathbf{q} \in \text{mBZ}} \sum_{\alpha, \beta = \text{LP, UP}} \frac{\nu_X^\alpha(\mathbf{Q}/2 + \mathbf{q}) \nu_X^\beta(\mathbf{Q}/2 - \mathbf{q})}{\omega - \varepsilon_\alpha(\mathbf{Q}/2 + \mathbf{q}) - \varepsilon_\beta(\mathbf{Q}/2 - \mathbf{q}) + i\eta}$$

$\varepsilon_{\text{LP/UP}}(\mathbf{k})$: energy of lower/upper polariton

$\nu_X^{\text{LP/UP}}(\mathbf{k})$: (hybrid) exciton contribution to lower/upper polariton (Hopfield coefficient)

- Scattering matrix has narrow resonance at the bound state energy



References

- [1] A. Julku, Phys. Rev. B **106**, 035406 (2022)
 [2] K. Winkler et al., Nature **441**, 853 (2006)



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